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E4-87-831

## V.A.Kuzmin, V.G.Soloviev

## EFFECT OF THE PARTICLE-PARTICLE INTERACTION ON THE GAMOW-TELLER $\beta^+$ -DECAY IN SPHERICAL NUCLEI

Submitted to "Письма в ЖЭТф"

Advances in describing some characteristics of collective lowlying states within the interacting boson model<sup>/1/</sup>, that takes into account particle-particle interactions, entailed studies of the role of particle-particle interactions. Taking interactions in the particlehole and particle-particle channels into account simultaneously allows one to fit the description of the two-neutrino double &-decay with the experiment<sup>/2/</sup>. This letter is devoted to the description of the Gamow-Teller  $\beta^+$ -decay of even spherical nuclei in the RPA taking into account interactions between quasiparticles in the particle-particle and particle-hole channels.

The QPNM Hamiltonian<sup>/3/</sup> contains particle-hole and particle-particle interactions between quasiparticles. From the general formulae of the model one can easily derive equations for the energies and wave functions of the Gamow-Teller 1<sup>+</sup> states in the RPA with particle-particle and particle-hole charge-exchange interactions with the constants  $G_{j}^{01}$  and  $\mathcal{B}_{j}^{02}$ , respectively. The creation operator of the charge-exchange phonon is written in the form

$$\Omega_{i}^{+} = \sum_{jp,jn} \left\{ \Psi_{jp,jn}^{i} A^{(jp,jn;10)} + \varphi_{jp,jn}^{i} A^{(jp,jn;10)} \right\},$$

where

$$A^{\dagger}(j_{p}j_{n};10) = \sum_{m_{p},m_{n}} \langle j_{p}m_{p}j_{n}m_{n}|10\rangle d^{\dagger}_{jp}m_{p}d^{\dagger}_{jn}m_{n}$$

Here  $d_{jm}$  is the quasiparticle creation operator,  $j_p m_p (j_n m_n)$  are quantum numbers of proton (neutron) single-particle states,  $\ell = 1, 2, 3, ...$  is the root number of the relevant secular equation.

The matrix element of the  $\mathfrak{g}^+$ -decay of the ground state of the doubly even nucleus with the wave function  $\Psi_o$  into the one-phonon 1<sup>+</sup>state of the doubly odd nucleus with the wave function  $\Omega_i^+ \Psi_o$  is

 $(\Psi_{0}^{*}\Omega_{i}H_{\beta}^{i}\Psi_{o}) = \sum_{i_{0},i_{0}} \langle j_{n} \| \Gamma_{\beta} \| j_{p} \rangle (\Psi_{j_{0}j_{n}}^{i_{0}} \mathcal{U}_{j_{n}} + \varphi_{j_{0}j_{n}}^{i_{0}} \mathcal{U}_{j_{0}} \mathcal{U}_{j_{n}}),$ 

where  $\langle jn || \int_{\mathcal{B}} || jp \rangle$  is the single-particle Gamow-Teller matrix element,  $\mathcal{U}_j$  and  $\mathcal{U}_i$  are the coefficients of the Bogolubov canonical transformation. The inclusion of the particle-particle interaction alongside with the particle-hole interaction leads to the increase of  $\mathcal{U}_{join}$  for the lowest states. Since the signs of the functions  $\mathcal{U}_{join}$  and  $\mathcal{U}_{join}$  are opposite, the increase in the absolute value of the particle-particle interaction constant leads to the suppression of the  $\mathbb{B}^+$  transition probabilities.

The results of calculations of  $\log \hat{ft}$  for the  $\beta^+$  transitions to the lowest 1<sup>+</sup> states of neutron-deficit nuclei calculated in the RPA and the relevant experimental sum values of log  $\tilde{ft}$ , defined as

 $(\tilde{\mathfrak{ft}})^{-1} = \sum_{\kappa} (\mathfrak{ft})_{\kappa}^{-1}$ 

of refs.  $(5^{-9})'$  are shown in the table. The same single-particle energies and wave functions of the Saxon-Woods potential and the pairing constants were used in the calculations as in ref. (4)'. The value of  $|\mathfrak{A}_{1}|'$  is 1.5 times as large as the value used in ref. (4)'' when studying nuclei of the stability zone. The constants of the particle-particle interaction  $G_{1}''=-0.2 \mathfrak{A}_{1}'''$  were used in the calculations. With increasing constant one can suppress twice the strength of B-transitions to the low-lying states without violating the applicability conditions for the RPA.

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Table.	Values	of log 🕂	for ß	-transitions	from	0	to	1+	states
		110 1	Alle	1 . + 1 3 x X		g.s.	1		

log ft	log ft(calculation)			
(experimenț)	$G_{1}^{01} = -0.2  \varkappa_{1}^{01}$	$G_1 = 0$		
3.4	3.5	3.1		
3.6	3.5	3.2		
1	3.7	3.4		
3.8	3.8	3.3		
3.3	3.3	3.0		
	(experiment) (3.4 3.6 (3.9 3.8 3.3	(experiment) (3.4) (3.4) (3.6) (3.6) (3.9) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8) (3.8)		

In more detail we consider the results of calculations of  $\beta^+$ -decay in <sup>148</sup>Dy. In comparison with the independent particle model the

inclusion of the particle-hole interaction results in a 2.65 decrease of the total  $\beta^{+}$  transition strength. The inclusion of the particle-particle interaction gives an additional twice suppression. At  $\mathcal{C}_{f}^{\sigma_{f}} = - \mathcal{Q}_{f} \mathcal{Q}_{f}^{\sigma_{f}}$ we get log ft=3.9 that coincides with the experimental value. With increasing  $\mathcal{C}_{f}^{\sigma_{f}}$  the total  $\beta$ -transition strength decreases on retention of the relevant sum rule.

We have calculated the matrix elements of the two-neutrino double  $\tilde{b}$ -decay for  $^{128,130}$ Te at  $G_{f}^{0f} = -0.2 \varkappa_{f}^{0f}$  that are in agreement with ref.<sup>/2/</sup> and the experimental value. The necessity of taking into account the interaction in the particle-particle channel to suppress the strength of the two-neutrino double  $\beta$ -decay is confirmed in ref.<sup>/10/</sup>.

In describing ß-decays of nuclei far from stability, the axialvector constant  $G_A$  of a weak interaction is renormalised. In order to obtain an agreement with the experimental ft-values, it has been assumed in refs.<sup>5,6/</sup> that  $|G_A/G_V| = 0.6-0.8$  and in ref.<sup>11/</sup> that  $|G_A/G_V| =$ = 0.7-1.0. Our calculations are performed with  $|G_A/G_V| = 1$ . The ft-values, that are in agreement with the experiment, have been obtained for  $|G_A/G_V| = 1.25$  with increasing  $G_4^{01}$  by 3%. Therefore, one cannot conclude on the renormalisation of the axial-vector constant  $G_A$  of the weak interaction from the calculations of  $\beta^+$ -decays of nuclei far from stability.

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Received by Publishing Department on November 24, 1987.

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Кузьмин В.А., Соловьев В.Г. Е4-87-831 Влияние взаимодействия в канале частица-частица на гамов-теллеровский β<sup>+</sup>-распад в сферических ядрах

В приближении хаотических фаз с учетом остаточных сепарабельных взаимодействий в каналах частица-частица и частица-дырка вычислены ft-величины гамов-теллеровских  $\beta^+$  -распадов ядер <sup>152</sup> Yb, <sup>150</sup> Er, <sup>148,146</sup> Dy и <sup>96</sup> Pd. При фиксированном отношении констант частично-частичного и частично-дырочного взаимодействий получено согласие с экспериментально определенными значениями ft для  $|G_A/G_V| = 1$ . Показано, что из расчетов вероятностей  $\beta$ -распадов ядер, удаленных от полосы стабильности, нельзя сделать вывод о величине перенормировки аксиально-векторной константы  $G_A$  слабого взаимодействия.

Работа выполнена в Лаборатории теоретической физики ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна 1987

Kuzmin V.A., Soloviev V.G. E4-87-831 Effect of the Particle-Particle Interaction on the Gamow-Teller  $\beta^+$ -Decay in Spherical Nuclei

In RPA approximation, taking account of residual separable interactions in the particle-hole and particleparticle channels, ft-values for the Gamow-Teller  $\beta$  -decays of <sup>152</sup> Yb, <sup>150</sup> Er, <sup>148,146</sup> Dy and <sup>96</sup>Pd are calculated. At a fixed constant ratio of the particle-particle and particle-hole interaction agreement with experimentally determined ft-values for  $|G_A/G_V| = 1$  is obtained. It is shown that from the calculations of probabilities of  $\beta$ -decays of nuclei far from stability one cannot conclude on the renormalisation of the axial-vector constant  $G_A$  of the weak interaction.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

Preprint of the Joint Institute for Nuclear Research. Dubna 1987