

94-315



ОБЪЕДИНЕННЫЙ  
ИНСТИТУТ  
ЯДЕРНЫХ  
ИССЛЕДОВАНИЙ  
ДУБНА

E7-94-315

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SPONTANEOUS FISSION  
OF LIGHT CALIFORNIUM ISOTOPES  
PRODUCED IN THE  $^{206,207,208}\text{Pb} + ^{34,36}\text{S}$  REACTIONS;  
NEW NUCLIDE  $^{238}\text{Cf}$

Submitted to «Nuclear Physics A»

1994

## 1. Introduction

Until recently, the 35.7-h  $\alpha$ -decaying isotope  $^{246}\text{Cf}$  [1] having a tiny spontaneous fission (SF) branch,  $b_{sf}=2\times 10^{-6}$ , remained to be the lightest one among the Cf isotopes whose SF decay mode was established and characterized. At the same time, seven lighter Cf isotopes down to  $^{239}\text{Cf}$  were already produced and identified by their  $\alpha$  decay [1,2].

We report here on our experiments designed to probe the SF decay of the lighter known even-even isotopes of Cf and to produce new, even more neutron-deficient Cf species. For the production of light Cf isotopes we used fusion-evaporation reactions induced by  $^{34}\text{S}$  and  $^{36}\text{S}$  projectiles on  $^{206,207,208}\text{Pb}$  target nuclei. In general, the available reaction systems  $^{206,207,208}\text{Pb}+^{32,33,34,36}\text{S}$  provide not only the possibility of producing the lightest new Cf nuclei but also an opportunity to study the influence of nuclear structure and nucleon composition of both reaction partners on cross sections of fusion-evaporation reactions leading to very heavy, highly fissionable evaporation residues. At present, experimental information of this kind is very scarce or virtually absent. Yet this knowledge is of far-reaching importance for both understanding the mechanism of fusion of two complex nuclei and finding out the most prolific ways for the synthesis of new heavy and superheavy nuclides. An especially interesting case is associated with the neutron-rich  $^{36}\text{S}$  projectile which involves the magic number of neutrons,  $N=20$ , and can be considered to be similar, in a sense, to the famous projectile  $^{48}\text{Ca}$ . However, the fusion-evaporation reactions between Pb and S nuclei have never been studied before. The cross section measurements performed in our work represent a first step in this direction.

## 2. Experimental technique

The irradiations were performed at the Dubna U400 cyclotron by using beams of  $^{34,36}\text{S}$  and  $^{40}\text{Ar}$  projectiles with incident energies of 215 and 225 MeV, respectively. For the present studies, it was important to obtain appropriately intense beams of  $^{34}\text{S}$  and  $^{36}\text{S}$  at a reasonably low consumption of isotopically enriched sulphur materials (the abundances of  $^{34}\text{S}$  and  $^{36}\text{S}$  in naturally occurring sulphur are 4.21 and 0.02%, respectively). A special technique was developed for the production and acceleration of these ions by using a solid ( $\text{ZnS}$  enriched in  $^{34}\text{S}$  to  $\approx 40\%$  or in  $^{36}\text{S}$  to  $\approx 24\%$ ) as a working material to feed the discharge chamber of a PIG ion source of sputtering type. This technique allowed us to achieve an average consumption of the working material of  $\approx 10 \text{ mg h}^{-1} \mu\text{A}^{-1}$  and to provide a possibility of  $\approx 60\%$  recovering of the material. Average intensities of  $^{34}\text{S}$  and  $^{36}\text{S}$  beams applied to Pb targets were about  $8 \times 10^{12}$  and  $3 \times 10^{12}$  pps, respectively.

In our experiments we employed the wheel system described in ref.[3]. A beam of accelerated particles struck tangentially the lateral surface of a cooled copper cylinder onto which 2 to 3  $\text{mg/cm}^2$  of the metallic target material was deposited. This cylindrical target (serving simultaneously as a recoil catcher) rotated at a constant velocity. For each particular bombardment, the period of the wheel revolution,  $T_{\text{rev}}$ , was chosen according to the expected half-life value of a SF activity under study. Mica fission fragment detectors arranged all around the rotating target cylinder (except for the beam input zone) were used for the detection of SF events. The metallic layers of isotopically enriched  $^{208}\text{Pb}$  (99.0%),  $^{207}\text{Pb}$  (93.2%), and  $^{206}\text{Pb}$  (94.9%) were deposited onto the target cylinder by evaporation in vacuum.

Earlier, this technique was extensively used in experiments aimed at synthesizing transfermium nuclides (see, e.g., refs.[3,4]) where it permitted the detection of SF species produced with picobarn cross sections. More recently, it was employed in the experiments that led to the discovery of a new region of  $\text{EC}(\beta^+)$ -delayed fission around  $^{180}\text{Hg}$ - $^{188}\text{Pb}$ - $^{196}\text{Po}$  [5], in the studies of the stability of the K-isomeric states

of  $^{250}\text{Fm}$  and  $^{254}\text{102}$  against SF [6], in searches for carbon radioactivity of  $^{114}\text{Ba}$  [7], as well as in other studies [8].

### 3. Results and discussion

To probe the SF stability of the 3.4-min isotope  $^{242}\text{Cf}$  [1], we employed the  $^{208}\text{Pb}(^{40}\text{Ar}, 2n)$  reaction leading to the  $\approx 1.2$ -s  $\alpha$ -decaying nucleus  $^{246}\text{Fm}$  ( $b_\alpha=92\%$ ) and thus to  $^{242}\text{Cf}$ . Since  $^{246}\text{Fm}$  possesses a SF branch,  $b_{sf}=8\%$  [1], its SF detection could be used to calibrate the yield of  $^{246}\text{Fm}$  and hence that of  $^{242}\text{Cf}$ . Two  $^{208}\text{Pb}+^{40}\text{Ar}$  bombardments were performed at two different periods of the target wheel revolution,  $T_{rev}=10$  s and 17 min, which were chosen according to the  $T_{1/2}$  values of  $^{246}\text{Fm}$  and  $^{242}\text{Cf}$ . In the first bombardment with the  $^{40}\text{Ar}$  beam dose of  $3 \times 10^{17}$  we observed 840 SF events distributed in time with  $T_{1/2}=1.5 \pm 0.1$  s. From this, the  $^{246}\text{Fm}$  production cross section was determined to be  $19 \pm 9$  nb. The second bombardment with the beam dose of  $1.4 \times 10^{18}$  was performed to search for SF of  $^{242}\text{Cf}$ . Only 5 SF events were detected and thus the SF stability of  $^{242}\text{Cf}$  was proven to be rather high. It is characterized by  $b_{sf} \leq 1.4 \times 10^{-4}$  and  $T_{sf} \geq 1.5 \times 10^6$  s.

The SF stability of  $^{240}\text{Cf}$  was probed by using the  $^{208}\text{Pb}+^{34}\text{S}$  reaction. In a bombardment performed with the  $^{34}\text{S}$  beam dose of  $1.2 \times 10^{18}$  at  $T_{rev}=5$  min we revealed a fission activity (65 events) with  $T_{1/2}=0.9 \pm 0.2$  min, at a cross section level of 20 pb (see table 1 and fig.1). We assigned this activity to the SF branch of the  $\alpha$ -decaying nuclide  $^{240}\text{Cf}$  on the basis of the following observations and arguments. First, its  $T_{1/2}$  value agrees with  $T_{1/2}=1.06 \pm 0.15$  min known for  $\alpha$  decay of  $^{240}\text{Cf}$  [1]. Second, the cold-fusion-type reactions  $^{206,207,208}\text{Pb}+^{34}\text{S}$  specified by the so-called minimum excitation energy  $E_{min}^*$  of the composite systems (i.e., the excitation energy at the Bass fusion barrier [11]) of about 33.5 MeV are similar to the well-studied  $^{206,207,208}\text{Pb}+^{40}\text{Ar}$  reactions characterized by  $E_{min}^* \approx 31.5$  MeV. Hence, 2n to 4n evaporation channels are expected to be the most probable ones in the  $^{208}\text{Pb}+^{34}\text{S}$  system, as it is the case in the reactions  $^{208}\text{Pb}(^{40}\text{Ar}, xn)$  (for a summary of measured cross sections of the  $^{206,207,208}\text{Pb}+^{40}\text{Ar}$  fusion-evaporation reactions, see ref.[10]). We note,

**Table 1**

Summary of experimental results on the production of SF activities in bombardments of  $^{206,207,208}\text{Pb}$  target nuclei with  $^{34}\text{S}$  and  $^{36}\text{S}$  projectiles

Reaction	$T_{\text{rev}}$	Beam dose $\times 10^{17}$	$N_{sf}^{a)}$	$T_{1/2}^{b)}$	Assignment	xn channel	$\sigma_{sf}^{c)}$ nb
$^{208}\text{Pb}+^{34}\text{S}$	5 min	12	65	$0.9\pm 0.2$ min	$^{240}\text{Cf}$	2n	$0.02\pm 0.01$
$^{206}\text{Pb}+^{36}\text{S}$	5 min	1.5	38	$0.8_{-0.2}^{+0.3}$ min	$^{240}\text{Cf}$	2n	$0.10\pm 0.05$
$^{208}\text{Pb}+^{34}\text{S}$	0.2 s	3	387	$20\pm 2$ ms	$^{238}\text{Cf}$	4n	$0.5\pm 0.2$
$^{207}\text{Pb}+^{34}\text{S}$	0.2 s	1.5	425	$26\pm 4$ ms	$^{238}\text{Cf}$	3n	$1.1\pm 0.5$
$^{206}\text{Pb}+^{34}\text{S}$	0.2 s	3	244	$25_{-6}^{+9}$ ms	$^{238}\text{Cf}$	2n	$0.3\pm 0.1$
$^{206}\text{Pb}+^{36}\text{S}$	0.2 s	0.5	100	$15_{-3}^{+4}$ ms	$^{238}\text{Cf}$	4n	$0.7\pm 0.3$
$^{207}\text{Pb}+^{34}\text{S}$	15 s	4	63	$2.4_{-0.4}^{+0.8}$ s	$^{237}\text{Cf}$	4n	$0.05\pm 0.02$
$^{206}\text{Pb}+^{34}\text{S}$	15 s	10	121	$1.9\pm 0.3$ s	$^{237}\text{Cf}$	3n	$0.05\pm 0.02$

a) Total number of detected SF events.

b) Deduced by using the maximum likelihood method [9]; the indicated errors of  $T_{1/2}$  reflect statistical uncertainties only. In calculating  $T_{1/2}$  values for  $^{238}\text{Cf}$ , small contributions from longer-lived SF activities of  $^{240}\text{Cf}$  and  $^{237}\text{Cf}$  were taken into account.

c) Cross sections for SF branches. The  $\sigma_{sf}$  values were estimated from thick-target yields assuming  $\Delta E_{FWHM}=9\pm 2$ ,  $10\pm 2$ , and  $12\pm 2$  MeV for the widths of the excitation functions of the 2n-, 3n-, and 4n-evaporation channels, respectively (see, e.g., refs.[6,8,10] and references cited therein).

In the case of  $^{238}\text{Cf}$  ( $b_{sf}\approx 1$ ), the  $\sigma_{sf}$  values give total production cross sections.

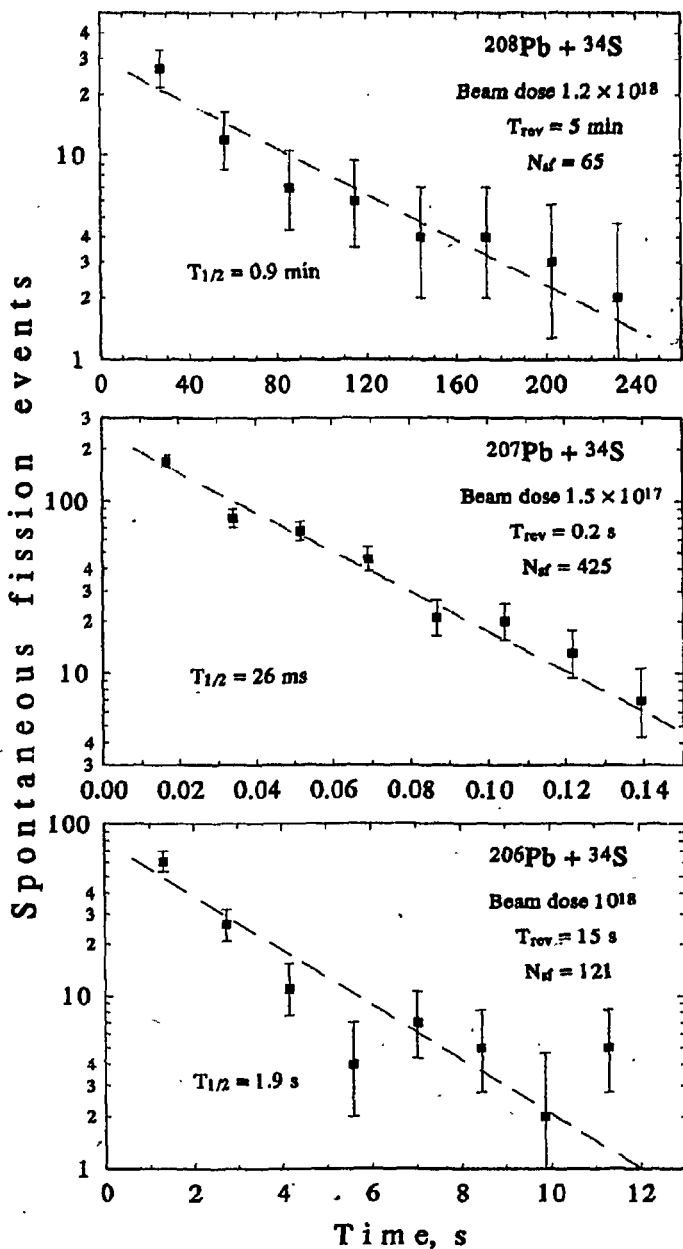


Fig.1. Time distributions of SF events detected in some particular bombardments  $^{208}\text{Pb}+^{34}\text{S}$ ,  $^{207}\text{Pb}+^{34}\text{S}$ , and  $^{206}\text{Pb}+^{34}\text{S}$ . See also table 1.

however, that the  $^{34}\text{S}$ -induced reactions on targets of  $^{206,207,208}\text{Pb}$  lead to Cf nuclides which are more neutron-deficient ( $N=138-142$ ) compared to Fm nuclides ( $N=142-146$ ) produced by the  $^{40}\text{Ar}$ -induced reactions. This fact should be reflected by a corresponding decrease in cross sections for xn channels in the  $^{34}\text{S}$  case. Finally, the odd-A nucleus  $^{239}\text{Cf}$  as well as the isotope  $^{238}\text{Cf}$  (see below) cannot be sources of the 0.9-min fission activity, whereas a possible contribution from EC( $\beta^+$ )-delayed fission with  $T_{1/2}=2.4\pm 0.1$  min in the decay chain  $^{238}\text{Bk} \xrightarrow{\text{EC}} ^{238}\text{Cm}$  [8,12] is expected to be small. Assuming for the  $^{208}\text{Pb}(^{34}\text{S}, 2n)$  reaction a cross section of  $\sim 1$  nb, we obtain an order-of-magnitude estimate  $b_{sf} \sim 2 \times 10^{-2}$  for  $^{240}\text{Cf}$ . Evidently, a more accurate  $b_{sf}$  determination for  $^{240}\text{Cf}$  requires an absolute measurement of the  $^{208}\text{Pb}(^{34}\text{S}, 2n)$  reaction cross section.

The observation of the SF decay mode of  $^{240}\text{Cf}$  was confirmed by the results of a bombardment of  $^{206}\text{Pb}$  with  $^{36}\text{S}$  that we performed at  $T_{rev}=5$  min. As table 1 shows, some 40 SF events distributed in time with  $T_{1/2}=0.8_{-0.2}^{+0.3}$  min were detected in this bombardment. This result provides also a possibility of making the interesting comparison between the 2n-evaporation cross sections of the complete fusion reactions  $^{206}\text{Pb}+^{36}\text{S}$  and  $^{208}\text{Pb}+^{34}\text{S}$  with  $E_{min}^*=29.8$  and 33.5 MeV, respectively, leading to the same compound nucleus  $^{242}\text{Cf}$ . For the  $^{36}\text{S}$ -induced reaction, the measured 2n-evaporation cross section proved to be  $\approx 5$  times larger as compared to the  $^{34}\text{S}$  case. This fact seems to be due to the lowered  $E_{min}^*$  value in the former reaction, which is somewhat more appropriate for the sub-barrier 2n-evaporation channel.

To explore the SF stability of still lighter, unknown isotopes of Cf, we carried out two  $^{207}\text{Pb}+^{34}\text{S}$  bombardments with the wheel revolution periods  $T_{rev}=0.2$  and 15 s. With  $T_{rev}=0.2$  s, a striking short-lived fission activity was discovered—we detected 425 fission events distributed in time according to  $T_{1/2}=26$  ms (see fig.1 and table 1). The yield of these events corresponds to a cross section of about 1 nb. With  $T_{rev}=15$  s, we observed 63 fission events distributed in time with  $T_{1/2}\approx 2$  s (table 1); these events appeared with a cross section of  $\approx 50$  pb.

Further, we made also two  $^{206}\text{Pb}+^{34}\text{S}$  bombardments, again with  $T_{rev}=0.2$  and 15 s. As table 1 and fig.1 show, the pattern of fission events observed was qualitatively similar to that of the  $^{207}\text{Pb}+^{34}\text{S}$  case.

Thus, we revealed two new fission activities. It is quite clear that neither  $\text{EC}(\beta^+)$ -delayed fission of neutron-deficient Bk or Am nuclides nor SF of light Cm or Pu nuclides can provide a source of the very short-lived fission activity with  $T_{1/2}\approx 25$  ms. As can be inferred from systematics [13], the appearance of a new spontaneously fissioning isomer with such a half-life value is also absolutely improbable in the region of neutron-deficient actinide nuclei. From these considerations and from the data presented in table 1 it follows that the short-lived SF activity should belong to the new Cf isotope,  $^{238}\text{Cf}$ , produced via the  $3n$  and  $2n$  evaporation channels on  $^{207}\text{Pb}$  and  $^{208}\text{Pb}$ , respectively. To verify this assignment once again, we performed a bombardment of  $^{206}\text{Pb}+^{34}\text{S}$  at  $T_{rev}=0.2$  s. Indeed, we again detected the 20-ms SF activity (387 events) with a cross section of  $\approx 0.5$  nb.

Finally, we clearly identified the  $\approx 20$ -ms SF activity in a bombardment of  $^{206}\text{Pb}$  with  $^{36}\text{S}$  carried out at  $T_{rev}=0.2$  s. This allows us to compare straightforwardly the cross sections of the fusion-evaporation reactions  $^{206}\text{Pb}(^{36}\text{S}, 4n)$  and  $^{208}\text{Pb}(^{34}\text{S}, 4n)$  leading to  $^{238}\text{Cf}$ . As seen from table 1, these cross sections proved to be virtually equal.

From predictions [14,15] and experimental systematics shown in fig.2,  $^{238}\text{Cf}$  is expected to have a partial  $\alpha$ -decay half-life of a few seconds. Therefore we conclude that the isotope  $^{238}\text{Cf}$  gives a new example of a short-lived spontaneously fissioning nucleus. By analyzing the whole set of data from the four reactions employed to produce  $^{238}\text{Cf}$  (see table 1), we determined its half-life to be  $21\pm 2$  ms.

Although some additional bombardments would be desirable to exclude completely few minor SF or  $\text{EC}(\beta^+)$ -delayed fission sources (e.g.,  $^{236}\text{Bk}$ ), the most probable origin of the SF activity with  $T_{1/2}=2.1\pm 0.3$  s seems to be due to a perceptible SF branch of the new, odd-A isotope  $^{237}\text{Cf}$  which is expected [14,15] to be predominantly an  $\alpha$  emitter. In this case, the evaluation of the SF branch is essentially



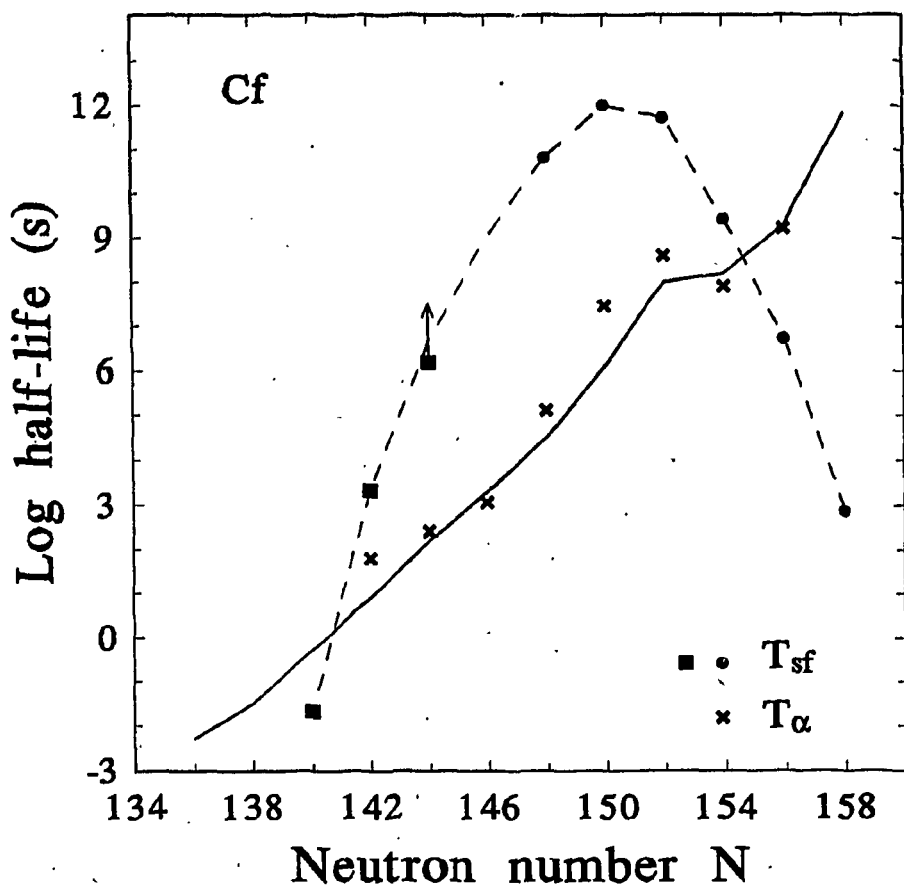


Fig.2. Experimental partial half-lives for  $\alpha$  and SF decay of the even-even Cf isotopes. The squares show the data of the present work. Solid line shows partial  $\alpha$  half-lives calculated by Smolanczuk and Sobiczewski [15].

dependent on the assumptions regarding the absolute cross sections of the  $^{207}\text{Pb}(^{34}\text{S}, 4n)$  or  $^{206}\text{Pb}(^{34}\text{S}, 3n)$  reactions. Using a cross section value of 0.5 nb, we obtain an order-of-magnitude estimate  $b_{sf} \sim 10^{-1}$  and, accordingly, a SF hindrance factor of  $\geq 10^3$ .

The decay properties of the light isotopes of Cf studied in the present work are summarized in table 2.

**Table 2**

Decay properties of the light Cf isotopes studied in the present work

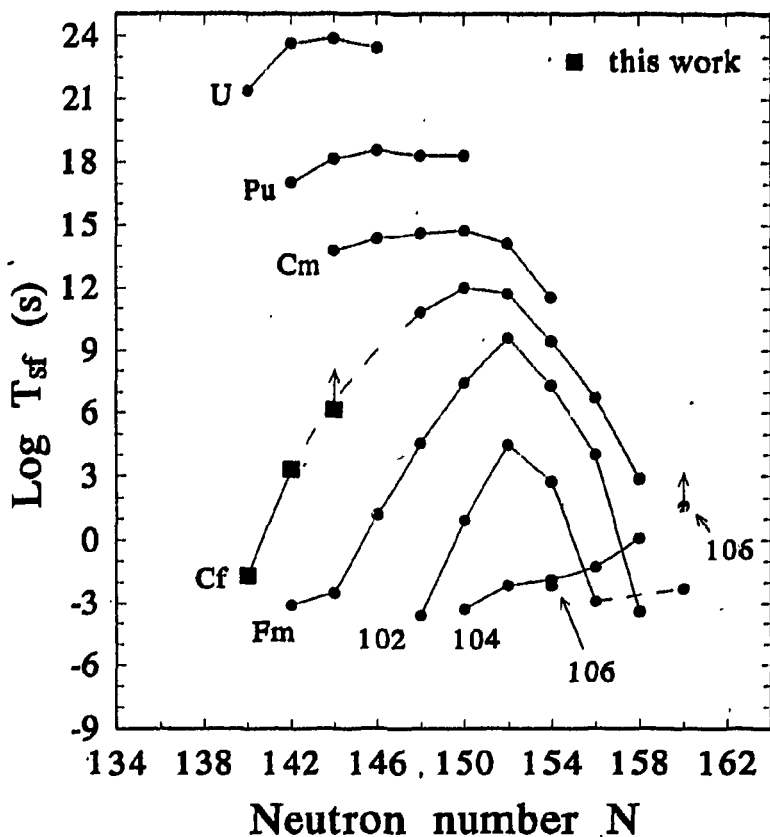
Isotope	Half-life	$b_{sf}$	$T_{sf}$
$^{242}\text{Cf}$	$3.4 \pm 0.2 \text{ min}^{a)}$	$\leq 1.4 \times 10^{-4}$	$\geq 1.5 \times 10^6 \text{ s}$
$^{240}\text{Cf}$	$1.06 \pm 0.15 \text{ min}^{a)}$	$\sim 2 \times 10^{-2}$	$\sim 3 \times 10^3 \text{ s}$
$^{238}\text{Cf}$	$21 \pm 2 \text{ ms}$	$\approx 1$	$\approx 21 \text{ ms}$
$^{237}\text{Cf}$	$2.1 \pm 0.3 \text{ s}$	$\sim 10^{-1}$	$\sim 20 \text{ s}$

<sup>a)</sup> Data from ref.[1].

#### 4. Conclusion

Thus, we explored the SF stability of the light Cf nuclei in a wide range of N. The dramatic effect of the neutron-deformed shell N=152 on the SF half-lives was demonstrated by revealing a  $T_{sf}$  decrease from  $6 \times 10^{10} \text{ s}$  for  $^{246}\text{Cf}$  down to  $2 \times 10^{-2} \text{ s}$  for  $^{238}\text{Cf}$  (see fig.3). Our experiments resulted in the production of the new nuclide  $^{238}\text{Cf}$  and gave an indication of the production of the new isotope  $^{237}\text{Cf}$ . The identification of  $^{238}\text{Cf}$  was confirmed in recent experiments [18] performed by using the Dubna gas-filled recoil separator [19].

We obtained also first experimental information about cross sections  $\sigma_{xn}$  of  $^{34}\text{S}$ - and  $^{36}\text{S}$ -induced fusion-evaporation reactions occurring on  $^{206,207,208}\text{Pb}$  target nuclei. The measured  $\sigma_{xn}$  values were found to be in the range of 0.3 to 1.1 nb for  $x=2$  to 4. A comparative study of the  $^{206}\text{Pb}(^{36}\text{S}, xn)$  and  $^{208}\text{Pb}(^{34}\text{S}, xn)$  reactions with  $x=2,4$  leading to the same compound nuclei and final products has shown that the  $\sigma_{2n}$  value is some 5 times larger in the  $^{36}\text{S}$  case, while  $\sigma_{4n}$  values are practically equal. An interesting extension of the present studies will be associated with the involvement of the  $^{32}\text{S}$  projectile which makes it possible also to attempt the further production of new, ultra-neutron-deficient Cf species.



**Fig.3.** Systematics of partial SF half-lives for even-even nuclei with  $Z=92$  through 106. For origins of the data points, see refs.[1,6,16,17] and references therein.

We wish to thank the U400 cyclotron staff led by B.N.Gikal for providing these experiments with the intense  $^{34}\text{S}$ ,  $^{36}\text{S}$  and  $^{40}\text{Ar}$  beams. We are indebted to the late V.M.Plotko for preparing excellent lead targets. Our thanks are due to L.V.Jolos and K.I.Merkina for their job of processing and scanning fission fragment detectors. We take pleasure in thanking R.Smolanczuk and A.Sobiczewski for performing  $\alpha$  half-life calculations for ultra-neutron deficient Cf isotopes and for their kind permission to include these results in the present paper.

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Received by Publishing Department  
on August 8, 1994