

ОБЪЕДИНЕННЫЙ  
ИНСТИТУТ  
ЯДЕРНЫХ  
ИССЛЕДОВАНИЙ

Дубна

96-160

E6-96-160

V.G.Chumin, V.I.Fominykh, K.Ya.Gromov,  
M.Ya.Kuznetsova, V.V.Tsupko-Sitnikov, M.B.Yuldashev

LEVEL STRUCTURE IN  $^{213}\text{Bi}$

Submitted to «Zeitschrift für Physik A»

1996

In our studies of  $(\alpha - \gamma)$ -coincidences at the decay of  $^{225}\text{Ac}$  and its daughter products [1,2,3] we had a possibility of obtaining some experimental data on  $\alpha$ -decay of  $^{217}\text{At}$  to levels of  $^{213}\text{Bi}_{130}$  which is expected to be a good example of the nucleus described by the shell model.

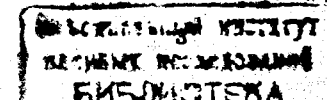
The available data on the  $^{213}\text{Bi}$  levels are scanty and somewhat contradictory [4]. On the basis of the results of Liang [5] and Dzhelepov et al. [6] levels of  $^{213}\text{Bi}$  with energies 258 keV, 465 keV and 593 keV were introduced. A few more  $\gamma$ -transitions were ascribed to  $^{217}\text{At}$  decay [4] but not placed in the decay scheme. Last year we established that the 6810 keV  $\alpha$ -line earlier ascribed to the  $^{217}\text{At}$  decay belonged to  $^{221}\text{Ra}$  arising in the  $^{225}\text{Ac}$  decay chain in the weak  $\beta$ -decay branch of the  $^{221}\text{Fr}$  decay [3] and therefore the 465 keV level does not exist in  $^{213}\text{Bi}$ .

We separated  $^{225}\text{Ac}$  ( $T_{1/2}=10$  days) [7] from 30  $\mu\text{Ci}$  activity of  $^{229}\text{Th}$  ( $T_{1/2} = 7340$  years). The samples obtained were vacuum evaporated on a thin Al foil. The  $^{225}\text{Ac}$  sources were placed between the  $\alpha$  and  $\gamma$  detectors in the  $180^\circ$  close geometry. The Si(Au)- $\alpha$ -detector,  $100\text{ mm}^2$  in area, has a 20 keV resolution. The HPGe  $\gamma$ -detector, 84 cc in volume, has a 1.0 keV resolution at the energy 150 keV. Single  $\alpha$  and  $\gamma$  spectra were taken together with  $4096 \times 4096$  channel coincidence measurements [8]. The  $(E_\alpha, E_\gamma, T)$ -coincidence events were written in a list mode. In the course of the analysis of the experimental data on very weak ( $\sim 0.01\%$  per decay) branches of the  $^{217}\text{At}$   $\alpha$ -decay it was necessary to take into account relatively weak instrumental effects caused by  $^{225}\text{Ac}$ ,  $^{221}\text{Fr}$  and  $^{213}\text{Bi}$  nuclei from the  $^{225}\text{Ac}$  decay chain the decay of which is accompanied by intense  $\alpha$  and  $\gamma$  radiation. These effects are: a) random summation of the amplitudes of the  $\alpha$ -pulses resulting in coincidences of intense  $\gamma$ -peaks throughout the  $\alpha$ -spectrum (for example weak coincidences with  $\gamma_{218\text{ keV } ^{221}\text{Fr}}$  are observed up to  $E_\alpha > 10$  MeV); b) summing of the pulses from the  $\alpha$ -particles and conversion electrons; c) summation of the pulses from cascade  $\gamma$  and  $\chi$ -rays in  $\gamma$ -spectra; d) remnants of the random coincidence peaks. The majority of the random events is eliminated in the process of sorting the coincidence events. We identified  $\gamma$ -rays with  $^{217}\text{At}$  decay by observing the changes in the area of the  $\gamma$ -peaks in the set  $\alpha$ -gates (Fig. 1a). For example maxima of the  $\gamma_{258\text{ keV}}$  peak area are observed at the energies  $\sim 6810$  keV and  $\sim 6480$  keV. The maxima of the

area of the 335 keV and 593 keV  $\gamma$ -peaks manifest themselves only at the energy  $\sim 6480$  keV. Thus the  $^{213}\text{Bi}$  258 keV and 593 keV levels (Fig.1b) proposed in [5,6] are confirmed. Besides, we observed 759 keV  $\gamma$ -rays with the area maximum in the  $E_\alpha=6340$  keV gate. This gate energy is the energy of the intense  $\alpha_0$ -line at the  $^{221}\text{Fr}$  decay ( $E_{\alpha_0} = 6343\text{ keV}$ ,  $I_{\alpha_0} = 85\%$ )

but we cannot attribute coincidences with  $\gamma_{759\text{ keV}}$  to the remnant of the random coincidences because in the  $^{217}\text{At}$   $E_{\alpha_0}=7069$  keV gate this peak is absent (Fig. 1.). The energy of  $\alpha$  particles calculated for the population of the  $^{213}\text{Bi}$  759 keV level is equal to 6322 keV. This allowed us to confirm excitation of the  $^{213}\text{Bi}$  759 keV level in the  $^{217}\text{At}$   $\alpha$ -decay. In the  $\gamma$  spectra measured in the gates of Fig.1 we observed a few more weak  $\gamma$  peaks. The analysis similar to the one in Fig.1 excludes their belonging to the  $^{217}\text{At}$  decay. For example, the 150 keV  $\gamma$ -peak area showed the maximum at  $E_\alpha=6610$  keV. These  $\alpha$ -particles, as mentioned above, arise from the  $^{221}\text{Ra}$  decay after the  $^{221}\text{Fr}$   $\beta^-$ -decay. The 440 keV  $^{213}\text{Bi}$  and 465 keV  $^{209}\text{Tl}$   $\gamma$ -peak areas have maxima in the  $^{221}\text{Fr}$   $E_{\alpha_0}$  and  $^{217}\text{At}$   $E_{\alpha_0}$  gates and therefore are connected with the random  $(\alpha - \gamma)$ -coincidences. In the gated  $\gamma$ -spectra we also observed weak peaks of the most intense  $\gamma$ -rays of  $^{225}\text{Ac}$  and  $^{221}\text{Fr}$ : 218 keV, 100 keV, 172 keV, 188 keV, etc. We ascribe them to the effect of random summation of the  $\alpha$ -particle pulses.

For determination of the exact values of the  $\gamma$ -ray energy and intensity (Table 1) the single  $\gamma$ -spectrum of the  $^{225}\text{Ac}$  equilibrium chain and the spectrum of  $\gamma$ -rays coinciding with  $\alpha$ -particles of energy  $E_\alpha = 5.9 \div 7.2$  MeV were used. The determination of the intensity of  $\gamma$ -rays in per cent per decay was made using the known data on the 218.2 keV  $\gamma$ -transition at the decay of  $^{221}\text{Fr}$ . According to [9] the 218.2 keV  $\gamma$ -transition is of E2 type. Its total intensity is 15.1% per decay [5,9]. Thus, using the calculated [10] conversion coefficients, we found the intensity of 218.2 keV  $\gamma$ -rays to be 11.2(2)% per decay. The upper limits estimated for the  $\gamma$ -ray intensity of the possible unobserved 501 keV and 166 keV  $\gamma$ -transitions from the 759.0 keV level are given in Table 1. Relative intensities of 258 keV  $\gamma$ -rays and  $K\chi$ -rays in the spectrum gated by  $E_{\alpha_{258}}=6810$  keV were used to determine  $\alpha_{K258}=0.45(4)$ . Comparison of this value with the calculated ones [10] allows the conclusion that the 258 keV  $\gamma$ -transition is of the (75% M1 + 25% E2) type.



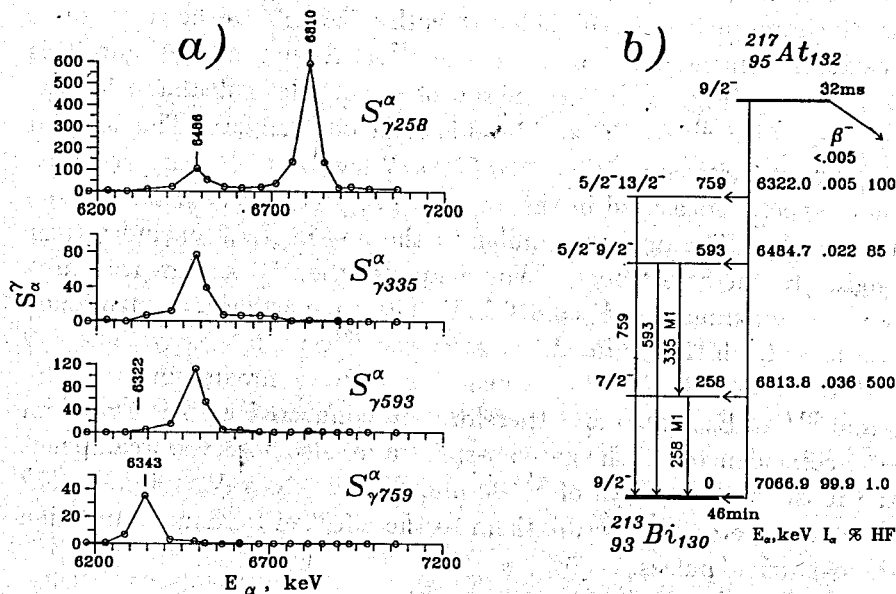


Fig. 1. a)  $(\alpha - \gamma)$ -coincidence at the decay of  $^{217}\text{At}$ . Gamma peak areas  $S_\gamma$  in the alpha gated spectra as a function of the gate middle energy  $E_\alpha$ . Statistical errors do not exceed the size of the points. b) Decay scheme of  $^{217}\text{At}$ .

Table 1. Alpha-decay of  $^{217}\text{At}$  (32 ms) on the levels of  $^{213}\text{Bi}_{130}$

Energy	(keV)	of	Intensity	(% per decay	$^{217}\text{At}$ ) of	HF
$\gamma$ -ray	level	$\alpha$ -particle	$\gamma$ -ray	$\gamma$ -transition	$\alpha$ -decay <sup>a)</sup>	
-	-	7066.9(16)	-	-	99.9	1.0
257.88(4)	257.88(4)	6813.8(16)	0.0287(7)	0.0450(25)	0.036(3)	500
335.33(10)	593.10(10)	6484.7(16)	0.0062(3)	0.0090(9)	0.021(2)	85
593.10(10)	" "	" "	0.0115(5)	0.012(1)		
758.9(1)	758.9(1)	6322.0(16)	0.0049(4)	0.005(1)	0.005(1)	100
(165.8)	" "	" "	<0.0002			
(501.0)	" "	" "	<0.0002			

<sup>a)</sup> Intensities of the  $\alpha$ -decay branches are calculated on the assumption of the absence of the level population from higher lying unknown levels.

The quantitative analysis of the  $(\alpha - \gamma)$ -coincidence spectra confirms the population of the  $^{213}\text{Bi}$  257.9 keV, 592.9 keV and 759.0 keV levels at the  $^{217}\text{At}$   $\alpha$ -decay. The total intensity of the 335 keV transition is determined to be 0.0090(9)% per decay. This value is consistent with M1 multipolarity of the 335 keV  $\gamma$ -transition. Weak population of the 592.9 keV and 759 keV levels by transitions from higher so far unknown levels of  $^{213}\text{Bi}$  is not excluded. The experimental data on the  $^{217}\text{At} \rightarrow ^{213}\text{Bi}$  decay (Fig.1b) are summed up in Table 1. Energies of  $\alpha$ -particles for the decay to excited  $^{213}\text{Bi}$  levels are calculated using the energy value for the ground-ground  $\alpha$ -transition  $E_{\alpha 0} = 7066.9(16)$  keV recommended by Rytz [11] and our values for the  $^{213}\text{Bi}$  level energies. The new fine structure line  $E_{\alpha 759} = 6322.0$  keV,  $I_{\alpha 759} = 0.005\%$  is discovered at the  $^{217}\text{At}$  decay. On the basis of more precise values for the  $\alpha$ -decay intensity hindrance factors for the  $\alpha$ -decay to  $^{213}\text{Bi}$  levels are determined. The experimental data allow the following interpretation of the excited states of  $^{213}\text{Bi}$ . The value of HF for the  $\alpha$ -transition between the  $^{217}\text{At}$  and  $^{213}\text{Bi}$  ground states shows that the following configurations must be ascribed to ground states of these nuclei:  $^{217}_{85}\text{At} - (\pi h_{9/2}^3, \nu g_{9/2}^6)_{9/2-}$  and  $^{213}_{83}\text{Bi} - (\pi h_{9/2}, \nu g_{9/2}^4)_{9/2-}$ . The high value HF=500 for the 258 keV level shows that this state has a configuration that includes shell model proton states  $f_{7/2}$  or  $i_{13/2}$ . The M1 character of the 258 keV transition excludes the latter possibility. Thus we ascribe the configuration  $(\pi f_{7/2}, \nu g_{9/2}^4)_{7/2-}$  to the 258 keV state. The majority of the wave functions of the 592.9 keV and 759 keV states is apparently connected with the coupling of the

