

С 346

R-79

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ОБЪЕДИНЕННЫЙ ИНСТИТУТ ЯДЕРНЫХ ИССЛЕДОВАНИЙ

R1 - 3129

4796/1

A.H. Rosenfeld, Paul Soding, W.J. Willis, C.G. Wohl
L.R. Price, Matts Roos, A. Barbaro-Galtieri, W.J. Podolsky,

Data on Particles and Resonant States

January 1967

CLEBSCH-GORDAN COEFFICIENTS AND SPHERICAL HARMONICS

Note: A $\sqrt{\quad}$ is to be understood over every coefficient; e.g., for $-8/15$ read $-\sqrt{8/15}$.

Notation: $\begin{matrix} J & J & \dots \\ M & M & \dots \\ m_1 & m_2 & \dots \\ m_1 & m_2 & \dots \\ \dots & \dots & \dots \\ m_1 & m_2 & \dots \end{matrix}$ Coefficients

$Y_0^0 = \sqrt{\frac{1}{4\pi}} \cos \theta$
 $Y_1^0 = -\sqrt{\frac{3}{8\pi}} \sin \theta \cos \theta$
 $Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2} \right)$
 $Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$
 $Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$

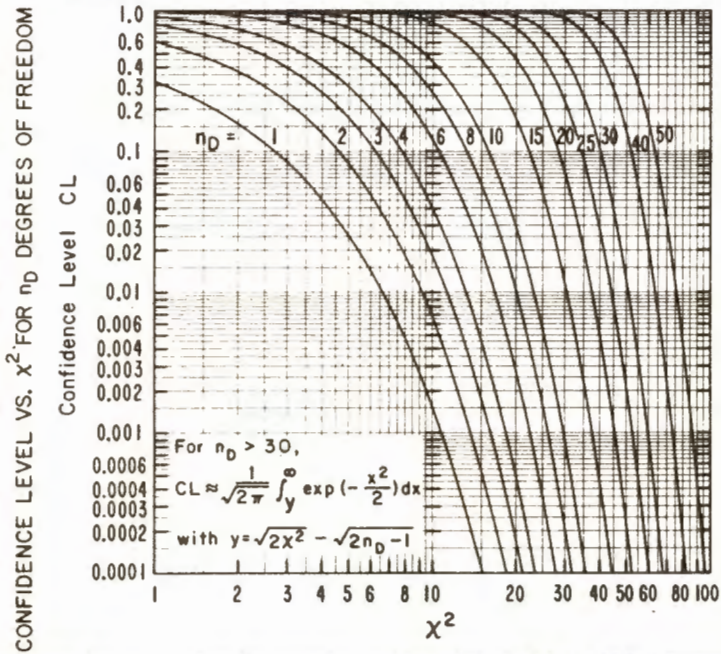
$Y_l^{-m} = (-1)^m Y_l^{m*}$

MU-28363

Atomic and Nuclear Properties of Materials

Material	Z	A	Cross Section σ barns	Collision Length L_{coll}		Minimum $-dE/dx^*$		Radiation Length L_{rad}		Density ρ g cm $^{-3}$
				cm	cm	MeV g $^{-1}$ cm 2	MeV cm $^{-1}$	g cm $^{-2}$	cm	
H $_2$	1	1.01	0.063	26.5	374	4.13	0.292	58.0	819	0.0708 ^b
D $_2$	1	2.01	0.100	33.4	202	0.342		116	703	0.168 ^b
He	2	4.00	0.16	42.0	336	1.94	0.242	85.4	683	0.125 ^b
Li	3	6.94	0.23	50.4	94.1	1.69	0.902	78.7	148	0.534
Be	4	9.01	0.28	55.0	29.9	1.60	2.96	63.7	34.7	1.848
C	6	12.01	0.33	60.4	1	1.78	1	42.4	1	=1.55 ^f
N $_2$	7	14.01	0.36	63.6	78.9	1.85	1.46	37.8	46.7	0.884 ^b
Al	13	26.98	0.57	79.2	29.3	1.62	4.37	24.0	8.9	2.70
Fe	26	55.85	0.92	101.2	12.8	1.48	11.6	13.9	1.8	7.87
Cu	29	63.54	1.00	105.4	11.8	1.44	12.9	12.0	1.34	8.96
Sn	50	118.69	1.55	129.7	37.8	1.28	9.4	8.89	1.22	7.31
W	74	183.85	2.02	150.8	7.80	1.17	22.6	6.89	0.36	19.3
Pb	82	207.19	2.20	156.2	13.8	1.13	12.8	6.52	0.58	11.35
U	92	238.03	2.42	163.6	16.63	1.09	20.6	6.13	0.12	18.95
Air				64.6	53620	1.81	0.0022	36.5	30290	0.001205 ^b
Freon (CF $_3$ Br)				87.1	498.0	1.52	2.3	16.6	111	1.5
H $_2$ (bubble chamber, 27°C)				26.5	442	4.13	0.248	58.0	970	=0.084 ^b
H $_2$ O				57.2	57.2	2.03	2.03	35.7	35.7	1.00
Urad Emulsion				103.0	27.0		5.49	11.2	2.91	3.815
LiF				63.8	24.2	1.69	4.46	19.0	14.8	2.64
Mylar (C $_5$ H $_8$ O $_2$)				59.1	42.8	1.91	2.64	19.6	28.7	1.18
Nal				119.0	32.4	1.32	4.84	9.58	2.61	3.17
Polyethylene (CH $_2$)				51.0	95.5	2.09	11.92	44.1	448	0.92
Polystyrene (CH)				54.9	95.3	2.03	12.14	43.4	44.1	1.05
Propane (C $_3$ H $_8$ bubble chamber)				48.9	119.3	2.28	0.935	44.6	109	0.41

- a. $\sigma = \sigma_{\text{natural}} = \pi^2 (h/m_0 c)^2 \times A^2 / \lambda^2 = 62.8 \text{ mb} \times A^2 / \lambda^2$
 b. $L_{\text{coll}} = A / (N \sigma_{\text{natural}}) = 26.5 \text{ g cm}^{-2} \times A^2 / \lambda^2$
 c. From W. H. Barkas and M. J. Berger, *Tables of Energy Losses and Ranges of Heavy Charged Particles*, NASA SP-1013 (1964)
 d. Mainly from *High Energy and Nuclear Physics Data Handbook*, W. Galbraith and W. S. C. Williams, Ed. (N. I. R. N. S., Rutherford Lab., Chilton, Didcot, Berks.) 1964
 e. boiling at 1 atmosphere f. density variable g. at 20°C
 h. May vary by about a 1%, depending on operation conditions



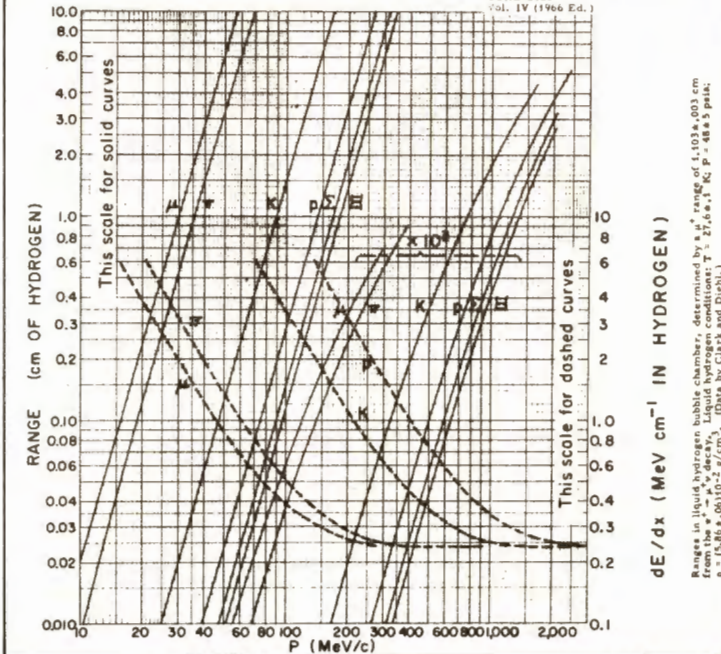
STIRLING'S APPROXIMATION
 $\sqrt{2\pi n} \left(\frac{n}{e}\right)^n < n! < \sqrt{2\pi n} \left(\frac{n}{e}\right)^n \left(1 + \frac{1}{12n}\right)$

Gaussianlike Distributions
 For $n > -1$ but not necessarily integral:
 $\int_0^\infty x^{2n+1} \exp\left[-\frac{x^2}{2\sigma^2}\right] dx = 2^n n! \sigma^{2n+2} \left(\frac{1}{2}\right)! \approx \sqrt{2}$

Relation between standard deviation σ and mean deviation \bar{x} :
 $2\sigma^2 = \pi \bar{x}^2$; $\sigma = 1.4826$ probable error.
 Odds against exceeding one standard deviation = 2.15:1;
 two, 21:1; three, 370:1; four, 16,000:1;
 five, 1,700,000:1

MULTIPLE COULOMB SCATTERING
 The rms projected angle θ due to multiple Coulomb scattering (only) of a particle of charge z , momentum P , velocity V is
 $\theta_{\text{proj}} = z \frac{15(\text{MeV})}{PV(\text{MeV})} \sqrt{\frac{L}{(\text{rad})}} (1 + \epsilon)$ radians;
 L = Length in scatterer;
 For $L \geq 1/10 L(\text{rad})$ ϵ is generally $< 1/10$. The distribution of θ is not truly Gaussian. The rms projected displacement y on traversing an absorber of thickness L is
 $y_{\text{rms}} = L \theta_{\text{proj}} / \sqrt{3}$

RADIOACTIVITY
 1 curie = 3.7×10^{10} disintegrations/sec
 $1 \text{ R} = 87.8 \text{ ergs/g air} = 5.49 \times 10^7 \text{ MeV/g air}$
 Fluxes (per cm 2) to liberate 1 R in carbon:
 3×10^7 minimum ionizing singly charged particles
 0.9×10^9 photons of 1 MeV energy.
 (These fluxes are actually correct to within a factor of two for all materials.)
 Natural background: 100 mR/year
 "Tolerance" 100 millirem/week [Note, 1 R may produce up to 10 "Rem." (R equivalent for man), depending on type of radiation.]



Particle or resonance	$I(J^P)$ = estab.	Beam π, K (BeV) (BeV/c)	Mass (MeV)	Γ (MeV)	$M^2 \approx \Gamma M$ (BeV ²)	Partial decay modes				
						Mode	Fraction (%)	Q (MeV)	p or p_{max}^2 (MeV/c)	$4\pi k^2$ (mb)
p	$1/2(1/2^+)$		938.3 939.6		0.880 0.883		See Table S			
$N(1400)$	$1/2(1/2^+)$ P_{11}	$T=0.43$ $p=0.55$	-1400^a	~ 200	1.96 ± 0.28	$N\pi$	70	322	367	36.3
$N(1525)$	$1/2(3/2^+)$ D_{13}	$T=0.62$ $p=0.75$	1525^a	105	2.33 ± 0.16	$N\pi$ $N\pi\pi$ $[\Delta(1236)\pi]^a$ [-20]	65 35 149	447 308 229	460 414 229	23.2
$N(1570)$	$1/2(1/2^+)$ S_{11}	$T=0.69$ $p=0.82$	1570^a	130	2.46 ± 0.20	$N\pi$ $N\eta$	~ 30 ~ 70	492 82	491 242	20.3
$N(1670)$	$1/2(5/2^+)$ D_{15}	$T=0.87$ $p=1.00$	1670^a	140	2.79 ± 0.23	$N\pi$ $N\pi\pi$ $[\Delta(1236)\pi]^a$ ΔK $N\eta$	40 dominant ^a [?] ^a small small	592 453 294 200	560 526 357 200	15.6
$N(1688)$	$1/2(5/2^+)$ F_{15}	$T=0.90$ $p=1.03$	1688^a	110	2.85 ± 0.19	$N\pi$ $N\pi\pi$ $[\Delta(1236)\pi]^a$ ΔK $N\eta$	65 dominant ^a [?] ^a small small	610 474 312 75 200	572 538 372 231 388	14.9
$N(1700)^c$	$1/2(1/2^+)$ S_{11}	$T=0.92$ $p=1.05$	1700^a	240	2.89 ± 0.41	$N\pi$	100	622	580	14.5
$N(2190)$	$1/2(7/2^+)$	$T=1.94$ $p=2.07$	2190	200	4.80 ± 0.14	$N\pi$ ΔK	30 ?	1112 577	888 710	6.21
$N(2650)$	$1/2(11/2^+)^b$	$T=3.12$ $p=3.26$	2650 ± 10	~ 300	7.02 ± 0.80	$N\pi$ ΔK	7 ?	1572 1037	1154 1022	3.67
$N(3030)^c$	$1/2(15/2^+)^b$	$T=4.26$ $p=4.40$	3030	400	9.18 ± 1.21	$N\pi$	0.7	1972	1377	2.62
$\Delta(1236)$	$3/2(3/2^+)$ F_{33}	$T=0.195$ $p=0.304$ $m_0 - m_{++} = 0.45 \pm 0.85$	(++) 1236.0 ± 0.6 $m_- - m_{++} = 7.9 \pm 6.8$	120 ± 2	1.53 ± 0.15	$N\pi$ $N\pi^+\pi^-$	100 0	158 18	231 89	91.9
$\Delta(1670)$	$3/2(1/2^+)$ S_{31}	$T=0.87$ $p=1.00$	1670^a	~ 180	2.79 ± 0.30	$N\pi$ $N\pi\pi$	40 ?	592 453	560 526	15.6
$\Delta(1920)$	$3/2(7/2^+)$	$T=1.35$ $p=1.48$	1920	200	3.69 ± 0.38	$N\pi$ ΣK	50 seen	842 229	722 423	9.37
$\Delta(2420)$	$3/2(11/2^+)^b$	$T=2.51$ $p=2.65$	2423 ± 10	~ 275	5.87 ± 0.67	$N\pi$ ΣK	10 ?	1345 732	1024 830	4.66
$\Delta(2850)$	$3/2(15/2^+)^b$	$T=3.71$ $p=3.85$	2850 ± 12	~ 300	8.12 ± 0.86	$N\pi$	3	1772	1266	3.09
$\Delta(3230)^c$	$3/2(19/2^+)^b$	$T=4.94$ $p=5.08$	3230	440	10.4 ± 1.4	$N\pi$	0.6	2152	1475	2.24
$Z_0(1865)^c$	$0(1/2^-)$	$p=1.15$	K^+p 1863	150	3.47 ± 0.28	NK	55 (if $J=1/2$)	432	579	14.6
Λ	$0(1/2^+)$		1115.6		1.24		See Table S			
$\Lambda(1405)^d$	$0(1/2^+)$	$p < 0$	K^+p 1405	35	1.97 ± 0.05	$\Sigma\pi$	100	68	142	
$\Lambda(1520)$	$0(3/2^-)$	$p=0.392$	1518.8 ± 1.5	16 ± 2	2.31 ± 0.02	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi\pi$	$S=1.7^*$ $\begin{cases} 39\pm 5 & 81 \\ 51\pm 6 & 182 \\ 10\pm 2 & 124 \end{cases}$	235 258 251	83.6	
$\Lambda(1670)^a$	$0(1/2^+)$	$p=0.74$	1670	18	2.79 ± 0.03	$\Lambda\eta$ $N\bar{K}$	$K^+p \rightarrow \Lambda\eta$ seen	6 233	66 410	28.5
$\Lambda(1700)$	$0(3/2^-)$	$p=0.80$	1700 ± 10	40 ± 10	2.89 ± 0.07	$N\bar{K}$ $\Sigma\pi$	20 seen	263 363	439 411	25.0
$\Lambda(1820)$	$0(5/2^+)$	$p=1.06$	1819.5 ± 3.5	83 ± 8	3.31 ± 0.15	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $\Lambda\eta$	70 11 18 ~ 1	382 482 295 155	541 502 362 349	16.5
$\Lambda(2100)$	$0(7/2^-)$	$p=1.68$	2100	160	4.41 ± 0.34	$N\bar{K}$ $\Sigma\pi$	29 seen	663 763	748 699	8.68
$\Lambda(2340)$	$0(1/2^-)$	$p=2.27$	2340 ± 20	105	5.48 ± 0.25	$N\bar{K}$ seen in g (total)	10 L (if $J=9/2$)	903	907	5.92
Σ	$1(1/2^+)$		(+)1189.5 (0)1192.6 (-)1197.4		1.41 1.42 1.43		See Table S			
$\Sigma(1385)$	$1(3/2^+)$	$p < 0$	K^+p (+)1382.2 ± 0.9 $S=1.6^*$ $S=2.1^*$ $S=4.9^*$ (-)1388.0 ± 3.0 $S=3.7^*$		1.92 ± 0.05	$\Lambda\pi$ $\Sigma\pi$	91 ± 3 9 ± 3	130 48	208 117	
$\Sigma(1660)^a$	$1(3/2^-)$	$p=8.72$	1660	50	2.76 ± 0.08	$\Lambda(1405)\pi$ $\Sigma\pi$ $\Lambda\pi$ $N\bar{K}$	large ? ? small	115 323 405 223	197 379 439 400	29.9
$\Sigma(1770)$	$1(5/2^-)$	$p=0.95$	1768 ± 4 $S=1.5^*$	89 ± 12 $S=2.0^*$	3.13 ± 0.16	$N\bar{K}$ $\Lambda\pi$ $\Sigma(1520)\pi$ $\Sigma(1385)\pi$ $\Sigma\eta$ $\Sigma\pi$	49 17 19 12 2 < 1	331 517 110 243 27 431	498 520 192 316 143 463	19.4
$\Sigma(1910)^c$	$1(5/2^+)$	$p=1.25$	1910 ± 10	60	3.65 ± 0.11	$N\bar{K}$ $\Lambda\pi$ $\Sigma\pi$	8 10 3	473 655 573	612 619 568	12.9
$\Sigma(2035)$	$1(7/2^+)$	$p=1.53$	2035 ± 15	160	4.14 ± 0.33	$N\bar{K}$ $\Lambda\pi$ $\Sigma\pi$	16 25 seen	598 784 698	703 703 655	9.83
$\Sigma(2260)^c$	$1(1/2^-)$	$p=2.06$	2260 ± 20	180	5.11 ± 0.41	$N\bar{K}$ seen in g (total)	14 L (if $J=9/2$)	821	855	6.66
Ξ	$1/2(1/2^+)$		(0)1314.7 (-)1321.2		1.73 1.75		See Table S			
$\Xi(1530)$	$1/2(3/2^+)$	p -wave	(0)1528.9 ± 1.1 (-)1533.8 ± 1.9	7.3 ± 1.7	2.34 ± 0.01	$\Xi\pi$	100	69	145	
$\Xi(1815)$	$1/2(1/2^-)$		1815 ± 3	16 ± 8 $S=2.2^*$	3.29 ± 0.03	$\Delta\bar{K}$ $\Xi\pi$ $\Xi\pi\pi$	~ 65 ~ 10 ~ 25	202 354 215	391 409 351	
$\Xi(1930)$	$1/2(1/2^-)$		1933 ± 16	140 ± 35	3.74 ± 0.27	$\Xi\pi$ $\Delta\bar{K}$	seen seen	472 320	501 504	
Ω^-	$0(3/2^+)$		1674		2.80		See Table S			

$N^*_{1/2}$
 $N^*_{3/2}$
 Z^*_0
 Y^*_0
 Y^*
 $\Xi^*_{1/2}$

a. See note in data listings.
 b. JP assignment based on straight-line Regge-trajectory recurrence hypothesis and supported by fits to π elastic scattering at 180°. See note following data listings.
 c. Evidence for the existence of the effect and/or for its interpretation as a resonance is open to some question.
 d. Λ is the minimum momentum that any of the particles in the decay modes into π particles has been found to have.
 e. Square brackets indicate a sub-reaction of the previous unbracketed decay mode.

BARYON STATES, UCRL-8030 Rev.

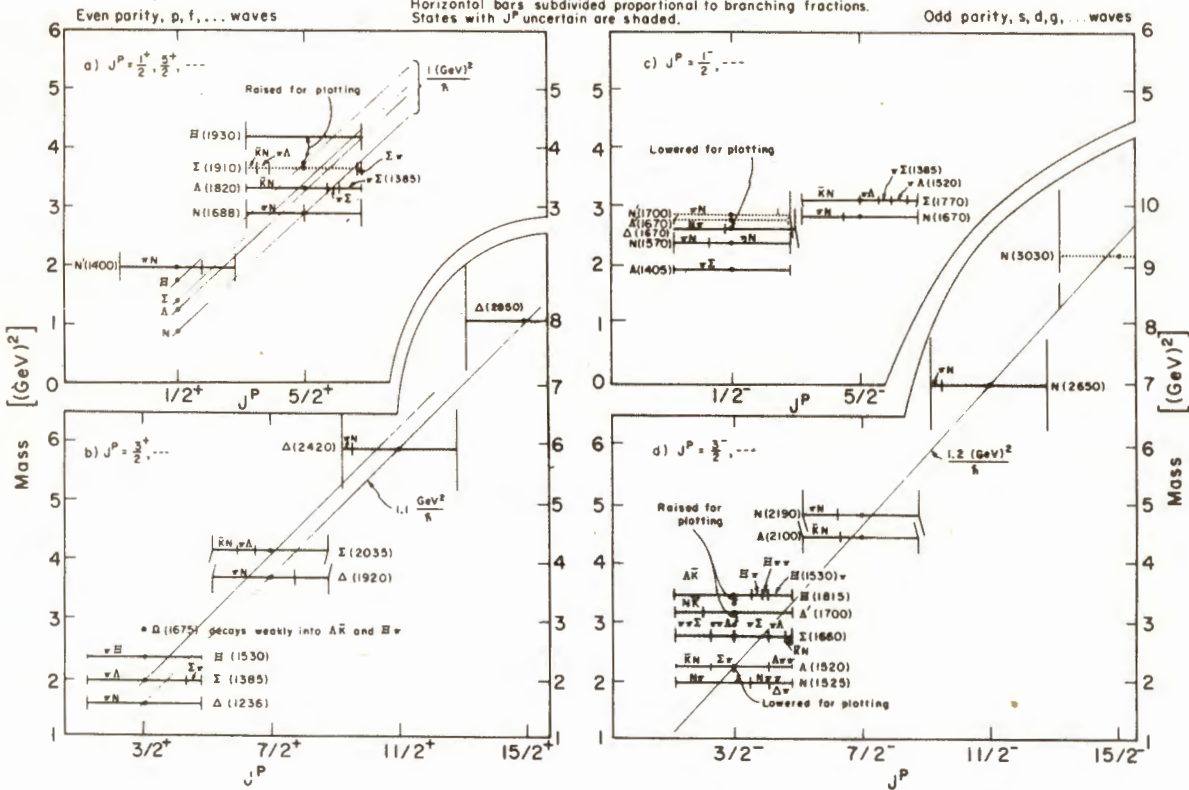
Jan. 1967

Possibly questionable states are dashed.

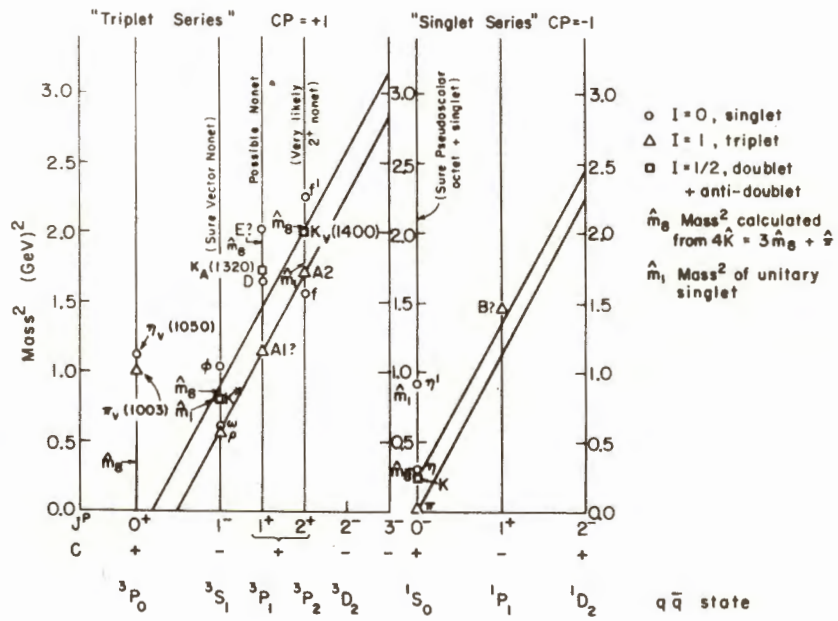
Vertical bars : $\Gamma (m^2) = 2m\Gamma =$ full width.

Horizontal bars subdivided proportional to branching fractions.

States with J^P uncertain are shaded.



MUB-12316



MUB-8630

$$\mu_{dp} = m_p dE = m_p v_{beam, lab} dp = m_p dp$$

Special Relativity

Notation: 4-vector in c.m. $p = (w, \vec{p})$; in lab $P = (W, \vec{P})$, $T = W - m$.
 Solid-angle element $d\omega = 2\pi d\cos\theta$; $d\Omega = 2\pi d\cos\Theta$.
 $p^2 = w^2 - \vec{p}^2 = m^2$ is an invariant. Cross section σ is invariant.

Lorentz Transformation

$$\begin{pmatrix} P_x \\ P_y \\ P_z \\ P_0 \end{pmatrix} = \begin{pmatrix} \gamma & 0 & 0 & 0 \\ 0 & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} P_x \\ P_y \\ P_z \\ P_0 \end{pmatrix} \quad \text{If } \theta \text{ and } \Theta \text{ are measured with respect to the transformation axis } x.$$

$$\frac{P_x}{P_0} = \tan\theta = \frac{P \sin\Theta}{-\gamma W + \gamma P \cos\Theta} \quad (1)$$

If particle 1 is beam, 2 is target. $(W_1, \vec{P}_1) = (m_1, \vec{0})$ and $\vec{W} = (W_1 + m_2)/\sqrt{s}$, $\vec{P} = \gamma \vec{P}_1 / \sqrt{s}$, $|\vec{P}_1| = |\vec{P}_2| = \bar{m}_2 = |\vec{P}_1| m_2 / \sqrt{s}$.

$$W = W_1 + m_2, \quad \vec{P}^2 = 1 + T_1 / 2m_1 \quad (2)$$

Invariants. Notation: $t + z = t' + z'$.

$$s = (p_1 + p_2)^2 = m_1^2 + m_2^2 + 2(W_1 W_2 - \vec{P}_1 \cdot \vec{P}_2) \quad (3)$$

$$t = (p_1 - p_2)^2 = m_1^2 + m_2^2 - 2(W_1 W_2 - \vec{P}_1 \cdot \vec{P}_2) \quad (4)$$

$$u = (p_1 - p_2)^2 = (p_2 - p_1)^2 \quad (5)$$

$$\text{General relation: } s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2 \quad (6)$$

In lab system $P_2 = (m_2, \vec{0})$, and writing $W = m + T$

$$s = m^2 + m_2^2 + 2W_1 m_2 = (m_1 + m_2)^2 + 2T_1 m_2 \quad (3, lab)$$

$$t = m^2 + m_2^2 - 2W_1 m_2 = (m_1 - m_2)^2 - 2T_1 m_2 \quad (4, lab)$$

$$\text{In c.m. system } dt = 2|\vec{p}_1| |\vec{p}_2| d\cos\Theta \quad (4, cm)$$

For elastic scattering ($m_1 = m_1', m_2 = m_2'$), (4) in c.m. simplifies to

$$t = -2p^2 (1 - \tau \cos\theta) = -4p^2 \left\{ \begin{matrix} \sin^2\theta/2 \\ \cos^2\theta/2 \end{matrix} \right. \quad (4, el)$$

For elastic scattering, using (4, lab), (4, el), and (2),

$$T_2 = \frac{2p^2}{s} \sin^2\left(\frac{\theta}{2}\right) \quad \text{(useful for calculating } \delta\text{-ray energies).} \quad (7)$$

Two-Body States, Energies and momenta in c.m.

$$W = \frac{s + m_1^2 - m_2^2}{2\sqrt{s}}, \quad P_1 = P_2 = \frac{1}{4s} [s - (m_1 + m_2)^2] [s - (m_1 - m_2)^2] \quad (8)$$

3- and 4-Body States, Let $m_{ij}^2 = (p_i + p_j)^2$, etc., then

$$\sum_{i < j} m_{ij}^2 = \sum m_i^2 + m_{23}^2 = \text{const.} \quad (i, j = 1, 2, 3) \quad \text{(follows from (2))} \quad (9)$$

$$= 2\sum m_i^2 + m_{123}^2 = \text{const.} \quad (i, j, k = 1, 2, 3, 4) \quad (10)$$

$$\sum_{i < j < k} m_{ijk}^2 = \sum m_i^2 + 2m_{123}^2 = \text{const.}$$

R_n Invariant Volume in n-Body Momentum Space

A useful invariant is $\int d^3p \delta(p^2 - m^2) = \frac{d^3p}{2w} = \frac{p^2 d|\vec{p}| dw}{2w} = \frac{1}{2} |\vec{p}| dw dw$.

$$R_2 = v |\vec{P}_1| / \sqrt{s}, \quad R_3 = v^2 \int d\omega_1 d\omega_2 = (v^2/4s) \int dm_{12}^2 dm_{23}^2 \quad \text{(using c.m. quantities)}$$

Recurrence Relation for Factoring R_n (see e.g., Hagedorn, p. 93^b)

$$\text{Write } N = 1, 2, \dots, k, k+1, \dots, n \quad (R_n) \quad \text{then}$$

$$\text{as } N = \left. \begin{matrix} K, k+1, \dots, n \\ 1, 2, \dots, k \end{matrix} \right\} R_n = \int dM_{KK}^2 R_{n-k+1} \quad (11)$$

Cross Section (or Decay Rate)

For $1 + 2 \rightarrow n$ particles (or $1 \rightarrow n$ particles)

$$\sigma = |M|^2 R_n \quad \text{(or } \Gamma = |M|^2 R_n)$$

where M is an invariant matrix element.

F is (Moller's) invariant flux factor, $F = (p_1 p_2)^2 - m_1^2 m_2^2$.

In every system where \vec{P}_1 and \vec{P}_2 are collinear, $F = w_1 w_2 |\vec{v}_1 - \vec{v}_2|$ ($\vec{v} = \vec{P}/W$).

If particle 1 is beam, 2 is target ($\vec{P}_2 = 0$), $F = |\vec{P}_1| m_2 = |\vec{P}_1| v_2$ [cf. (2)].

The rate (= number per unit 4-dimensional volume $d^4r = dt d^3r$) is

$$d^4N/d^4r = \sigma p_1 p_2 |\vec{v}_1 - \vec{v}_2|, \quad p_i = \text{volume density of particles (i = 1, 2).}$$

A. R. Hagedorn, Relativistic Kinematics, (W. A. Benjamin, New York, 1964).

(P(LAB) (MEV/C))	-- INVARIANT MASS -- (MEV)				-- (P(CMS) (MEV/C))			
ep	wp	Kp	pp	ep	wp	Kp	pp	
0	939	1078	1432	1877	0	0	0	0
20	958	1079	1432	1877	20	17	13	10
40	977	1083	1433	1877	38	35	26	20
60	996	1089	1434	1877	56	52	39	30
80	1015	1097	1436	1878	74	68	52	40
100	1033	1105	1439	1875	91	85	65	50
120	1051	1116	1441	1880	107	101	78	60
140	1069	1127	1445	1882	123	117	91	70
160	1087	1139	1449	1883	138	132	104	80
180	1104	1152	1453	1885	153	147	116	90
200	1121	1165	1457	1887	167	161	129	99
220	1137	1178	1462	1889	182	175	141	109
240	1154	1192	1468	1892	195	189	153	119
260	1170	1206	1474	1894	209	202	166	99
280	1186	1219	1480	1897	222	215	178	138
300	1201	1233	1486	1900	234	228	198	148
320	1217	1247	1493	1902	247	241	201	168
340	1232	1261	1500	1906	259	253	213	167
360	1247	1274	1507	1910	271	265	224	177
380	1262	1288	1514	1913	282	277	235	186
400	1277	1302	1522	1917	294	288	247	196
420	1292	1315	1530	1921	305	300	258	205
440	1306	1329	1538	1925	316	311	268	214
460	1320	1342	1546	1929	327	322	279	224
480	1335	1356	1554	1933	337	332	290	233
500	1349	1369	1563	1938	348	343	300	242
520	1362	1382	1572	1943	358	353	310	251
540	1376	1395	1580	1947	368	363	321	260
560	1390	1408	1589	1952	378	373	331	269
580	1403	1421	1598	1957	388	383	341	278
600	1416	1434	1607	1962	397	393	350	287
620	1430	1446	1616	1967	407	402	360	296
640	1443	1459	1625	1973	416	412	370	304
660	1456	1472	1634	1978	425	421	379	313
680	1468	1484	1644	1984	434	430	388	322
700	1481	1497	1653	1989	443	439	397	330
720	1494	1509	1662	1995	452	448	406	339
740	1506	1521	1671	2001	461	457	415	347
760	1519	1533	1681	2007	470	466	424	355
780	1531	1545	1690	2013	478	474	433	364
800	1543	1557	1699	2019	486	482	442	372
820	1555	1569	1709	2025	495	490	450	380
840	1567	1580	1718	2031	503	499	459	388
860	1579	1592	1728	2037	511	507	467	396
880	1591	1604	1737	2043	519	515	475	404
900	1603	1615	1747	2049	527	523	483	412
920	1615	1627	1756	2056	535	531	492	420
940	1626	1638	1765	2062	542	538	500	428
960	1638	1649	1775	2069	550	546	508	435
980	1649	1661	1784	2075	558	554	515	443
1000	1660	1672	1794	2082	565	561	523	451
1020	1672	1683	1803	2088	573	569	531	459
1040	1683	1694	1812	2095	580	576	538	466
1060	1694	1705	1822	2102	587	583	546	473
1080	1705	1716	1831	2109	594	591	553	481
1100	1716	1726	1840	2115	601	598	561	488
1120	1727	1737	1850	2122	609	605	568	495
1140	1738	1748	1859	2129	616	612	575	502
1160	1748	1758	1868	2135	622	619	583	510
1180	1759	1769	1877	2142	629	626	590	517
1200	1770	1780	1887	2149	636	633	597	524
1220	1780	1790	1896	2156	643	639	604	531
1240	1791	1800	1905	2163	650	646	611	538
1260	1801	1811	1914	2170	656	653	618	545
1280	1812	1821	1923	2177	663	660	624	552
1300	1822	1831	1932	2184	669	666	631	559
1320	1832	1841	1941	2191	676	673	638	565
1340	1843	1851	1950	2198	682	679	645	572
1360	1853	1862	1959	2205	689	685	651	579
1380	1863	1872	1968	2212	695	692	658	585
1400	1873	1881	1977	2219	701	698	664	592
1420	1883	1891	1986	2226	708	704	671	599
1440	1893	1901	1995	2233	714	711	677	605
1460	1903	1911	2004	2240	720	717	684	612
1480	1912	1921	2013	2247	726	723	690	618
1500	1922	1930	2022	2254	732	729	696	624
1520	1932	1940	2031	2261	738	735	702	631
1540	1942	1950	2039	2268	744	741	709	637
1560	1951	1959	2048	2275	750	747	715	643
1580	1961	1969	2057	2282	756	753	721	650
1600	1970	1978	2065	2289	762	759	727	656
1620	1980	1988	2074	2296	768	765	733	662
1640	1989	1997	2083	2303	774	770	739	668
1660	1999	2006	2091	2311	779	776	745	674
1680	2008	2014	2100	2318	785	782	751	680
1700	2017	2021	2109	2324	790	787	756	686
1720	2027	2034	2117	2332	796	793	762	692
1740	2036	2041	2126	2339	802	799	768	698
1760	2045	2053	2134	2344	807	805	774	704
1780	2054	2062	2143	2353	813	810	779	710
1800	2064	2071	2151	2360	818	816	785	716
1820	2073	2080	2159	2367	824	821	791	721
1840	2082	2089	2168	2374	829	827	796	727
1860	2091	2098	2176	2381	835	832	802	733
1880	2100	2107	2184	2388	840	837	808	739
1900	2108	2115	2193	2395	845	843	813	744
1920	2117	2124	2201	2402	851	848	818	750
1940	2126	2133	2209	2409	856	853	824	756
1960	2135	2142	2217	2416	861	859	829	761
1980	2144	2150	2226	2423	867	864	835	767
2000	2153	2160	2234	2430	872	869	840	772
2020	2161	2168	2242	2437	877	874	845	778
2040	2170	2176	2250	2444	882	879	851	783
2060	2179	2185	2258	2451	887	885	856	789
2080	2187	2194	2266	2458	892	890	861	794