

259

✓

JOINT INSTITUTE FOR NUCLEAR RESEARCH
Laboratory of Theoretical Physics

P - 259

G. I. Kopylov

THE SIMULATING OF P-P INTERACTION AT 10 BEV

Dubna, 1959.

G. I. Kopylov

THE SIMULATING OF P-P INTERACTION AT 10 BEV

Abstract

Simplified methods of numerical simulating the reactions involving 5-6 secondary particles is described, the isobar model is being taken into account. The table of random stars for p-p interaction at 10 Bev containing 201 stars in the center of mass system and in the laboratory system is given.

I.

A method of constructing a table of random stars is given in [1]. It needs an electronic computer. An attempt was made of constructing such a table using the hand calculation to study the possibilities of this method. We succeeded in doing it for the reaction of pp interaction at 10 Bev basing in the following assumptions: a) the isobar model of the multiple production is considered to be correct; b) the reactions involving more than 6 secondary particles are not taken into account; c) in some cases the mass of a pion is considered equal to zero. Besides the table of 201 random stars the simplified methods of its constructing are given in this paper.

The analysis of data from the table will be made in another paper.

All the notations are taken from [1].

1. Basic formulae

The preliminary calculations have shown that the methods of constructing the table from [1] when there are more than 4 secondary particles becomes not effective for the hand calculation.

The necessity of too frequent rejection of the sets of the momentum components which do not satisfy the inequalities (2.18) leads to many useless calculations. To increase the efficiency of the methods it is necessary to decrease the number of the limitations on the momenta of the particles. So, e.g., if there are five secondary particles, the inequality (2.18) in rejection technique is verified twice (after the sampling of the momenta of the second and the third particles.) Instead of this one may restrict oneself by single rejection after the sampling of the second particle momentum. One should try to choose of the momentum the third particle using the direct method (method B from [1]). A direct method is more difficult than the rejection technique but it does not require the rejection of the values of the momentum components already chosen.

The direct method can be applied to the sampling of the momenta only if one makes use of the assumptions a) and c). Let us pass to the derivation of the corresponding formulas.

We start from the formula for the statistical weight of the system of n particles

$$S_n(E, \vec{P}) = \int d\vec{p}_1, d\vec{p}_2, \dots, d\vec{p}_{n-2} S_2(E_{n-1}, \vec{P}_{n-1}) \quad (1)$$

where (see (2.25))

$$S_2 = \frac{4}{3} \pi p_{n-1}^* \frac{\sqrt{E_{n-1}^* E_n^* E_{n-1}^2 - p_{n-1}^{*2} P_{n-1}^2}}{M_{n-1}^3} \quad (2)$$

Write (1) in the form of a chain of the recursion formulae

$$S_{\nu+1}(E_{n-\nu}, \vec{P}_{n-\nu}) = \int d\vec{p}_{n-\nu} S_\nu(E_{n-\nu+1}, \vec{P}_{n-\nu+1}), \quad (3)$$

where it is assumed that $E_i = E, \vec{P}_i = \vec{P}$, whereas $\nu (\nu=2,3,\dots,n-1)$ are integers. The limits of the integration in (3) may be found quite simply if $m_{n-\nu} = 0$. In this case (3) is convenient to be written as follows:

$$S_{\nu+1}(E_{n-\nu}, \vec{P}_{n-\nu}) = 2\pi \int d\eta_{n-\nu} \Lambda_\nu(E_{n-\nu}, \vec{P}_{n-\nu}, \eta_{n-\nu}), \quad (4)$$

where

$$\Lambda_\nu = \int_0^{p_{n-\nu}^*} p_{n-\nu}^2 S_\nu(E_{n-\nu} - e_{n-\nu}, \vec{P}_{n-\nu} + \vec{p}_{n-\nu}) dp_{n-\nu}, \quad (5)$$

whereas

$$\Pi_\nu = \frac{M_{n-\nu}^2}{2(E_{n-\nu} + P_{n-\nu} \eta_{n-\nu})} \equiv \frac{M_{n-\nu}^2}{2U_{n-\nu}}. \quad (6)$$

Let the distribution over $\vec{p}_1, \dots, \vec{p}_{n-\nu}$ with the required density $S_{\nu+1}$ be fulfilled in some way. Then it follows from (4) that $\Lambda_\nu(E_{n-\nu}, \vec{P}_{n-\nu}, \eta_{n-\nu})$ presents density of the distribution (not normalized) over $\vec{p}_1, \dots, \vec{p}_{n-\nu+1}, \eta_{n-\nu}$. It follows from (5) that the density of the distribution over $\vec{p}_1, \dots, \vec{p}_{n-\nu-1}, \eta_{n-\nu}, p_{n-\nu}$ is $p_{n-\nu}^2 S_\nu(E_{n-\nu+1}, \vec{P}_{n-\nu+1})$.

Further we shall designate through α the random numbers uniformly distributed in (0,1). Making use of the direct method of sampling (see^[1], page 8), for the sampling $\eta_{n-\nu}$ we obtain the equation

$$2\pi \int_0^{p_{n-\nu}^*} d\eta_{n-\nu} \Lambda_\nu = 2\pi \alpha \int_0^1 d\eta_{n-\nu} \Lambda_\nu \equiv \alpha S_{\nu+1}, \quad (7)$$

whereas for the sampling $p_{n-\nu}$ - the equation

$$\int_0^{p_{n-\nu}^*} dp_{n-\nu} \cdot p_{n-\nu}^2 S_\nu = \alpha \int_0^{p_{n-\nu}^*} dp_{n-\nu} \cdot p_{n-\nu}^2 S_\nu \equiv \alpha \Lambda_\nu \quad (8)$$

Introduce the following notations:

$$E_{n-\nu} + P_{n-\nu} \eta_{n-\nu} = U_{n-\nu}, \quad (9)$$

$$2\pi \int \Lambda_\nu d\eta_{n-\nu} = \bar{\Lambda}_\nu(\eta_{n-\nu}), \quad (10)$$

$$\int d\rho_{n-\nu} \cdot \rho_{n-\nu}^2 S_\nu = \bar{S}_\nu(\rho_{n-\nu}, \eta_{n-\nu}) \quad (11)$$

(in (10) and (11) are omitted $E_{n-\nu}$ and $P_{n-\nu}$ among the arguments Λ_ν and S_ν). Then the equations (7) and (8) may be finally rewritten in the form:

$$\bar{\Lambda}_\nu(\eta_{n-\nu}) = \alpha \bar{\Lambda}_\nu(1) + (1-\alpha) \bar{\Lambda}_\nu(-1), \quad (7')$$

$$\bar{S}_\nu(\rho_{n-\nu}, U_{n-\nu}) = \alpha \bar{S}_\nu(\Pi_\nu, U_{n-\nu}). \quad (8')$$

To make use of the formulas (7') and (8') for the sampling $\eta_{n-\nu}$ and $\rho_{n-\nu}$ it is necessary to know the concrete form of the functions $\bar{\Lambda}_\nu$ and \bar{S}_ν . We shall present the corresponding expressions for $\nu=2$ and $\nu=3$ under the assumption $m_{n-1} = m_n = 0$. Then the initial formula - the expression for S_2 - is of the form:

$$S_2 = \frac{\pi}{6} (3E_{n-1}^2 - P_{n-1}^2). \quad (12)$$

$\nu=2$. Assume $m_{n-2} = 0$. From (11) and (12) we obtain

$$\bar{S}_2(\rho, U) = \frac{\pi}{6} \left(\frac{2E^2 + M^2}{3} - \frac{2E+U}{2} \rho + \frac{2}{5} \rho^2 \right) \rho^3 \quad (13)$$

(the index $n-2$ is meant everywhere in the magnitudes without it).

Then it follows from (5)

$$\bar{\Lambda}_2(\eta) \equiv \bar{S}_2(\Pi_2, U) = \frac{\pi}{6} \left(\frac{8E^2 + M^2}{12} - E\Pi_2 + \frac{2}{5} \Pi_2^2 \right) \Pi_2^3 \quad (14)$$

while from (10)

$$\bar{\Lambda}_2(\eta) = \frac{\pi^2 M^2}{5} \left(\frac{8E^2 + M^2}{24} - \frac{E\Pi_2}{3} + \frac{\Pi_2^2}{10} \right) \Pi_2^2 \quad (15)$$

Finally, from (4)

$$S_3(E, P) \equiv \bar{\Lambda}_2(1) - \bar{\Lambda}_2(-1) = \frac{\pi^2}{240} EM^2(7E^2 - 3P^2). \quad (16)$$

$\nu=3$. In the same order as for $\nu=2$, can be obtained¹⁾:

$$\begin{aligned} \bar{S}_3(P, u) = & \frac{\pi^2}{240} \left\{ \frac{EM^2(4E^2 + 3M^2)}{3} - \frac{3M^2(M^2 + 4E^2) + 4uE(3M^2 + 2E^2)}{4} \rho + \right. \\ & \left. + \frac{12}{5} [EM^2 + (2E^2 + M^2)u + Eu^2] \rho^2 - \frac{2}{3} (M^2 + 6Eu + 3u^2) \rho^3 + \frac{8}{7} u \rho^4 \right\} \rho^3; \\ \Lambda_3(\eta) = & \frac{\pi^2}{240} \cdot \frac{M^2}{2} \left[\frac{(10E^2 + 3M^2)E}{15} - \frac{12E^2 + M^2}{10} \Pi_3 + \frac{4}{5} E \Pi_3^2 - \frac{4}{21} \Pi_3^3 \right] \Pi_3^3; \end{aligned} \quad (18)$$

$$\bar{\Lambda}_3(\eta) = \frac{\pi^3}{120} \frac{M^4}{20P} \left(\frac{10E^2 + 3M^2}{6} E - \frac{12E^2 + M^2}{6} \Pi_3 + E \Pi_3^2 - \frac{4}{21} \Pi_3^3 \right) \Pi_3^2; \quad (19)$$

$$S_4(E, \bar{P}) = \frac{\pi^3}{2 \cdot 8!} M^4 (33E^4 - 18E^2 P^2 + P^4). \quad (20)$$

One may get the formulas for an arbitrary ν using the same method as used in [2] for $S_\nu(E_{n-\nu+1}, 0)$. The obtained expressions are, however, too big, we do not intend to present them here. We write down only the formula for $S_\nu(E_{n-\nu+1}, \bar{P}_{n-\nu+1})$ (the index is omitted);

$$\begin{aligned} S_\nu(E, \bar{P}) = & - \frac{(2\pi)^{\nu-1}}{(2\nu-1)!} \frac{1}{2} \frac{d}{dP^2} \sum_{m=0}^{2\nu-1} \frac{(2\nu+m-1)!}{(2\nu-m-1)!(\nu+m-1)!(2m)!!} \times \\ & \times [(E-P)^{\nu+m-1} - (-1)^{2\nu-m-1} (E+P)^{\nu+m-1}] P^{2\nu-m-1} \end{aligned} \quad (21)$$

Further we shall need the magnitudes of the maximum of the function

$$\Phi_\nu = S_\nu \prod_{i=1}^{n-\nu} \rho_i^2 \quad (22)$$

Since Φ_ν depends upon $\rho_2, \dots, \rho_{n-\nu}$ only through $P_{n-\nu+1}^2$, then one of the equations determining the position of the maximum, is of the form:

$$P_{n-\nu+1}^2 = 0 \quad (23)$$

¹⁾The index $n-3$ is everywhere meant in the magnitudes without an index.

It follows from here, by the way, that there exists an infinite set of the values ν_i , which maximized Φ_ν .

On the other hand, Φ_ν depends upon $\rho_1, \dots, \rho_{n-\nu}$ both through $E_{n-\nu+1}, \bar{P}_{n-\nu+1}$ and directly. This leads to the following system of the equations

$$\frac{\rho_1^2}{e_1} = \frac{\rho_2^2}{e_2} = \dots = \frac{\rho_{n-\nu}^2}{e_{n-\nu}} = \frac{2}{3\nu-4} E_{n-\nu+1} \quad (24)$$

It is convenient to solve it using the iteration method. If we assume as a first approximation, for instance, its solution when $m_1 = \dots = m_{n-\nu} = \frac{1}{n-\nu} \sum_{i=1}^{n-\nu} m_i$ i.e. the positive root $e = e^{(0)}$ of the quadratic equation

$$e^2 \left[1 + \frac{2(n-\nu)}{3\nu-4} \right] - \frac{2}{3\nu-4} E e - \frac{1}{n-\nu} \sum_{i=1}^{n-\nu} m_i = 0, \quad (25)$$

it is necessary to calculate $\tilde{\chi}_i = \frac{\rho_i^2}{e_i} - \frac{2}{3\nu-4} E_{n-\nu+1}$

at $e_i = e^{(0)}$. Then one may determine the correction for the previous approximation from the linear system

$$\sum_{j=1}^{n-\nu} \left[\frac{2}{3\nu-4} + \left(1 + \frac{m_i^2}{e_i^2} \right) \delta_{ij} \right] de_j = -\tilde{\chi}_i^{(0)}. \quad (26)$$

Let us designate the roots (24) through \bar{e}_i .

The magnitude of the maximum Φ_ν is as follows:

$$\Phi_{\max} = \frac{\pi^{\nu-1}}{2^{\nu-1}} (2\pi)^{n-\nu} \left(\frac{2}{3\nu-4} \right)^{n-\nu} \frac{(4\nu-4)! (2\nu-1)}{[(2\nu-1)!]^2 (3\nu-4)!} \bar{E}_{n-\nu+1} \cdot \prod_{l=1}^{n-\nu} \bar{e}_l \quad (27)$$

(formula (3.16) is used from [2]; $\bar{E}_{n-\nu+1} \equiv E - \sum_{i=1}^{n-\nu} \bar{e}_i$).

Note that at great ν the surface Φ_ν has a sharp peak in the point of the maximum. This restricts the efficiency of the proposed method.

2. Methods of Sampling

Formulas (13)-(16) make it possible to get 5-particle stars, whereas formulas (13)-(20) - 6-particle ones. Let us describe the procedure of the sampling of 6-particle stars in more detail.

1. ρ_1 and ρ_2 respectively in $(0, \tilde{\rho}_{1, \max}), (0, \tilde{\rho}_{2, \max})$ are picked uniformly. So is η_2 in $(-1, +1)$. Only such values of the numbers ρ_1, ρ_2, η_2 are left for which $M_3 \geq \mu_2$.

2. Φ_4 is calculated by (20), (22) and compared with a random number d from the interval $(0, \Phi_{4, \max})$. Only such sets of ρ_1, ρ_2, η_2 for which $\Phi_4 \geq d$, are left.

The further procedure of the sampling is carried out without a rejection. Therefore,

the calculations according to the points 1 and 2 are repeated until a necessary number of sets is obtained.

3. The equation (7') where (19) is substituted, is being solved with respect to Π_3 . From (6) one may find u_3 and η_3 .

4. The system (8') (17) with respect to P_3 is being solved.

After the calculations according to the points 3-4 the momentum of the third particle \vec{p}_3 is wholly defined. Therefore, it is possible to determine E_4, P_4, M_4 . It is not necessary to check the inequality $M_4 \geq \mu_3$, it must be fulfilled automatically. (Note that now $\mu_3 = \mu_4 = \mu_5 = \mu_6 = 0$).

Then one passes to the sampling of the components \vec{p}_4 .

5. (7')-(15) are being solved with respect to $\Pi_2; \eta_4$ and u_4 are calculated.

6. (8')-(13) are being solved with respect to $\rho_4; E_5$ and P_5 are calculated.

Thus, the components \vec{p}_4 are determined. The sampling of the module of the momentum P_5 is made as is shown in [1], i.e.

7. The equation (2.27) is being solved, from where e_5 is found, then ρ_5 . The remaining components of the 5 and 6 particles may be found from the laws energy and momentum conservation

$$\eta_5 = \frac{E_6^2 - m_6^2 - P_5^2 - P_5^2}{2P_5\rho_5} = \frac{M_5^2 + m_5^2 - m_6^2 - 2E_5e_5}{2P_5\rho_5} \quad (28)$$

$$\vec{p}_6 = -\sum_{i=1}^5 \vec{p}_i \quad (29)$$

When the stars with $n = 5$ are sampled the points 3 - 4 are omitted, while in the points 5-7 the indeces in all magnitudes are decreased by a unit.

When the stars with $n = 4$ are sampled after point 2 immediately follows 7, with the decrease of all the indeces by two.

The sampling of the stars with $n = 3$ is given in [1].

Before enumerating further operations with the obtained sets $p_1, \dots, p_{n-1}, \eta_2, \dots, \eta_{n-1}$ the following should be noted.

The inequality $M_4 \geq \mu_3$ is fulfilled practically in a great majority of cases after the calculations by 4. It is so even if in the calculation of E_4 we assume $m_3 \neq 0$ and in the calculation of μ_3 assume $m_4, m_5, m_6 \neq 0$. The same holds for the fulfillment of other similar inequalities. Therefore, though the formulae of the previous paragraph were derived under the assumption that the mass of a pion is equal to zero the mass of a pion was not practically neglected in the calculations. Only in several (five-six) events this has led to imaginary values of M_4 or to the violation of the inequalities of $M_4 \geq \mu_3$.

type. In these cases the calculation was made with $m_{n-\nu} = 0$ for the second time. The use at $m_{n-\nu} \neq 0$ of the formulas correct at $m_{n-\nu} = 0$ would, evidently, distort in some way the characteristics of the low energy pions containing in the table. This distortion, however, must be small. At the same time the condition $m_{n-\nu} \neq 0$ made it possible to avoid the difference in the characteristics of the mesons produced in the reactions with $n = 5-6$ as well as in the reactions with $n = 3-4$ (in the latter ones the condition $m_{n-\nu} = 0$ is not obligatory).

Note also, that the indeces 1 and 2 should be given to heavy secondary particles, while other indeces are meant for the light ones. The reactions involving more than two secondary baryons cannot be chosen using this method by means of the hand calculation.

The efficiency of this method of the sampling is characterized by Table 1. Roughly speaking, this method makes it possible to keep the efficiency at a level of 10% with the increase of a number of particles.

Enumerate further operations with the quantities $p_1, \dots, p_{n-1}, \eta_2, \dots, \eta_{n-1}$.

8. η_1 in $(-1, 1)$ and $\varphi_1, \dots, \varphi_{n-1} \ln(0, 2\pi)$ are picked. The components of the momenta in the rectangular coordinate system are calculated by the formulas (2.9)-(2.12). At this to check the calculations the formulas

$$\left. \begin{aligned} x_\kappa^2 + y_\kappa^2 + z_\kappa^2 &= p_\kappa^2, \\ X_\kappa^2 + Y_\kappa^2 + Z_\kappa^2 &= P_\kappa^2, \\ X_{\kappa+1} &= X_\kappa + x_\kappa, \quad Y_{\kappa+1} = Y_\kappa + y_\kappa, \quad Z_{\kappa+1} = Z_\kappa + z_\kappa \end{aligned} \right\} \quad (30)$$

are used.

9. The rectangular components \bar{p}_n are calculated by (29). As a control calculation the verification of the fulfillment of the law of energy conservation is used.

10. If an isobar is produced in the reaction the sampling of the isobar decay should be also made. When making the table it was assumed that: 1) the mass of the isobars is constant and equal to 1.24 BeV/c²; 2) the products of the isobar decay are distributed in the space isotropically. Both these assumptions may be easily replaced by other ones more real.

Under the assumption I the characteristics of the decay in the system of reference where the isobar rests are the constants. Having chosen uniformly in $(-1, 1)$ the cosine η^* of the emergence angle of a nucleon we obtain the longitudinal and transverse nucleon momentum by the formulas

$$\left. \begin{aligned} p_{N'} &= p^* \sqrt{1 - \eta^{*2}} \cong 0,233 \sqrt{1 - \eta^{*2}} \\ p_{N \rightarrow} &= \frac{e_{N'} p^* \eta^* + p_N e_N^*}{m_{N'}} \cong 0,188 e_{N'} \eta^* + 0,781 p_{N'} \end{aligned} \right\} \quad (31)$$

(the quantities with the index N' are referred to an isobar).

It is simply to get from here in the L-system the emergence angle of a nucleon with respect to the initial direction of an isobar (characterized by the vector $\{x_{N'}, y_{N'}, z_{N'}\}$).

Having chosen uniformly in $(0, 2\pi)$ the azimuth φ^* of the direction of the nucleon emergence one may make use of the formulas (2.9)-(2.12) for obtaining the components of the nucleon momentum in the L-system. The substitution leads to the following sequence of the formulas convenient for the calculation:

$$P_N^2 = P_{N\rightarrow}^2 + P_{N\uparrow}^2; \quad R_{N'}^2 = x_{N'}^2 + y_{N'}^2; \quad (32_{1-2})$$

$$u = \frac{x_{N'}}{P_{N'}}; \quad v = \frac{y_{N'}}{P_{N'}}; \quad W = \frac{z_{N'}}{P_{N'}} \cos \varphi^*; \quad \xi = \frac{x_{N'}}{R_{N'}}; \quad \zeta = \frac{y_{N'}}{R_{N'}}; \quad (33_{1-5})$$

$$\left. \begin{aligned} A_N &= W\xi - \zeta \sin \varphi^*; \\ B_N &= W\zeta + \xi \sin \varphi^*; \\ C_N &= (R_{N'}/P_{N'}) \cos \varphi^*; \end{aligned} \right\} \quad (34_{1-3})$$

Control:

$$A_N^2 + B_N^2 + C_N^2 = 1; \quad (35)$$

$$x_N = u P_{N\rightarrow} + A_N P_{N\uparrow};$$

$$y_N = v P_{N\rightarrow} + B_N P_{N\uparrow};$$

$$z_N = W P_{N\rightarrow} - C_N P_{N\uparrow};$$

$$x_N^2 + y_N^2 + z_N^2 = P_N^2$$

(36₁₋₃)

Control:

It is convenient to calculate the components of the pion momentum by formulae

$$\vec{p}_\pi = \vec{p}_{N'} - \vec{p}_N \quad (37)$$

Control:

$$e_N + e_\pi = e_{N'} \quad (38)$$

11. Finally, the momenta of the particles are transformed in L-system.

3. Distribution by Charge States

Simplified methods for sampling the stars for different reactions is described in the previous paragraphs. Now there remains to enumerate those reactions which were taken into account in the course of the tabulation and to give the used distributions by the charge states.

At 10 Bev in p-p interactions both multiple production of pions and the production of antiparticles and strange particles may occur. The latter process on the photographs may be separated from the other ones. Therefore only the production processes of pions and

antinucleons are included into the table. The statistical weights of these processes at $T_p^{kin} = 10$ Bev are given in Table 2. The unpublished calculations which are the basis of paper [3]* were used there.

When making the table of random stars we restricted ourselves by the first 14 reactions from Table 2, i.e. by the reactions $\rho\rho \rightarrow NNn\pi, \bar{N}\bar{N}n\pi, N'\bar{N}'n\pi$ ($n=1-4$) and reactions $\rho\rho \rightarrow 3N\bar{N}(\bar{N}')$. The rejected reactions of elastic scattering, of the production of more than 4 mesons and some other are less than 4% of the weight of the reactions taken into account.

Starting from the general length of the table consisting of 200 lines the number of lines for each reaction was taken proportional to its statistical weight (Table 2).

One could not resort to a more correct procedure of the random sampling of the number of lines (see § 3 from [1]). At this the necessity of including the reactions with 7 secondary particles might appear but their sampling cannot be realized by the hand calculation.

Thus, the number of lines for each reaction does not fluctuate with respect to the value expected by the isobar model.

Each of the reactions included into the table may be concretely realized by different charge states. So, for instance, the reaction $\rho\rho \rightarrow NN\pi$ may have 8 realizations:

$\rho^+\rho^-\pi^-, \rho^+n\pi^0, \rho^0\rho^-\pi^0, \rho^0n\pi^+, \rho^-\rho^-\pi^+, n^+n\pi^+, n^+\rho^-\pi^0, n^0\rho^-\pi^+$
(ρ^+ , for instance, designates an isobar disintegrated into ρ and π^+). The number of lines for each charged state has been already randomly chosen. The statistical weights of the states were calculated starting from the assumption about the equal probability of each isotopic state [4,5]. The charge distributions from [5] are not applicable for the table of random stars since there the isobar decay has been already taken into account. Therefore, the calculation of the charge distribution had to be carried out a new for all the reactions where the isobars appear. The methods used in the calculations in general coincide with those stated in [5].

They are as follows.

Let the formula be obtained which presents the probability that the system with the isotopic spin T and its Z -component t is composed of two subsystems: the subsystem of two nucleons (isobars) in the state (T_1, t_1) and the subsystem of the n -mesons with a definite sign.

We shall write it in the form

$$(T, t) = \sum_i C_i^{(n)}(T_i, t_i; n\pi). \quad (39)$$

Here $C_i^{(n)}$ are the probabilities (unnormalized). To obtain an analogous formula with $(n+1)$ mesons from here it is sufficient to have some formulas of the form

* The author expresses his gratitude to V.S. Barashenkov and V.M. Maltsev for the data which they kindly gave him.

$$(T, t) = \sum_{\kappa} C_{\kappa}^{(n)} (T_{\kappa}, t_{\kappa}; \mathcal{I}) \quad (40)$$

for all $T = T_i$ and $t = t_i$. Substituting (40) into (39) it is possible to obtain the formula we were looking for by multiplying and summing $C^{(n)}$ and $C^{(l)}$

$$(T, t) = \sum_i C_i^{(n)} \sum_{\kappa} C_{\kappa}^{(l)} (T_{\kappa}, t_{\kappa}; (n+1)\pi) \quad (39')$$

If in addition one has a set of formulas which present the probability of different charge states of the subsystem NN or N'N or N'N' in the state (T_{κ}, t_{κ}) then from the formulas of (39') type we obtain the probabilities of different charge states of the system involving 2 nucleons (isobars) and $(n+1)$ mesons.

The coefficients from formula (40) are given in table 3. Those from (39) obtained in the transition from n to $n+1$ are given (for n from 1 up to 4) in table 4. Table 5 contains the probabilities that the subsystems (T, t) are in different charge states. The substitution of the coefficients from Table 5 into Table 4 gives respective (unnormalized) probabilities of different charged states. They are presented in Table 6. It was used for the sampling of the distribution by the charge states inside each reaction (according to the methods used in § 3|1|)

In accordance with the results of the sampling in each random star the data concerning neutral particles were crossed out.

4. Table of Random Stars

The data concerning only charged particles are included into the table presented in this paper. The stars are shuffled at random. The following data on stars are given in the table: the observed particles (p-proton, $\pm - \pi^{\pm}$), the momentum components in the center of mass system (p_x^*, p_y^*, p_z^*) , the momentum p^* in the center-of-mass system, Z-component of the momentum p_z and the momentum P in the L-system. The center-of-mass system is considered to move along the axis Z of the L-system. It was assumed in the calculations that the incident protons have the kinetic energy $T = 10$ Bev, the mass of a nucleon is taken equal to $0,94 \text{ Bev}/c^2$, the mass of a meson $0,14 \text{ Bev}/c^2$. The figures are given in Mev/c . The accuracy up to $1 \text{ Mev}/c$ for the magnitudes less than 1000, and $10 \text{ Mev}/c$ for the magnitudes more than 1000 (in the calculations the accuracy was 4-5 signs). The random numbers were taken from [6].

The calculations were made in the calculation bureau of the Joint Institute for Nuclear Research, by A.V. Modestova, N.V. Demina, L.A. Shustrova, L.A. Isajeva. The author expresses his gratitude to them.

References

1. G.I. Kopylov GETP, 35, 6 (12), 1426, (1958); preprint of the Joint Institute for Nuclear Research, P-205 (1958);
Yu.N. Blagoveschensky, G.I. Kopylov, preprint of the Joint Institute for Nuclear Research, P-213 (1958).
2. R.H. Milburn, Rev. Mod. Phys., 27, 1 (1955).
3. V.S. Barashenkov, B.M. Barbashov, E.G. Bubelev "Statistical theory of multiple production" preprint of the Joint Institute for Nuclear Research;
V.S. Barashenkov, V.M. Maltsev, preprint of the Joint Institute for Nuclear Research, P - 98 (1957).
4. E. Fermi, Phys.Rev. 92, 452 (1933); 93, 1434 (1954).
5. S.Z. Belenky, A.I. Nikishev, V.M. Maksimenko, L.L. Rozentel, Uspekhi Fiz.Nauk., 62, N 2, 1 (1957).
6. "A Million of Random Digits", Wiley Publ. in Statistics.

T A B L E OF RANDOM STARS

The number Particle	p_x^*	p_y^*	p_z^*	p^*	p_z	p	The number Particle	p_x^*	p_y^*	p_z^*	p^*	p_z	p
1. p	+ 54	- 41	+ 397	403	3360	3360	12. p	- 410	+ 224	- 417	626	1560	1630
+	+ 655	+ 271	- 404	816	892	1140	+	- 150	+ 11	+ 309	344	1630	1640
2. p	- 298	- 19	+ 132	326	2630	2640	-	- 230	- 343	- 319	522	446	607
-	+ 98	- 12	- 200	223	104	143	+	- 223	- 249	- 540	630	130	359
-	+ 89	- 153	- 199	266	194	263	13. p	- 418	- 907	+ 386	1071	4260	4370
+	- 245	- 349	+ 511	665	2850	2880	+	- 50	- 64	- 168	186	116	142
+	- 455	+ 199	+ 66	501	1370	1450	+	+ 7	- 248	- 170	301	339	420
+	+ 881	+ 280	- 400	1007	1340	1630	-	- 105	+ 439	+ 320	553	2120	2170
3. p	- 234	- 177	+ 48	298	2400	2410	-	+ 161	+ 41	+ 227	281	1300	1310
-	- 77	+ 73	- 78	132	248	270	+	+ 423	+ 231	+ 176	513	1670	1740
+	+ 45	+ 98	- 4	108	239	262	14. p	- 493	- 527	- 246	762	2170	2290
+	- 90	+ 165	- 10	188	408	449	-	- 270	+ 95	+ 44	290	854	901
4. p	- 653	-1143	+ 206	1332	4280	4480	+	+ 543	- 73	- 14	548	1270	1380
+	- 135	- 50	- 170	223	179	230	+	+ 171	- 4	- 3	171	387	423
5. p	+ 136	+ 263	- 191	352	1840	1860	15. p	+ 348	- 274	+ 451	632	3750	3770
+	- 40	- 84	- 973	977	- 205	225	p	- 291	+ 45	+ 634	699	4290	4300
6. p	-1367	- 342	- 33	1410	3820	4080	+	+ 111	+ 75	- 142	195	197	238
+	- 60	+ 84	- 138	172	166	195	-	- 509	- 414	- 138	670	1230	1400
-	+ 253	+ 562	- 5	617	1450	1580	16. p	-1196	+ 296	+ 594	1367	5320	5460
+	+ 325	+ 113	+ 165	382	1350	1390	+	- 95	+ 134	+ 345	382	1810	1810
7. +	+ 111	+ 88	- 248	286	112	181	-	+ 130	- 367	- 152	418	635	745
+	- 266	+ 272	- 546	666	198	429	+	+ 196	+ 110	- 310	383	161	277
8. p	+ 530	- 454	+ 191	724	3220	3290	17. p	- 823	- 438	- 770	1209	1600	1850
p	+ 668	+ 507	+ 140	850	3280	3380	-	+ 347	- 61	+ 168	390	1380	1420
-	- 153	- 169	+ 28	230	691	728	+	+ 326	+ 3	+ 67	333	1000	1060
+	- 467	- 304	+ 578	803	3330	3380	+	- 726	- 298	- 045	786	1730	1900
9. p	-1060	+ 531	- 778	1418	1970	2300	18. +	+ 129	+ 783	+ 431	903	3190	3290
+	+ 247	- 357	+ 365	567	2260	2310	+	- 318	+ 139	+ 508	615	2740	2760
-	+ 707	+ 309	+ 117	780	2120	2260	19. p	- 899	- 855	+1627	2046	9290	9370
+	+ 2	- 77	- 198	212	90	118	+	+ 147	+ 76	- 170	237	206	264
10. p	- 306	+ 117	- 110	346	2030	2060	20. p	- 652	+1133	+1184	1764	7590	7700
+	- 54	- 18	- 336	341	4	57	+	+ 66	- 494	-1092	1201	43	501
-	+ 51	- 186	+ 797	820	3920	3930	21. p	- 710	+ 151	- 602	944	1560	1720
+	+ 336	- 327	+ 179	502	1650	1720	p	- 191	- 426	+ 478	668	3860	3890
11. p	- 122	+ 241	- 12	270	2230	2240	+	+ 79	+ 180	- 155	250	272	336
+	- 138	- 25	- 275	309	90	167	+	- 205	- 238	+ 494	585	2630	2650
-	- 688	- 441	+ 15	816	1950	2110	-	+ 714	- 20	+ 142	728	2070	2190
+	+ 220	+ 348	+ 702	814	3670	3700	-	+ 313	+ 353	- 357	592	505	691

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P
22. p	+ 30	+ 694	- 170	715	2300	2400	33. p	- 170	- 370	+ 10	407	2390	2420
+	+ 76	+ 489	+ 152	517	1620	1690	+	+ 71	+ 79	+ 917	923	4460	4460
+	- 42	+ 266	- 62	276	560	620	34. p	+ 81	+ 736	+ 197	766	3290	3380
-	+ 510	-1493	- 69	1579	3490	3830	p	+ 55	+ 129	- 730	744	929	940
23. p	+ 63	- 144	+ 337	372	3180	3190	35. +	- 354	+ 174	- 38	396	876	951
+	+ 298	- 77	- 17	308	737	799	+	- 389	- 372	- 632	829	351	643
24. p	+ 262	+ 490	+ 19	556	2570	2630	36. +	+ 272	- 162	+ 46	320	923	976
p	+ 698	+ 109	+ 562	902	4420	4470	+	+1047	+ 752	- 550	1402	1870	2270
-	- 406	- 116	+ 14	423	1060	1140	37. p	- 129	- 609	+ 247	670	3280	3340
+	- 258	- 213	+ 90	346	1090	1140	+	+ 206	+ 12	+ 18	207	622	655
25. p	- 149	- 187	+ 265	357	2980	2990	-	+ 559	+ 15	- 345	657	683	883
-	+ 159	- 90	- 94	205	336	382	+	-1296	+ 262	+ 446	1395	4360	4550
+	+ 491	- 309	+ 554	802	3270	3320	38. p	+ 315	- 496	- 95	596	2330	2400
+	- 414	+ 156	+ 112	457	1380	1450	p	- 339	- 113	+ 159	391	2750	2770
26. p	- 39	+ 399	+ 471	618	3780	3800	+	+ 949	- 64	- 4	952	2210	2400
p	+ 607	+ 235	- 276	707	2020	2120	-	- 788	+ 965	- 166	1257	2500	2790
+	- 102	+ 188	+ 500	544	2550	2560	39. p	-1284	- 385	- 506	1433	2680	2980
-	- 474	- 172	- 508	716	406	647	p	+ 954	+ 450	- 235	1081	4140	4280
27. p	+ 434	- 706	- 136	840	2570	2700	+	- 64	+ 43	+ 107	132	714	718
+	+ 194	+ 82	+ 77	224	803	830	-	+ 394	- 108	+ 633	754	3360	3390
+	- 430	- 112	+ 28	446	1150	1230	40. p	+ 2	+ 123	+ 553	566	3920	3920
-	+ 316	+ 139	+ 208	402	1500	1540	p	- 832	+ 861	- 457	1281	2520	2790
28. +	+ 137	+ 23	+ 56	150	614	630	+	- 66	- 255	+ 224	346	1430	1450
-	+ 139	- 316	+ 45	348	978	1037	-	+ 279	- 125	+ 232	384	1530	1560
+	+ 366	+ 63	- 358	516	334	499	-	+ 309	- 282	- 645	769	181	456
+	+ 380	- 229	+ 416	608	2490	2530	+	+ 307	- 321	+ 93	454	1330	1400
29. p	- 204	- 884	- 69	910	2840	2990	41. p	+ 502	+1120	+ 402	1292	-4700	4860
+	- 104	- 245	- 323	419	206	337	p	- 755	-1140	- 516	1464	2720	3040
30. p	+ 709	- 883	- 622	1292	2120	2400	+	+ 201	+ 218	- 174	344	420	514
p	- 186	+ 38	- 251	315	1660	1670	+	+ 268	- 406	+ 211	530	1800	1860
+	+ 426	+ 173	+ 224	512	1790	1850	-	-109	+ 15	+ 100	148	721	729
-	- 378	+ 579	+ 573	898	3540	3600	-	- 107	+ 196	- 24	224	551	595
31. p	+ 6	- 428	- 58	432	2240	2280	42. p	+ 77	+ 106	- 443	462	1300	1310
-	+ 7	- 400	+ 184	440	1530	1580	-	+ 162	+ 342	+ 128	400	1300	1350
+	+ 515	- 11	- 246	571	731	894	+	- 186	- 35	+ 84	208	790	812
+	- 247	+ 459	+ 391	652	2520	2570	+	- 158	- 301	+ 249	421	1650	1680
32. p	-1180	+ 520	- 223	1308	3160	3410	-	- 37	- 220	+ 614	652	3080	3090
-	+ 461	- 448	- 160	663	1160	1330	+	+ 17	- 144	+ 148	207	948	959
+	- 111	- 83	+ 141	198	916	926							
+	+ 537	+ 228	+ 150	602	1800	1900							

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P
							53. p	+ 122	+ 157	- 595	627	1110	1130
							+	- 242	- 873	- 448	1010	1230	1530
43. p	- 18	+ 18	- 572	572	1100	1100	54. p	- 406	-1413	- 870	1708	2310	2740
p	+ 351	+ 648	+ 689	1009	4910	4970	-	+ 144	+ 64	- 125	201	250	295
-	+ 265	- 101	- 316	424	236	369	+	+ 53	- 74	+ 62	111	568	575
+	+ 145	+ 184	+ 96	253	911	942	+	- 300	+ 143	+ 521	618	2770	2790
44. p	- 794	+ 1	+ 79	798	3040	3140	55. p	- 73	+ 616	+1238	1385	6980	7000
p	+ 186	+ 780	- 138	814	2520	2640	+	- 854	- 776	- 649	1324	1440	1840
-	+ 14	- 89	- 170	192	122	152	56. p	+ 495	+ 492	-1007	1225	1030	1240
+	+ 194	+ 90	- 554	593	15	214	+	- 604	+ 171	+ 859	1063	4640	4680
45. p	- 382	+ 8	+ 74	390	2530	2560	57. p	- 76	-1031	+ 466	1134	4570	4680
+	- 84	- 241	+ 138	293	1090	1120	-	- 147	+ 274	- 53	316	664	733
+	- 574	+ 552	+ 783	1117	4560	4630	+	+ 14	+ 414	+ 183	453	1560	1610
-	+ 632	- 168	- 716	970	460	799	+	+ 282	+ 225	- 702	789	85	370
46. +	- 126	+ 86	- 17	154	437	463	58. p	+ 422	+ 285	+ 274	578	3240	3280
+	- 69	- 831	- 214	861	1470	1690	+	- 799	- 588	+ 899	1339	5360	5460
47. p	+ 404	+ 372	+1017	1156	5990	6020	59. p	- 326	-1342	+ 384	1433	4920	5110
p	+ 10	- 433	- 59	437	2240	2280	p	+ 117	+ 559	- 481	746	1560	1660
+	+ 191	+ 95	- 100	236	381	436	+	+ 337	+ 481	+ 186	616	1930	2010
-	- 337	- 174	- 635	740	141	405	-	- 180	+ 693	+ 66	719	1860	1990
48. p	- 445	- 569	+ 141	736	3110	3190	60. p	+ 37	- 432	+ 145	457	2780	2810
+	+ 341	+ 308	- 521	695	326	564	+	- 201	- 30	+ 236	312	1380	1400
49. p	+ 739	+ 262	+ 20	784	2880	2980	61. p	- 622	+ 320	+ 358	786	3730	3790
p	- 51	+ 28	+ 101	117	2440	2440	+	+ 12	+ 177	- 131	221	274	327
+	+ 228	+ 75	+ 312	394	1750	1770	62. p	- 52	+ 192	+ 303	371	3090	3100
+	- 574	- 344	+ 213	703	2190	2290	p	+ 45	- 18	- 629	631	1030	1030
-	- 433	+ 28	- 560	708	258	505	-	+ 151	- 172	+ 48	246	774	807
-	+ 250	+ 210	- 72	334	656	733	+	+ 594	+ 812	- 19	1006	2300	2510
50. p	- 219	+ 372	- 392	582	1570	1620	-	- 129	- 455	+ 266	542	1960	2020
p	+ 483	- 50	- 405	632	1600	1670	+	- 610	- 359	+ 31	708	1750	1890
+	+ 116	+ 117	+ 136	214	932	946	63. p	-1063	+ 209	- 472	1181	2300	2540
-	- 330	- 424	- 682	868	316	623	+	+ 31	- 304	+ 479	568	2550	2570
51. +	- 182	+ 62	- 114	224	322	375	-	+ 8	+ 41	+ 118	125	734	735
+	+ 16	- 547	- 171	573	931	1080	+	+ 485	+ 242	+ 432	693	2720	2770
-	+ 224	+ 243	+ 262	422	1680	1720	64. p	- 290	- 563	- 62	637	2460	2540
+	- 424	+ 895	+ 372	1058	3400	3540	-	+ 365	- 249	+ 252	509	1850	1900
52. p	+ 155	+ 327	- 469	592	1380	1430	+	- 125	+ 152	+ 73	210	768	793
+	+ 49	+ 192	- 509	547	22	199	+	- 168	+ 52	- 93	199	329	373
+	- 331	- 163	+ 255	449	1720	1760	-	+ 289	+ 363	- 558	726	303	554
-	+ 613	- 383	+1208	1408	6300	6340	+	- 548	+ 235	- 164	618	1050	1210

	* P _x	* P _y	* P _z	* P	P _z	P		* P _x	* P _y	* P _z	* P	P _z	P
65. p	- 679	- 508	- 71	857	2750	2880	^{76.} p	+1053	- 145	- 121	1070	2980	3170
+	+ 49	- 10	- 187	193	81	95	p	- 515	+ 648	- 419	928	1990	2160
-	+ 318	- 356	+ 106	488	1440	1520	-	- 265	- 309	+1174	1242	5830	5850
+	- 79	- 220	+ 450	507	2340	2360	+	- 273	- 194	- 634	717	90	346
66. p	+ 615	- 616	- 213	897	2460	2610	^{77.} p	- 524	+ 889	+ 17	1032	3260	3420
p	- 542	- 464	- 749	1035	1340	1520	+	- 353	- 802	- 748	1152	798	1180
+	+ 350	+ 187	+ 373	545	2240	2270	^{78.} p	+ 578	- 658	+ 307	928	3820	3920
-	- 151	+ 347	- 33	380	852	932	+	- 188	- 474	+ 683	852	3710	3740
67. p	+ 359	+ 462	- 72	590	2380	2450	^{79.} p	+ 379	- 927	- 354	1062	2380	2580
+	- 383	+ 472	- 139	624	1130	1280	p	- 403	+ 181	+ 75	449	2590	2630
-	- 127	+ 147	+ 8	194	469	508	+	+ 203	- 82	+ 142	261	1040	1060
+	+ 129	+ 184	- 89	241	333	402	-	- 404	+ 252	+ 138	496	1540	1610
68. p	- 441	- 78	- 700	831	1135	1220	+	+ 249	+ 46	- 391	466	140	290
p	+ 427	+1113	+ 804	1438	5980	6100	-	+ 192	+ 347	- 54	400	843	932
+	- 261	- 155	- 608	679	72	312	^{80.} p	+ 226	+ 49	+ 472	526	3670	3680
-	+ 275	- 879	+ 503	1050	3710	3820	+	- 92	- 370	+ 152	410	1380	1440
69. p	- 311	+ 532	- 155	636	2230	2310	+	+ 48	+ 6	- 957	958	- 171	178
+	- 296	- 64	- 29	304	700	763	-	+ 285	+ 324	- 504	664	297	523
+	+ 111	- 39	- 160	198	160	199	^{81.} p	+ 80	+ 121	+ 404	429	3400	3404
-	- 620	- 473	+ 76	783	2030	2170	-	+ 121	+ 174	- 123	245	341	401
70. p	- 267	+ 542	- 338	693	1840	1940	+	- 111	- 145	- 89	203	345	390
+	- 407	+ 293	- 103	512	966	1090	+	+ 238	+ 67	- 502	560	71	257
-	+ 190	- 238	+ 140	335	1100	1150	+	+ 293	+ 104	+ 185	362	1360	1400
+	+ 496	+ 510	+ 877	1129	4830	4880	-	-1037	+ 87	+ 265	1073	3160	3330
71. p	- 799	+ 84	+ 422	908	4130	4210	^{82.} p	+ 864	+ 328	- 378	998	2210	2400
+	+1511	-1022	+ 7	1824	4240	4610	+	- 513	-1211	+1088	1707	6690	6820
72. p	+ 33	- 120	- 203	238	1730	1730	^{83.} p	- 406	+ 452	+ 93	614	2820	2890
-	+ 235	- 120	+ 34	266	780	823	p	- 230	- 417	+ 548	726	4120	4140
+	+ 120	+ 155	- 283	344	147	245	+	+ 339	- 674	- 176	773	1370	1560
+	+ 227	+ 383	- 84	453	885	991	-	- 104	+ 494	- 49	507	1090	1200
73. p	+ 668	- 841	- 515	1191	2210	2450	^{84.} p	+ 31	+ 407	- 75	415	822	918
p	- 380	- 526	- 629	904	1420	1570	+	- 363	+ 304	+1429	1505	7080	7100
+	- 44	+ 291	+ 560	632	2900	2920	^{85.} p	+ 584	- 387	+ 510	867	4230	4290
-	- 245	+1076	+ 585	1249	4370	4510	p	- 779	+ 177	- 326	863	2120	2270
74. p	+ 94	+ 220	+ 224	328	2860	2870	+	+ 745	+ 675	- 380	1075	1540	1840
p	+ 429	- 203	+ 848	973	5250	5280	-	- 348	- 241	+ 475	637	2700	2730
75. p	- 26	- 191	- 356	405	1470	1480	^{86.} p	- 484	- 118	+ 627	801	4420	4450
+	-1504	- 335	+ 327	1575	4470	4730	-	- 164	- 143	+ 45	222	720	752
							+	+ 459	+1387	+ 105	1465	3658	3940
							+	+ 63	- 128	- 63	+ 156	316	347

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P
87. p	+ 438	+ 372	- 294	646	1890	1980							
	+ - 7	+ 275	- 327	428	216	350							
	- - 23	+ 270	+ 84	283	940	978	98. p	- 426	- 790	- 594	1077	1800	2010
	+ + 98	- 476	+ 123	502	1510	1590		- - 33	+ 326	- 76	336	649	727
88. p	+ 611	+ 185	-1022	1205	231	679		+ + 457	+ 64	+ 575	737	3180	3210
	+ - 499	+ 462	+ 674	957	3930	3990		+ - 225	+ 35	+ 148	272	1080	1100
89. p	- 239	- 718	-1028	1277	1070	1310	99. p	- 413	- 264	+ 468	678	3850	3880
	+ + 113	+ 6	+ 457	471	2280	2290		+ - 186	+ 11	+ 475	511	2420	2430
	+ - 6	+ 363	+ 171	402	1410	1460		- + 154	+ 311	+ 123	368	1220	1270
	- - 82	+ 292	+ 54	308	919	967		+ - 120	- 972	- 447	1076	1380	1690
90. p	- 7	+ 277	- 207	346	1790	1812	100. p	+ 219	- 426	- 443	652	1530	1600
	p - 295	- 63	+ 72	311	2460	2480		- + 92	- 550	+ 202	593	1880	1960
	+ + 255	+ 155	- 240	383	338	451		+ - 52	+1246	- 439	1323	1950	2310
	- + 467	- 189	- 544	741	372	626		+ + 17	+ 312	- 48	316	602	678
	- - 274	- 229	+ 702	788	3600	3630	101. p	- 51	- 574	+ 716	919	4830	4860
	+ + 244	+ 137	+ 51	284	862	906		p + 935	+ 362	- 431	1091	2240	2450
91. p	- 288	+ 582	- 244	694	2080	2180		- - 136	- 191	+ 594	638	3000	3010
	+ + 80	+ 117	- 104	176	258	294		+ - 673	+ 282	- 824	1100	487	877
	- + 426	+ 406	+ 208	624	2000	2080	102. p	- 305	- 715	+ 0	778	2820	2920
	+ + 199	- 903	- 669	1141	988	1350		+ + 127	- 461	- 22	479	1100	1200
92. p	+ 782	- 68	+ 192	808	3340	3430		+ + 40	- 8	- 30	51	268	271
	p - 280	- 74	- 226	367	1770	1790		- + 982	+ 518	+ 383	1175	3690	3860
	+ + 19	- 17	+ 278	279	1420	1420	103. p	+ 435	- 70	- 664	797	1180	1260
	- -1217	+ 257	+ 272	1274	3640	3850		+ - 32	+ 199	- 21	203	515	553
93. p	+ 130	- 41	+ 340	366	3180	3180		- - 0	+ 180	+ 52	187	671	695
	+ + 260	- 139	+ 333	444	1910	1940		+ - 511	- 113	- 315	611	654	838
94. p	+ 439	- 328	+ 299	624	3360	3400	104. p	- 178	+ 78	+ 675	703	4400	4410
	p - 89	+ 439	- 235	506	1870	1920		+ + 120	+ 627	- 313	711	886	1090
95. p	+ 922	+ 120	+1045	1399	6520	6580	105. p	- 73	+ 196	+ 306	371	3100	3110
	+ + 30	- 6	+ 189	191	1020	1020		- + 524	- 2	+ 2	524	1260	1360
	- - 554	+ 469	- 674	990	612	949		+ + 195	+ 148	- 99	264	441	504
	+ - 165	- 312	+ 44	356	929	993		+ - 108	+ 63	+ 196	232	1120	1130
96. p	- 544	- 430	+ 493	851	4160	4220		- - 960	- 141	+ 152	982	2670	2840
	p + 108	- 437	- 684	819	1160	1240		+ - 206	- 283	- 188	397	499	609
	- - 315	- 491	+ 42	584	1490	1600	106. p	- 467	+ 121	- 70	487	2270	2320
	+ + 541	+ 811	- 11	975	2240	2450		p - 450	+ 584	+ 315	802	3640	3720
97. p	- 294	+ 89	+ 530	612	3920	3930		+ - 272	+ 168	+ 125	343	1170	1210
	p + 106	- 119	- 489	514	1240	1250		- +1372	- 784	- 624	1699	2360	2840
	+ - 124	- 248	+ 324	426	1850	1870	107. p	- 236	+ 47	- 425	488	1380	1400
	- + 908	- 189	- 767	1204	870	1270		+ - 119	- 189	+ 94	242	882	910
								+ - 51	+ 312	- 52	320	677	747
								- + 36	+ 687	+ 514	859	3300	3370

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P
108. p	+ 340	+ 87	- 136	377	1990	2020							
p	+ 532	+ 269	+ 348	690	3560	3610							
	+ - 210	+ 552	- 814	1006	296	660	119. p	+ 197	- 870	+ 286	937	3780	3880
	- - 558	- 574	- 9	801	1850	2020		+ 291	- 96	+ 187	359	1360	1390
109. p	+ 388	- 139	- 590	720	1250	1310		- - 6	+ 293	+ 738	794	3720	3730
p	+ 60	- 768	+ 322	835	3710	3790		+ - 599	+ 186	- 883	1083	299	695
	+ - 153	- 112	- 297	353	128	229	120. p	- 315	+ 392	- 339	607	1730	1800
	- - 9	+ 371	- 48	374	800	882		p + 701	- 696	- 84	991	2940	3100
	- - 301	+ 262	- 289	493	455	605		+ - 327	+ 837	- 23	899	2040	2230
	+ + 15	+ 386	+ 903	982	4560	4580		- + 52	- 260	+ 17	266	735	781
110. p	+ 115	- 936	- 60	945	2920	3070	121. p	+ 149	- 294	+ 393	513	3460	3470
	- + 100	+ 319	- 112	353	595	682		p + 352	+ 393	+ 106	538	2770	2820
	+ + 449	- 794	- 57	914	1990	2190		- - 201	+ 94	+ 8	222	625	663
	+ - 252	+ 407	+ 23	480	1210	1300		+ + 238	- 248	- 929	990	- 28	345
111. p	- 170	+ 92	- 627	657	1070	1090	122. p	- 161	- 693	- 38	713	2620	2720
	- + 67	- 88	+ 74	133	631	641		+ + 275	+ 256	- 74	383	753	842
	+ + 188	- 119	- 320	390	153	270	123. p	- 18	- 94	+ 241	260	2860	2860
	+ + 980	- 109	- 289	1028	1670	1940		+ - 961	+ 782	+ 260	1265	3590	3800
112. p	+ 580	- 598	+ 269	875	3640	3730	124. p	- 695	+ 370	- 316	848	2130	2270
	+ - 112	- 113	+ 199	255	1170	1180		- + 30	+ 67	- 38	83	278	289
	+ + 130	+ 137	- 282	340	140	235		+ - 345	- 85	+ 89	369	1130	1190
	- - 536	- 188	- 498	755	520	770		+ - 747	+ 188	+ 289	814	2650	2760
113. p	- 702	+ 235	- 243	779	2210	2330	125. p	+ 941	- 829	- 389	1312	2750	3020
	+ - 311	+ 322	- 311	545	517	684		+ - 94	+ 106	+ 400	424	2040	2050
	- + 226	- 73	- 37	240	548	597	126. p	+ + 570	- 390	- 348	774	938	1165
	+ - 266	- 42	- 9	269	598	656		+ - 677	+ 537	+ 881	1234	5080	5150
	- + 454	+ 136	+ 147	496	1510	1590	127. p	- 167	- 349	- 45	389	2230	2270
	+ + 585	+ 169	- 12	609	1380	1500		p - 1	+ 180	+ 47	186	2330	2340
114. p	- 87	+ 440	+ 61	453	2560	2600		+ - 97	+ 158	+ 134	229	955	973
	+ + 205	+ 130	+ 217	326	1360	1380		- - 572	+ 279	+ 182	662	2020	2120
	- - 248	- 201	+ 73	327	1000	1060	128. p	- 844	+ 506	+ 214	1007	3720	3840
	+ + 164	- 808	- 134	835	1620	1810		p - 456	+ 633	+ 678	1034	4930	4990
115. p	+ 643	+ 80	- 386	754	1810	1920	129. p	- 53	+ 124	- 30	138	2120	2120
	+ + 343	+ 41	- 994	1052	- 49	349		p + 143	+ 154	+ 203	292	2780	2790
116. p	- 462	- 1287	- 353	1413	3030	3320		- + 686	- 545	+ 494	998	3570	3670
	+ - 93	- 190	+ 190	285	1210	1230		+ - 779	+ 411	- 785	1180	767	1170
117. p	+ 133	+ 564	+ 604	837	4420	4460	130. p	- 60	+ 45	- 103	127	1930	1930
	p + 333	+ 284	+ 607	748	4300	4320		- p - 498	+ 357	- 87	619	2380	2460
	+ - 46	- 133	+ 24	142	522	541		+ - 459	- 13	- 86	467	910	1020
	- + 399	+ 36	- 38	402	887	973		- - 60	- 215	+ 523	569	2670	2680
118. p	+ 442	- 31	- 641	779	1210	1290		+ + 528	+ 79	+ 95	543	1530	1620

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P					
131.	+	-	6	-1316	+ 389	1372	4160	4360	143.	+	+	0	- 149	+ 187	239	1110	1120	
	+	-	28	+ 154	- 459	485	11	157		+	+	295	- 61	- 571	646	89	314	
132.	p	-	968	- 364	+ 240	1061	3870	4010	144.	p	+	+	590	- 587	+ 854	1193	5650	5710
	+	+	330	- 260	+ 71	427	1210	1280		+	+	28	+ 62	- 55	88	242	251	
133.	p	+	301	+ 938	+ 152	996	3540	3670		+	+	434	+ 583	- 15	728	1670	1820	
	-	-	101	- 69	+ 78	145	1300	1310		-	+	177	+ 298	- 205	403	469	583	
	+	+	106	- 508	- 371	637	573	773	145.	p	+	+	466	- 123	- 168	511	2040	2100
	+	-	74	- 523	+ 696	874	3790	3830		p	-	-	392	- 438	- 378	699	1750	1850
134.	p	+	474	+ 211	+ 128	535	2820	2870		+	+	183	+ 400	- 63	444	913	1010	
	+	-	116	+ 117	+ 190	252	1130	1140		+	+	274	- 69	+ 506	580	2650	2660	
135.	p	-	500	+ 116	+ 671	843	4600	4630		-	-	273	- 100	- 233	372	333	442	
	+	-	185	- 447	+ 356	600	2320	2370		-	-	258	+ 206	+ 601	686	3130	3140	
136.	p	-	318	- 265	- 614	740	1220	1280	146.	p	-	-	168	+ 493	- 407	661	1630	1710
	-	-	362	+ 431	+ 176	589	1840	1920		+	+	132	+ 421	+ 186	478	1620	1680	
	+	-	326	- 755	- 5	822	1910	2080		-	+	548	- 262	- 64	611	1280	1420	
	+	-	179	+ 135	+ 96	244	890	918		+	-	758	- 337	- 651	1055	817	1165	
137.	p	+	244	+ 57	- 904	938	793	832	147.	p	-	-	199	+ 501	+ 291	612	3320	3360
	p	-	246	+ 161	+ 651	714	4360	4370		+	-	243	-1040	-1430	1785	536	1200	
	+	+	311	+ 191	- 192	412	521	636	148.	p	-	-	1271	+ 219	+ 206	1306	4230	4420
	-	-	914	- 548	+ 178	1080	2960	3150		+	-	243	- 253	+ 618	710	3220	3240	
138.	-	+	180	+ 124	+ 32	221	684	718	149.	p	+	+	774	+ 140	+ 13	786	2860	2970
	+	-	273	- 55	- 674	730	19	279		p	-	-	923	+ 805	+ 234	1247	4190	4370
	+	+	59	+ 241	- 117	274	418	486		+	+	636	+ 105	- 125	656	1230	1390	
	+	+	84	- 70	+ 35	115	513	525		-	-	184	- 668	+ 99	700	1890	2020	
139.	p	-	447	- 347	- 274	629	1920	2000	150.	p	-	-	376	+ 122	+ 404	566	3550	3570
	+	-	100	+ 176	+ 805	830	3970	3970		-	+	171	+ 141	+ 147	266	1060	1080	
	-	+	268	- 233	- 238	427	441	566		+	+	1120	- 225	+ 361	1198	3690	3860	
	+	+	185	- 371	- 158	444	678	794		+	-	211	- 565	- 390	718	711	932	
140.	p	-	631	- 458	+ 672	1029	4910	4970	151.	p	-	-	544	- 430	+ 493	851	4160	4220
	p	-	116	- 180	+ 23	215	2280	2290		p	+	+	108	- 437	- 684	819	1160	1240
	+	+	81	- 141	+ 338	375	1770	1780		-	-	315	- 491	+ 42	584	1490	1600	
	-	+	436	+ 445	+ 193	652	2030	2120		+	+	541	+ 811	- 11	975	2250	2450	
141.	p	+	883	- 906	+ 458	1345	4940	5100	152.	p	+	+	196	- 225	+ 97	313	2530	2550
	p	+	271	+ 605	-1117	1299	889	1110		+	-	82	+ 198	- 23	216	535	576	
	+	-	18	+ 64	+ 49	83	499	503		+	+	108	+ 531	+ 672	863	3710	3750	
	-	-	1136	+ 237	+ 610	1311	4580	4720		-	-	630	- 290	-1032	1242	289	750	
142.	p	-	334	+ 377	+ 163	529	2900	2940	153.	p	-	-	862	- 708	- 156	1126	2990	3190
	p	+	375	- 740	- 758	1123	1470	1690		-	+	106	+ 120	+ 255	301	1410	1420	
	+	+	182	+ 31	+ 138	230	968	985		+	+	47	+ 534	- 568	781	403	671	
	-	+	81	- 52	+ 258	275	1360	1365		+	+	512	+ 79	+ 162	543	1700	1780	

	P_x^*	P_y^*	P_z^*	P^*	P_z	P		P_x^*	P_y^*	P_z^*	P^*	P_z	P
154. p	+ 680	+ 79	+ 134	698	3040	3110	166. p	+ 658	+ 800	+ 748	1278	5540	5640
p	- 843	- 87	+ 39	849	3020	3130	p	- 661	- 0	- 646	925	1420	1560
-	- 311	- 283	- 392	575	381	568	+	+ 227	+ 21	+ 304	380	1700	1710
+	- 281	+ 401	+ 91	498	1420	1500	-	+ 265	- 262	- 510	631	210	428
155. +	+ 337	+ 95	- 166	387	533	638	167. p	+ 217	- 81	- 375	441	1450	1470
+	- 542	- 253	- 329	683	777	980	+	- 219	+ 74	- 138	270	354	424
156. p	- 515	- 399	+ 878	1093	5530	5570	-	- 572	- 645	+ 214	888	2613	2751
+	+ 153	- 238	- 49	287	617	679	+	- 431	+ 419	+ 437	743	2840	2910
-	- 394	+ 390	- 252	609	809	980	168. p	+ 7	+ 340	+ 498	603	3830	3850
+	+ 128	- 8	+ 517	533	2570	2580	+	- 215	- 32	- 30	219	525	568
157. p	- 23	+1331	- 515	1428	2650	2960	+	- 242	- 152	+ 123	311	1100	1130
+	+ 159	+ 49	+ 13	167	536	561	-	- 536	+ 800	- 380	1035	1450	1740
-	- 298	+ 389	+ 71	496	1370	1450	169. p	+ 408	- 373	- 680	876	1250	1370
+	+ 364	- 551	+ 254	707	2300	2400	+	+ 75	+ 58	- 18	96	347	360
158. p	- 504	- 579	+ 25	768	2860	2960	+	- 510	+ 362	+ 511	807	3180	3240
+	+ 197	+ 569	+ 68	606	1610	1720	-	+ 516	+ 189	- 802	972	249	603
159. p	- 881	- 288	+ 107	933	3330	3450	170. p	- 300	+ 868	- 305	967	2340	2520
+	+ 145	+ 375	+ 41	404	1090	1160	-	- 132	+ 240	- 35	276	626	684
-	+1027	+ 271	- 22	1062	2420	2640	+	+ 129	- 674	- 53	689	1490	1640
+	- 392	- 710	+ 186	832	2420	2550	+	+ 845	- 238	+ 606	1066	4000	4100
160. p	+ 786	- 392	- 694	1119	1630	1850	171. p	-1111	-1445	- 339	1854	3950	4350
+	+ 116	+ 124	+ 46	176	635	657	p	+ 469	+ 477	+ 12	669	2690	2770
-	+ 152	+ 9	+ 94	179	762	777	-	+ 51	- 57	+ 24	80	432	439
+	- 45	+ 174	+ 961	978	4690	4700	+	+ 351	+ 281	+ 257	518	1885	1940
161. p	+ 641	- 190	- 109	677	2400	2490	172. p	+ 76	- 176	- 389	434	1410	1420
p	-1011	- 624	+ 935	1512	6460	6560	+	- 8	- 369	- 214	427	497	620
+	+ 199	+ 201	- 192	342	369	465	+	- 661	- 369	+ 246	796	2460	2570
-	- 45	+ 40	- 760	762	- 123	137	-	+ 207	+ 495	+ 96	540	1530	1620
162. p	- 79	- 412	+ 691	808	4600	4620	173. p	- 892	+ 310	- 272	982	2450	2630
-	+ 180	+ 197	+ 250	365	1530	1550	-	+ 203	+ 208	- 84	303	559	630
+	+ 120	- 32	- 360	379	32	128	+	- 43	+ 303	- 103	323	553	632
+	- 888	- 318	- 672	1158	1000	1370	+	+ 861	- 243	- 53	896	1960	2150
163. p	- 44	+ 502	- 597	781	1320	1410	174. p	+ 502	+ 618	- 609	1002	1640	1820
+	+1176	-1316	+ 327	1794	4970	5280	-	- 11	+ 39	+ 258	261	1330	1330
164. p	+ 742	+ 366	- 73	831	2710	2830	+	- 292	- 383	- 169	510	796	930
+	- 317	- 289	- 416	597	370	566	+	- 68	- 553	+ 125	571	1670	1760
-	+ 295	+ 36	+ 829	881	4140	4150	175. p	+ 520	+ 21	+ 223	567	3090	3140
+	- 618	- 140	+ 25	634	1560	1680	+	+ 457	- 109	- 52	474	1010	1110
165. p	+ 510	- 176	+ 389	665	3640	3680	-	+ 36	- 174	+ 74	192	735	756
-	+ 51	- 38	+ 202	212	1100	1100	+	- 591	- 233	- 8	636	1480	1610

	P_x	P_y	P_z	P^*	P_z	P	P_x^*	P_y^*	P_z^*	P^*	P_z	P	
							186. p	+ 493	+ 881	- 288	1050	2530	2720
							-	- 94	- 280	+ 22	297	813	865
176. p	- 42	+ 159	- 407	438	1370	1380	f	- 538	+ 5	+ 398	669	2580	2630
	+ - 567	- 846	+ 951	1393	5620	5713	+	- 76	+ 330	+ 410	532	2300	2320
177. p	+ 74	+ 266	+1152	1185	6390	6390	187. p	+ 179	- 12	- 150	234	1860	1870
	+ 115	+ 369	- 678	781	126	406	-	+ 77	- 140	+ 121	200	867	882
178. p	- 519	+ 434	+ 257	723	3380	3450	+	- 64	+ 124	- 302	333	74	158
	p + 478	+ 342	- 237	634	2020	2100	+	-1098	- 23	+ 494	1204	4040	4180
	- 424	- 17	+ 204	471	1650	1700	188. p	- 286	+ 151	- 276	425	1690	1720
	+ - 105	+ 55	+ 130	176	846	854	+	- 202	+ 50	+ 196	286	1230	1250
	- 219	- 166	- 559	623	68	283	+	- 30	- 100	+ 764	771	3730	3730
	+ 788	- 648	+ 204	1040	2940	3110	-	+ 159	+ 793	- 604	1009	833	1160
179. +	- 92	- 84	- 34	129	353	374	189. +	+ 140	- 90	+ 224	279	1280	1290
	- + 12	+ 559	+ 73	564	1520	1620	+	- 412	+ 836	- 754	1199	888	1290
	+ - 243	+ 122	- 564	626	62	279	190. +	+ 14	- 0	- 5	14	311	311
	+ - 63	- 0	+ 626	629	3060	3060	+	+ 865	- 101	- 309	924	1380	1630
180. p	- 331	+ 536	+ 568	849	4350	4390	191. p	+ 406	+ 57	- 291	503	1730	1780
	+ - 72	+ 87	- 154	191	160	196	+	+ 39	+ 10	+ 210	214	1120	1120
	- + 189	+ 595	- 577	850	536	823	+	+ 152	+ 273	+ 204	373	1430	1470
	+ - 334	- 889	+ 98	955	2470	2650	-	- 12	- 653	- 413	773	774	1010
181. p	+1081	+ 58	+ 477	1183	4690	4810	192. p	+ 658	- 700	- 395	1038	2240	2440
	p + 101	- 418	+ 347	553	3390	3410	p	-1487	+ 199	+ 682	1648	6090	6280
	+ + 113	+ 271	- 661	724	39	296	193. p	+ 338	+ 43	- 95	354	2080	2110
	- - 363	- 129	- 51	389	824	910	p	+ 548	- 281	+ 689	924	4770	4810
182. p	+ 541	+ 326	+ 226	671	3230	3290	+	+ 198	+ 240	+ 162	351	1280	1320
	- 201	- 142	+ 60	253	818	854	-	- 108	+ 258	+ 346	445	1930	1950
	+ + 41	- 135	- 95	170	271	306	+	- 502	- 259	- 782	964	281	631
	+ - 135	+ 380	- 11	404	957	1040	-	- 474	- 1	- 321	572	554	729
183. p	-1133	+ 644	+ 54	1305	3840	4060	194. p	+ 230	+ 470	+ 744	909	4890	4920
	+ - 99	+ 228	- 206	323	294	385	-	+ 75	- 325	- 33	335	754	824
	- + 775	+ 358	+ 105	860	2280	2430	+	+ 497	- 206	- 64	541	1130	1250
	+ + 18	- 420	- 296	514	485	642	+	- 683	+ 628	+ 230	956	2810	2960
184. p	+ 918	+ 179	+ 189	954	3570	3690	195. p	+ 842	+ 372	- 458	1028	2060	2260
	+ - 388	+ 211	- 15	442	1030	1120	p	- 340	- 548	+ 254	693	3330	3390
185. p	+ 294	+ 253	- 129	409	2040	2080	+	- 56	- 50	+ 48	89	504	510
	+ + 83	+ 18	- 344	354	14	86	+	- 114	+ 116	- 88	185	314	354
	+ - 541	- 8	+ 178	570	1800	1880	-	+ 390	+ 32	- 545	672	213	446
	+ - 248	- 146	- 240	375	322	432	-	- 722	+ 77	+ 789	1072	4480	4540
	- + 160	- 83	+ 44	185	649	674	196. p	+ 96	+ 100	- 31	142	2110	2120
	- + 514	- 267	+ 142	596	1770	1860	p	- 13	+ 932	- 897	1293	1430	1710
							+	- 30	+ 335	+ 27	338	912	972
							-	- 139	- 952	+ 952	1354	5530	5620

	P_x^*	P_y^*	P_z^*	P^*	P_z	P	P_x^*	P_y^*	P_z^*	P^*	P_z	P	
<u>197.</u> P	- 474	+ 757	+ 100	899	3250	3370							
p	- 592	- 446	-1172	1387	919	1180							
-	+ 539	- 18	+ 976	1115	5050	5080	<u>200.</u> P	- 348	+ 272	- 371	577	1610	1670
+	+ 527	- 293	+ 96	610	1680	1790	-	+ 187	+ 59	- 237	308	183	268
<u>198.</u> P	- 202	+ 506	- 752	929	1160	1280	+	+ 307	- 667	- 277	785	1140	1360
p	+ 1	- 248	+ 545	599	3940	3950	+	- 490	- 500	+ 485	852	3200	3280
-	- 243	- 141	- 5	281	712	765	<u>201.</u> P	- 702	- 174	- 117	733	2460	2560
+	+ 572	+ 420	+ 226	745	2320	2420	+	- 196	- 46	- 193	279	235	309
<u>199.</u> P	+ 107	- 336	- 173	392	1915	1950	+	+ 78	- 56	- 179	203	120	154
p	- 230	+ 374	+1277	1350	7000	7020	-	- 69	- 171	+ 177	255	1120	1130
+	- 10	- 191	- 366	413	85	209							
-	- 64	+ 35	- 834	838	- 139	157							

Appendix

TABLE I
Efficiency of Sampling

Reaction	n	γ	Efficiency	Reaction	n	γ	Efficiency
$N'N\pi$	3	-	30 ± 15	$\tilde{N}3N$	4	2	10 ± 9
$N'N'\pi$	3	-	30 ± 14	$N'N3\pi$	5	3	10 ± 2
$NN\pi$	3	-	40 ± 30	$N'N'3\pi$	5	3	10 ± 3
$N'N'2\pi$	4	2	10 ± 2	$NN3\pi$	5	3	15 ± 5
$N'N'2\pi$	4	2	10 ± 2	$N'N'4\pi$	6	3	5 ± 3
$NN2\pi$	4	2	10 ± 3	$NN4\pi$	6	3	10 ± 6
$\tilde{N}'3N$	4	2	25 ± 17	$N'N4\pi$	6	4	6 ± 2

($\tilde{P}_1, \dots, \tilde{P}_{n-\gamma}$ - were chosen using the rejection technique
 $\tilde{P}_{n-\gamma+1}, \dots, \tilde{P}_{n-1}$ - using the direct method)

TABLE 2

the number	Reactions taken into account			
	Products of the reaction	Statistical weight	Statistical weight normalized to 100	Number of events in the table
1	$N'N2\pi$	15,17	19,19	38
2	$N'N3\pi$	13,92	17,61	35
3	$N'N'2\pi$	11,70	14,80	30
4	$N'N'3\pi$	9,28	11,74	24
5	$NN3\pi$	6,21	7,86	15
6	$NN2\pi$	5,25	6,64	14
7	$N'N4\pi$	4,43	5,60	11
8	$N'N'\pi$	3,53	4,46	9
9	$N'N\pi$	3,15	3,98	8
10	$N'N'4\pi$	2,09	2,64	5
11	$NN4\pi$	2,04	2,58	5
12	$NN\pi$	1,00	1,26	3
13	$\tilde{N}'3N$	0,80	1,01	2
14	$\tilde{N}3N$	0,49	0,62	2

TOTAL: 79,06 100 201

Rejected reaction

number	weight	number	weight
15	$N'N5\pi$ 0,54	20	$N'\tilde{N}2N$ 0,27
16	$N'N'5\pi$ 0,49	21	$\tilde{N}3N\pi$ 0,20
17	$NN5\pi$ 0,47	22	$\tilde{N}'3N\pi$ 0,18
18	$\tilde{N}'N'2N$ 0,36	23	$N'\tilde{N}2N\pi$ 0,06
19	$\tilde{N}N2N'$ 0,36		и т.д.

TABLE 3
Coefficients $c_k^{(1)}$

$T \backslash T_k = t_k$	0	1	1	1	2	2	2	2	2	3	3	3	3	3	3	3			
0	0	$\frac{1}{3}^+$	$\frac{1}{3}$	$\frac{1}{3}^-$															
1	-1	1^-	$\frac{1}{2}$	$\frac{1}{2}^-$	$\frac{3}{5}^+$	$\frac{3}{10}$	$\frac{1}{10}^-$												
1	0	1	$\frac{1}{2}^+$	$\frac{1}{2}^-$	$\frac{3}{10}$	$\frac{2}{5}$	$\frac{3}{10}^-$												
1	1	1^+	$\frac{1}{2}^+$	$\frac{1}{2}$		$\frac{1}{10}^+$	$\frac{3}{10}$	$\frac{3}{5}^-$											
2	-2		1^-		$\frac{2}{3}$	$\frac{1}{3}^-$			$\frac{5}{7}^+$	$\frac{5}{21}$	$\frac{1}{21}^-$								
2	-1		$\frac{1}{2}$	$\frac{1}{2}^-$	$\frac{1}{3}^+$	$\frac{1}{6}$	$\frac{1}{2}^-$		$\frac{10}{21}^+$	$\frac{8}{21}$	$\frac{1}{7}^-$								
2	0		$\frac{1}{6}$	$\frac{2}{3}$	$\frac{1}{6}^-$	$\frac{1}{2}^+$	$\frac{1}{2}^-$			$\frac{2}{7}^+$	$\frac{3}{7}$	$\frac{2}{7}^-$							
2	1		$\frac{1}{2}$	$\frac{1}{2}$		$\frac{1}{2}^+$	$\frac{1}{6}$	$\frac{1}{3}^-$		$\frac{1}{7}^+$	$\frac{8}{21}$	$\frac{10}{21}^-$							
2	2			1^+			$\frac{1}{3}$	$\frac{2}{3}$			$\frac{1}{21}^+$	$\frac{5}{21}$	$\frac{5}{7}^-$						
$T \backslash T_k = t_k$	2	2	2	2	2	3	3	3	3	3	3	3	4	4	4	4	4	4	4
3	-3	1^-				$\frac{3}{4}$	$\frac{1}{4}^-$						$\frac{7}{9}^+$	$\frac{7}{36}$	$\frac{1}{36}^-$				
3	-2	$\frac{1}{3}$	$\frac{2}{3}^-$			$\frac{1}{4}$	$\frac{1}{3}$	$\frac{5}{12}$						$\frac{7}{12}^+$		$\frac{1}{12}^-$			
3	-1	$\frac{1}{3}^+$	$\frac{8}{15}$	$\frac{2}{3}$		$\frac{5}{12}$	$\frac{1}{12}$	$\frac{1}{2}^-$							$\frac{1}{12}$	$\frac{5}{12}$	$\frac{1}{6}^-$		
3	0		$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}^-$		$\frac{1}{2}$	$\frac{1}{2}^-$							$\frac{5}{18}$	$\frac{4}{9}$	$\frac{5}{18}^-$		
3	1		$\frac{8}{15}$	$\frac{1}{3}$	$\frac{1}{15}^-$		$\frac{1}{2}$	$\frac{1}{12}$	$\frac{5}{12}$						$\frac{5}{12}$	$\frac{1}{6}^+$	$\frac{5}{18}$	$\frac{5}{12}^-$	
3	2		$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$		$\frac{1}{12}$	$\frac{5}{12}$	$\frac{1}{4}^-$						$\frac{1}{12}$	$\frac{5}{12}$	$\frac{7}{12}$	$\frac{7}{12}^-$	
3	3			1^+			$\frac{1}{4}$	$\frac{3}{4}$							$\frac{1}{36}$	$\frac{1}{36}$	$\frac{7}{36}$	$\frac{7}{9}^-$	

Notations
(taking as an example
the 1 line)
 $(0,0) = \frac{1}{3}(I, -I; \pi^+) +$
 $+ \frac{1}{3}(I, 0; \pi^0) +$
 $+ \frac{1}{3}(I, I; \pi^-).$

TABLE 3
(continuation)

$T \backslash T_x$	3	3	3	3	3	3	3	4	4	4	4	4	4	4	4	4	5	5	5	5	5	5	5	5	5	5	
t	-3	-2	-1	0	1	2	3	-4	-3	-2	-1	0	1	2	3	4	-5	-4	-3	-2	-1	0	1	2	3	4	5
4 -4	1^-							$\frac{4}{5}$	$\frac{1^-}{5}$							$\frac{9^+}{11}$	$\frac{9}{55}$	$\frac{1^-}{55}$									
4 -3	$\frac{1}{4}$	$\frac{3^-}{4}$						$\frac{1^+}{5}$	$\frac{9}{20}$	$\frac{7^-}{20}$							$\frac{36^+}{55}$	$\frac{16}{55}$	$\frac{3^-}{55}$								
4 -2	$\frac{1^+}{28}$	$\frac{3}{7}$	$\frac{15^-}{28}$					$\frac{7^+}{20}$	$\frac{1}{5}$	$\frac{9^-}{20}$							$\frac{28^+}{55}$	$\frac{21}{55}$	$\frac{6^-}{55}$								
4 -1		$\frac{3^+}{28}$	$\frac{15}{28}$	$\frac{5^-}{14}$				$\frac{9^+}{20}$	$\frac{1}{20}$	$\frac{1^-}{2}$								$\frac{21^+}{55}$	$\frac{24}{55}$	$\frac{2^-}{11}$							
4 0		$\frac{3^+}{14}$	$\frac{4}{7}$	$\frac{3^-}{14}$				$\frac{1^+}{2}$		$\frac{1^-}{2}$								$\frac{3^+}{11}$	$\frac{5}{11}$	$\frac{3^-}{11}$							
4 1			$\frac{5^+}{14}$	$\frac{15}{28}$	$\frac{3^-}{28}$			$\frac{1^+}{2}$	$\frac{1}{20}$	$\frac{9^-}{20}$								$\frac{2^+}{11}$	$\frac{24}{55}$	$\frac{21^-}{55}$							
4 2				$\frac{15^+}{28}$	$\frac{3}{7}$	$\frac{1^-}{28}$			$\frac{9^+}{20}$	$\frac{1}{5}$	$\frac{7^-}{20}$							$\frac{6^+}{55}$	$\frac{21}{55}$	$\frac{28^-}{55}$							
4 3					$\frac{3^+}{4}$	$\frac{1}{4}$				$\frac{7^+}{20}$	$\frac{9}{20}$	$\frac{1^-}{5}$								$\frac{3^+}{55}$	$\frac{16}{55}$	$\frac{36^-}{55}$					
4 4						1^+					$\frac{1^+}{5}$	$\frac{4}{5}$											$\frac{1^+}{55}$	$\frac{9}{55}$	$\frac{9^-}{11}$		

TABLE 4

Coefficients $C_k^{(n)}$ for the state (I, I)

n	T_k		mesons																
	t_k	t_k	0 0	1-1	1 0	1 1	2-2	2-1	2 0	2 1	2 2	3-3	3-2	3-1	3 0	3 1	3 2	3 3	
1	+	+	1		$\frac{1}{2}$				$\frac{1}{10}$										
	0	0				$\frac{1}{2}$				$\frac{3}{10}$									
	-	-									$\frac{3}{5}$								
2	+	+		$\frac{3}{5}$				$\frac{1}{5}$						$\frac{1}{35}$					
	+	0	1		$\frac{4}{5}$				$\frac{2}{5}$						$\frac{3}{35}$				
	+	-				$\frac{6}{5}$				$\frac{2}{5}$						$\frac{2}{35}$			
	0	0					$\frac{2}{5}$			$\frac{1}{5}$						$\frac{4}{35}$			
	0	-									$\frac{4}{5}$						$\frac{2}{7}$		
3	+	+						$\frac{3}{7}$						$\frac{3}{28}$					
	+	+		$\frac{6}{5}$					$\frac{24}{35}$						$\frac{33}{140}$				
	+	+	$\frac{9}{5}$		$\frac{6}{5}$					$\frac{18}{35}$						$\frac{9}{70}$			
	+	0	$\frac{6}{5}$		$\frac{5}{10}$					$\frac{39}{70}$						$\frac{9}{35}$			
	+	0				$\frac{12}{5}$					$\frac{48}{35}$					$\frac{33}{70}$			
	+	-									$\frac{9}{7}$						$\frac{9}{28}$		
	0	0					$\frac{3}{10}$				$\frac{3}{14}$					$\frac{3}{35}$			
	0	0										$\frac{33}{35}$					$\frac{3}{7}$		
	0	-																$\frac{27}{28}$	
	4	+	+												$\frac{1}{3}$				
+		+					$\frac{9}{7}$							$\frac{4}{7}$					
+		+		$\frac{12}{7}$					$\frac{172}{175}$						$\frac{2}{7}$				
+		+		$\frac{66}{35}$					$\frac{188}{175}$						$\frac{3}{5}$				
+		0	$\frac{24}{5}$		$\frac{132}{35}$					$\frac{78}{35}$						$\frac{32}{35}$			
+		+				$\frac{18}{7}$					$\frac{9}{7}$					$\frac{3}{7}$			
+		0		$\frac{6}{5}$		$\frac{36}{35}$				$\frac{24}{35}$						$\frac{38}{105}$			
+		0					$\frac{132}{35}$				$\frac{12}{5}$						$\frac{6}{5}$		
+		0										$\frac{27}{7}$						$\frac{12}{7}$	
0		0																	$\frac{4}{3}$
0	0	0				$\frac{9}{35}$					$\frac{6}{35}$					$\frac{4}{35}$			
	0	0										$\frac{36}{35}$					$\frac{4}{7}$		
	0	-																$\frac{4}{7}$	
	0	0																	$\frac{11}{7}$
	0	0																	

Example: $(1,1) = (0,0; \pi^+) + \frac{1}{2}(1,0; \pi^+) + \frac{1}{10}(2,0; \pi^+) + \frac{1}{2}(1,1; \pi^0) + \frac{3}{2}(2,1; \pi^0) + \frac{3}{2}(2,2; \pi^-)$

Probabilities of charge states
of the systems N'N (or N'N')

state	N'N														Tt		
	00	1-1	10	11	2-2	2-1	20	21	22	3-3	3-2	3-1	30	31		32	33
p ⁺ p									1								
p ⁺ n				$\frac{3}{4}$				$\frac{1}{4}$									
p ⁰ p				$\frac{1}{6}$				$\frac{1}{2}$									
p ⁰ n			$\frac{1}{3}$					$\frac{1}{3}$									
p ⁻ p			$\frac{1}{6}$					$\frac{1}{6}$									
p ⁻ n		$\frac{1}{12}$						$\frac{1}{4}$									
n ⁺ p				$\frac{1}{12}$				$\frac{1}{4}$									
n ⁺ n			$\frac{1}{6}$					$\frac{1}{6}$									
n ⁰ p			$\frac{1}{3}$					$\frac{1}{3}$									
n ⁰ n		$\frac{1}{6}$						$\frac{1}{2}$									
n ⁻ p		$\frac{1}{4}$						$\frac{1}{4}$									
n ⁻ n						1											
N'N'																	
p ⁺ p ⁺																	1
p ⁺ p ⁰									$\frac{2}{3}$							$\frac{2}{3}$	
p ⁺ p ⁻				$\frac{1}{5}$				$\frac{1}{3}$		$\frac{1}{3}$				$\frac{2}{15}$		$\frac{1}{3}$	
p ⁺ n ⁺									$\frac{1}{3}$								
p ⁺ n ⁰				$\frac{2}{5}$				$\frac{2}{3}$						$\frac{4}{15}$		$\frac{1}{3}$	
p ⁺ n ⁻	$\frac{1}{9}$		$\frac{9}{10}$					$\frac{1}{2}$					$\frac{1}{10}$				
p ⁰ p ⁰				$\frac{8}{45}$										$\frac{4}{15}$			
p ⁰ p ⁻	$\frac{1}{9}$		$\frac{1}{45}$					$\frac{1}{9}$					$\frac{1}{5}$				
p ⁰ n ⁺				$\frac{8}{45}$										$\frac{4}{15}$			
p ⁰ n ⁰	$\frac{2}{9}$		$\frac{2}{45}$					$\frac{2}{9}$					$\frac{2}{5}$				
p ⁰ n ⁻		$\frac{2}{45}$						$\frac{2}{3}$					$\frac{4}{15}$				
p ⁻ p ⁻		$\frac{2}{45}$											$\frac{1}{15}$				
p ⁻ n ⁺	$\frac{1}{18}$		$\frac{1}{90}$					$\frac{1}{18}$					$\frac{1}{10}$				
p ⁻ n ⁰		$\frac{8}{45}$											$\frac{4}{15}$				
p ⁻ n ⁻						$\frac{1}{3}$									$\frac{1}{15}$		
n ⁺ n ⁺				$\frac{2}{45}$													
n ⁺ n ⁰	$\frac{1}{9}$		$\frac{1}{45}$					$\frac{1}{9}$					$\frac{1}{5}$				
n ⁺ n ⁻		$\frac{1}{45}$						$\frac{1}{3}$					$\frac{2}{15}$				
n ⁰ n ⁰		$\frac{5}{45}$											$\frac{4}{15}$				
n ⁰ n ⁻													$\frac{2}{15}$				
n ⁻ n ⁻						$\frac{2}{3}$									$\frac{2}{3}$		

Example: $(2, -1) = \frac{2}{3}(p^0n^-) + \frac{1}{3}(n^+n^-)$

TABLE 6

Probabilities of charge states of the systems N^+Nn^- and $N^+N'n^-$
(S is a normalization factor)

1. $N^+N'n^-$; $S=1/5$

state	probabili- ty state	state	probabili- ty state	state	probabili- ty state
1 p^+n^-+	1	5 n^+n^0+	2/15	9 p^0n^+0	4/45
2 p^0p^-+	2/15	6 p^+p^-0	1/5	10 n^+n^+0	1/45
3 p^-n^++	1/15	7 p^+n^00	2/5	11 p^+p^0-	2/5
4 p^0n^0+	4/15	8 p^0p^00	4/45	12 p^+n^+-	1/5

2. $N^+N'^2n^-$; $S=1/7$

1 p^+n^-+0	10/7	9 p^0n^-++	8/21	17 p^+n^000	34/105
2 p^0p^-+0	4/21	10 n^+n^-++	4/21	18 p^0p^000	32/315
3 p^-n^++0	2/21	11 p^+p^-+-	8/21	19 p^0n^+00	32/315
4 p^0n^0+0	8/21	12 p^+n^0+-	16/21	20 n^+n^+00	8/315
5 n^+n^0+0	4/21	13 p^0p^0+-	8/35	21 p^+p^0-0	76/105
6 p^-p^-++	1/35	14 p^0n^++-	8/35	22 p^+n^+-0	38/105
7 p^-n^0++	4/35	15 n^+n^++-	2/35	23 p^+p^+--	3/7
8 n^0n^0++	4/35	16 p^+p^-00	17/105		

3. $N^+N'^3n^-$; $S=1/18$

1 p^+n^-+-	9/4	12 p^-n^0++0	29/105	23 p^0p^0000	8/105
2 p^0p^-+-	13/42	13 n^0n^0++0	29/210	24 p^0n^+000	8/105
3 p^-n^++-	13/84	14 p^0n^-++0	1	25 n^+n^+000	2/105
4 p^0n^0+-	13/21	15 n^+n^-++0	1/2	26 p^-n^-+++	5/28
5 n^+n^0+-	13/42	16 p^+p^-+-0	1	27 n^0n^-+++	5/14
6 p^+n^-+00	12/7	17 p^+n^0+-0	2	28 p^+p^0+-	15/14
7 p^0p^-+00	4/15	18 p^0p^0+-0	58/105	29 p^+n^++-	15/28
8 p^-n^++00	2/15	19 p^0n^++-0	58/105	30 p^+p^0-00	32/35
9 p^0n^0+00	8/15	20 n^+n^++-0	29/210	31 p^+n^+-00	16/35
10 n^+n^-00	4/15	21 p^+p^-000	1/7	32 p^+p^+-0	27/28
11 p^-p^-++0	29/420	22 p^+n^0000	2/7		

4. $N^*N^*4\pi$; $S = 1/46$

1	p^+n^-++0	7	16	p^-p^-++00	13/105	31	p^+p^-0000	13/105
2	p^0p^-++0	22/21	17	p^-p^0++00	52/105	32	p^+n^00000	26/105
3	p^-n^+--0	11/21	18	n^0n^0++00	52/105	33	p^0p^00000	8/105
4	p^0n^0++0	44/21	19	p^0n^-++00	856/525	34	p^0n^+0000	8/105
5	n^+n^0++0	22/21	20	n^+n^-++00	428/525	35	n^+n^+0000	2/105
6	p^+n^-+000	40/21	21	p^+p^-++-	1	36	p^-n^-+++0	13/21
7	p^0p^-+000	32/105	22	p^+n^0++-	2	37	n^0n^-+++0	26/21
8	p^-n^++000	16/105	23	p^0p^0++-	4/7	38	p^+p^0+-0	26/7
9	p^0n^0+000	64/105	24	p^0n^+--	4/7	39	p^+n^+--0	13/7
10	n^+n^0+000	32/105	25	n^+n^+--	1/7	40	p^+p^0-000	16/15
11	p^-p^-+++	2/21	26	p^+p^-+00	12/7	41	p^+n^+-000	8/15
12	p^-n^0+++	8/21	27	p^+n^0+00	24/7	42	p^+p^+-00	11/7
13	n^0n^0+++	8/21	28	p^0p^0+00	104/105	43	p^+p^-+--	4/3
14	p^0n^-+++	744/525	29	p^0n^+00	104/105	44	n^-n^-++++	1/3
15	n^+n^-+++	372/525	30	n^+n^+00	26/105			

5. $N^*N\pi$; $S=1/2$

1	p^0n^+	1/5	4	n^0p^+	1/5	7	n^+p^0	7/60
2	n^+n^+	1/10	5	p^+n^0	9/20	8	p^+p^-	3/5
3	p^-p^+	1/10	6	p^0p^0	7/30			

6. $N^*N2\pi$; $S=1/5$

1	p^+p^0-	4/5	6	p^0n^+0	2/5	11	n^+n^+0	1/5
2	p^+n^+-	1	7	p^-p^+0	1/5	12	n^0p^+0	2/5
3	p^+n^00	7/20	8	p^-n^++	1/10	13	n^0n^++	1/5
4	p^0p^+-	2/5	9	n^+p^+-	1/5	14	n^-p^++	1/2
5	p^0p^00	1/6	10	n^+p^00	1/12			

7. $N^*N3\pi$; $S=1/12$

1	p^-n^++0	19/70	8	p^0n^+++	4/7	15	p^+n^000	39/140
2	n^0n^++0	19/35	9	n^+n^+++	2/7	16	p^0p^000	11/70
3	n^-p^++0	15/14	10	p^-p^+++	2/7	17	n^+p^000	11/140
4	p^0n^+00	17/35	11	n^0p^+++	4/7	18	n^-n^+++	3/7
5	n^+n^+00	17/70	12	p^+n^+-0	15/7	19	p^+p^-00	33/35
6	p^-p^+00	17/70	13	p^0p^+-0	38/35	20	p^+p^+--	9/7
7	n^0p^+00	17/35	14	n^+p^+-0	19/35			

8. $N^*N \ 4\pi$; $S = 1/30$

1	$p^-n \ +\ +\ -$	67/175	10	$n^0p \ +\ +\ -0$	2	19	$p^0p \ +\ -00$	64/35
2	$n^0n \ +\ +\ -$	136/175	11	$p^0n \ +000$	4/7	20	$n^+p \ +\ -00$	32/35
3	$n^-p \ +\ +\ -$	268/175	12	$n^+n \ +000$	2/7	21	$p^+n \ 0000$	33/140
4	$p^-n \ +\ +00$	149/350	13	$p^-p \ +000$	2/7	22	$p^0p \ 0000$	9/70
5	$n^0n \ +\ +00$	149/175	14	$n^0p \ +000$	4/7	23	$n^+p \ 0000$	9/140
6	$n^-p \ +\ +00$	589/350	15	$p^+n \ +\ +\ -$	9/4	24	$n^-n \ +\ +\ +0$	9/7
7	$p^0n \ +\ +\ -0$	2	16	$p^0p \ +\ +\ -$	15/14	25	$p^+p \ +\ -\ -0$	27/7
8	$n^+n \ +\ +\ -0$	1	17	$n^+p \ +\ +\ -$	15/28	26	$p^+p \ -000$	36/35
9	$p^-p \ +\ +\ -0$	1	18	$p^+n \ +\ -00$	24/7			

9. $3 \ N \ \tilde{N}$; $g = 1/3$

$ppn\tilde{n}$	9/4
$ppp\tilde{p}$	3/4

10. $3 \ N \ \tilde{N}'$; $S = 1/3$

1	$ppn\tilde{n}^+$	9/5	4	$ppp\tilde{p}^0$	1/5
2	$ppn\tilde{n}^0$	3/5	5	$ppp\tilde{n}^-$	1/10
3	$ppn\tilde{p}^+$	3/10			

The Russian variant of this paper was received by Publishing Department on October 27, 1958.