# ОБЬЕАИНЕННЫЙ ИНСТИТУт <br> ЯАЕРНЫХ <br> ИССАЕАОВАНИЙ 

AY5HA
R.arlt, G.Beyer, E.Herrmann, A.Jasinski, . O.Knotek, G.Musiol, H.-G.Ortlepp,
H.-U.Siebert, H.Tyrroff

THE DECAY SCHEME OF ${ }^{127} \mathrm{Ba}$
AND HALF-LIVES OF SOME LOW-LYING EXCITED STATES IN $123,125,127,129 \mathrm{Cs}$

1325

R.Arlt, ${ }^{\mathbf{1}}$ G.Beyer, ${ }^{2}$ E.Herrmann, A.Jasinski, ${ }^{\text {T }}$ O.Knotek, G.Musiol,' H.-G.Ortlepp, H.-U.Siebert, H.Tyrroff ${ }^{2}$

THE DECAY SCHEME OF $\quad{ }^{127}$ Ba
AND HALF-LIVES OF SOME LOW-LYING EXCITED STATES IN $123,125,127,129 \mathrm{Cs}$

Submitted to Nuclear Physics
${ }^{1}$ Technical University, Dresden, DDR.
${ }^{2}$ ZfK, Rossendorf, DDR.
${ }^{3}$ Address from July, 1976: Institute for Nuclear Research, Swierk k/Otwocka, Poland.


## 1. Introduction

The aim of the present study is to obtain further experimental information on the low-lying levels in ${ }^{123-129} \mathrm{Cs}$ excited in the beta decay of ${ }^{123-129} \mathrm{Ba}$. These neutrondeficient isotopes belong to the $50<\mathrm{Z}, \mathrm{N}<82$ region where a stable deformation has been predicted/1/. This new region of deformation has been confirmed by theoretical calculations $/ 2,3 /$ and by experimental data obtained from in-beam spectroscopy (even nuclei) ${ }^{/ 4-7 / /}$ ). The sign of this deformation is not yet well determined. The theoretical calculations in ref. $/ 3$ suggest that oblate shapes are more stable than prolate ones but the coexistence of both prolate and oblate states is also possible. Support for these predictions seems to be provided by the observation of the excited states populated in the decay of the $11 / 2^{-}$isomeric states observed $/ 8-10 /$ in odd ${ }^{127} \mathrm{Cs}$ and ${ }^{129}$ La isotopes. Tentative Nilsson orbital assignments for these states suggest that the $5 / 2^{+}$first excited state in ${ }^{127} \mathrm{Cs}$ can be treated/8/ as a shape isomer.

It is thus of special interest to extend our knowledge of these nuclei by studying low-spin states populated in beta decay processes. In the present paper the decay scheme of ${ }^{127} \mathrm{Ba}$ and the half-lives of some of the excited states in $123,125,127,129 \mathrm{Cs}$ are reported. It is suggested that in all four cases the states connected by the $E 2$ transitions ( $5 / 3^{+} \rightarrow 1 / 2^{+}$) have thesame deformation, and the possibility of the shape isomerism, occuring as discussed in ref. $/ 8 /$, is questionable. As regards the isomeric states in ${ }^{127} \mathrm{Ba}$, marked discrepancies exist between the data reported in ref. ${ }^{11 /}$ and those cited below.

Preliminary results of the present investigations have been published previously / 12-16/

## 2. Source Preparation

Suspensions of 3 g GeO 2 in 5 ml 0.02 M HCl solution (size of $\mathrm{CeO}_{2}$ grain was about $5 \mu$ ) were irradiated for 20 min by the external 660 MeV proton beam of the JINR synchrocyclotron. Because of the high recoil energy, $20-40 \%$ of the spallation reaction products are stabilized in the liquid phase and can be easily separated from the target material by filtration by analogy with the $\mathrm{Ta}_{2} \mathrm{O}_{5}$ suspension technique as described in ref. $/ 17$. After filtration the filtrate was stired during 60 sec with about 10 mg cation exchange resin (Dowex 50 x 5 , resin size $20 \mu$ or Aminex A5). The rare earths ( $\mathrm{Pr}, \mathrm{Ce}$ and La ), Ba and partially Cs were adsorbed on the resin, which was transformed intc the $\mathrm{NH}_{4}^{+}$-form. The rare earths including a few times 100 mg of a Ce carrier from the target material were eluted with 0.8 alpha-hydroxy-iso-butyric-acid (alpha-HIB). In order to elute a narrow Ba peak it was necessary to transform the $\mathrm{NH}_{4}^{+}$-cation exchanger into the $\mathrm{NH}_{4}^{+}$-form by washing the column with 1 M HCI. Under these conditions C s was eluted. Using $6 \mathrm{M} \mathrm{NHO}_{3}$ a very narrow peak of pure Ba was washed from the column (fig. 1). To transform the nitrates into the oxides the Ba fraction (1-2 drops $6 \mathrm{M} \mathrm{HNO}_{3}$ ) was evaporated until dryness on a $5 \times 5 \mathrm{~mm}^{2}$ Ta foil and then heated to about 800 K . The Ba -isotope separation was carried out at the YASNAPP-facility / 18 / using the pipe type surface ionization ion source ${ }^{19 /}$. The activity of the ${ }^{127} \mathrm{Ba}$ samples (measured usually 20 min after the end of irradiation) was of the order of 0.5 miCi . To investigate also the ${ }^{123} \mathrm{Ba}$ and ${ }^{125} \mathrm{Ba}$ isotopes, the samples of the Ba fraction were prepared during about 1 min from a 100 mg La target in $\mathrm{HNO}_{3}$ solution, irradiated for 20 sec. The Ba isotopes were collected as $\mathrm{BaSO}_{4}$ by carrier precipitation.

Samples were also prepared by the method of fast extraction of isobars by direct electromagnetic massseparation from proton ir radiated tantalum targets $/ 20 /$.
3. Singles Gamma-Ray Spectra Measurements of ${ }^{127} \mathrm{Ba}$

The singles gamma-ray spectra (figs. 2 and 3) of ${ }^{127} \mathrm{Ba}$ (mass-separated sources) were measured using $2,4 c c$ and $41 c c \mathrm{Ge}(\mathrm{Li})$ detectors with system resolutions of 0.7 keV at 122 keV and 2.4 keV at 1332 keV , respectively. The energies of the gamma-rays were determined by taking combination spectra of ${ }^{127} \mathrm{Ba},{ }^{133} \mathrm{Ba}$ and ${ }^{126}$ Ra sources. The results obtained from a few series of measurements are given in table 1. These results should be compared with the data of ref. ${ }^{11 / /}$ where two isomeric states in ${ }^{127} \mathrm{Ba}$ with half-lives of ( $10 \pm 1$ ) and $(18 \pm 1) \mathrm{min}$ were reported. The value of 10 min can be in agreement (in two standard er ror limits) with the halflife adopted here (see insert in fig. 2). The 18 min activity (in the 10-2500 keV gamma energy region) in this experiment was not observed.

## 4. Conversion-Electron Spectra Measurements of ${ }^{127} \mathrm{Ba}$

The measurements of the conversion-electron spectra (fig. 4) were carried out using a $\mathrm{Si}_{\mathrm{i}}(\mathrm{Li})$ detector ( $1.6 \mathrm{~cm} \mathbf{x}$ $\mathbf{x} 0.3 \mathrm{~cm}$ ) with a 2.5 keV energy resolution at 100 keV . To check the purity of the sources simultaneously with the conversion electron measurements the gamma-ray spectra were recorded with a 38 cc $\mathrm{Ge}(\mathrm{Li})$-detector. The final results obtained from three series of measurements are given in table 2.

## 5. Gamma-Gamma Coincidence Measurements of ${ }^{127} \mathrm{Ba}$

The gamma-gamma coincidence measurements were carried out using 41 and $45 c c \operatorname{Ge}(\mathrm{Li})$ detectors with energy resolutions of 2.4 and 2.9 keV , respectively, for the $1332 \mathrm{keV} \gamma$-ray of ${ }^{60} \mathrm{Co}$. The detectors were griented at $90^{\circ}$. Pulses from the coincidencesystem ${ }^{/ 22}$ ( 50 ns resolving time ) were used to open the gates of two 4096-channel ADC units connected to an HP 2116 C computer. The coincidence spectra obtained with the gates set on the more intense photopeaks are shown in fig. 5.

Table 1
Gamma-rays observed in the decay of ${ }^{123} \mathrm{Ba}$

| E $\boldsymbol{\gamma}$ | $I^{\prime}$ | $\mathrm{E}^{\boldsymbol{y}}$ | I $y$ |
| :---: | :---: | :---: | :---: |
| KX | $470 \pm 50$ | $1385.2 \pm 0.5$ | $0.8 \pm 0.2$ |
| $66.3 \pm 0.3$ | $17.1 \pm 1.7$ | $1437.5 \pm 1.0$ | $\sim 0.1$ |
| $72.8 \pm 0.5$ | $6.1 \pm 0.5$ | $1448.8 \pm 0.5$ | $0.4 \pm 0.1$ |
| $114.8 \pm 0.3$ | $74.7 \pm 3$ | $1500.1 \pm 0.3$ | $3.0 \pm 0.3$ |
| $139: 0 \pm 0.8$ | $0.8 \pm 0.4$ | $1511.2 \pm 1.0$ | $1.0 \pm 0.1$ |
| $180.8 \pm 0.3$ | 100 | $1522.0 \pm 0.7$ | $0.6 \pm 0.1$ |
| $429.3 \pm 0.6$ | $2.2 \pm 0.4$ | $1566.0 \pm 0.3$ | $3.1 \pm 0.3$ |
| $441.0 \pm 1.0$ | $0.4 \pm 0.2$ | $1576.3 \pm 1.0$ | $0.5 \pm 0.1$ |
| $451.5 \pm 1.0$ | $0.7 \pm 0.2$ | $1618.0 \pm 0.3$ | $2.0 \pm 0.3$ |
| 511.0 | $877 \pm 70$ | $1697.0 \pm 0.8$ | $0.4 \pm 0.2$ |
| $523.5 \pm 0.7$ | $3.5 \pm 0.8$ | $1753.6 \pm 0.3$ | $2.0 \pm 0.3$ |
| $532.1 \pm 0.7$ | $0.4 \pm 0.1$ | \% $800.1 \pm 0.6$ | $0.6 \pm 0.2$ |
| $567.5 \pm 0.3$ | $2.8 \pm 0.3$ | $1842.2 \pm 0.6$ | $0.9 \pm 0.2$ |
| $573.9 \pm 0.5$ | $0.7 \pm 0.2$ | $1915.3 \pm 0.6$ | $0.9 \pm 0.3$ |
| $578.0 \pm 0.3$ | $5.2 \pm 0.5$ | $1920.6 \pm 0.8$ | $0.3 \pm 0.1$ |
| $619.0 \pm 1.0$ | -0.1 | $1950.8 \pm 0.6$ | $1.1 \pm 0.2$ |
| $621.5 \pm 0.8$ | $0.2 \pm 0.1$ | $1 y 62.8 \pm 0.8$ | $0.5 \pm 0.1$ |
| $625.5 \pm 0.7$ | $0.4 \pm 0.1$ | $1981.8 \pm 0.3$ | $1.8 \pm 0.2$ |
| $647.1 \pm 0.8$ | $0.4 \pm 0.1$ | $1951.9 \pm 0.6$ | $1.0 \pm 0.1$ |
| $641.9 \pm 0.7$ | $0.4 \pm 0.1$ | $2028.2 \pm 0.7$ | $0.5 \pm 0.1$ |
| $713.5 \pm 0.8$ | $\sim 0.1$ | $2057.0 \pm 0.6$ | $1.0 \pm 0.2$ |
| $872.5 \pm 0.5$ | $0.8 \pm 0.1$ | $2075.0 \pm 0.6$ | $0.9 \pm 0.2$ |
| $1012.3 \pm 0.5$ | $0.4 \pm 0.1$ | $2089.8 \pm 0.4$ | $1.3 \pm 0.2$ |
| $1019.8 \pm 0.5$ | $1.3 \pm 0.1$ | $2100.3 \pm 0.5$ | $1.1 \pm 0.2$ |
| $1062.0 \pm 1.0$ | $0.4 \pm 0.1$ | $2141.0 \pm 0.8$ | $0.3 \pm 0.1$ |
| $1084.9 \pm 0.5$ | $3.5 \pm 0.3$ | $2172.0 \pm 0.6$ | $1.1 \pm 0.2$ |
| $1108.3 \pm 0.5$ | $0.9 \pm 0.2$ | $2182.0 \pm 0.3$ | $1.8 \pm 0.2$ |
| $1135.2 \pm 1.0$ | $\sim 0.1$ | $2189.0 \pm 0.7$ | $0.3 \pm 0.1$ |
| $1150.7 \pm 0.7$ | $1.5 \pm 0.2$ | $2222.4 \pm 0.7$ | $0.5 \pm 0.1$ |
| $1201.0 \pm 0.3$ | $13.0 \pm 1.5$ | $2238.0 \pm 1.0$ | $0.4 \pm 0.1$ |
| $1222.9 \pm 0.9$ | $0.2 \pm 0.1$ | $2321.2 \pm 0.5$ | $1.2 \pm 0.2$ |
| $1289.3 \pm 0.4$ | $1.0 \pm 0.1$ | $2467.8 \pm 0.7$ | $1.2 \pm 0.2$ |

$乙$ əTqe山




Fig. 1. Radio-chromatogram of the Ba -separation from the spallation products (RE and Cs) of a 660 MeV proton irradiated $\mathrm{CeO}_{2}$ target. Micro-cation exchange-column (Aminex A5; 2x15 mm).


Fig. 2. Low energy part of $\gamma$-radiation from mass-separated ${ }^{12}$ Ba sources observed with a 2.4 cc $\mathrm{Ge}(\mathrm{Li})$-detector. Decay curves
shown in the insert.




Fig. 5. Gamma-ray spectra in coincidence with the 66.3, 72.8 and 114.8 keV transitions. Gamma-spectra coincident with the Compton background are shown for comparison.

## 6. Beta-Gamma Coincidence Measurements of ${ }^{127} \mathrm{Ba}$

The beta-gamma coincidence measurements were carried out using $35 c c \quad \mathrm{Ge}(\mathrm{L} . \mathrm{i})$ and $\mathrm{NaI}(\mathrm{Tl})$ detectors as the gating crystals and a $1 \mathrm{~cm}^{2} \times 1.5 \mathrm{~cm} \mathrm{Si}(\mathrm{Li})$ detector with a resolution of 8 keV for the 975 keV conversion electron line of ${ }^{207} \mathrm{Bi}$ to obtain the coincidence beta-spectra. The detectors oriented at $180^{\circ}$ were coupled to the same coincidence system used for the gamma-gamma coincidence measurements (sect. 5). The continuous betaspectra distorted by the back-scattering and summing effects in the crystal were corrected using the method described in ref. $/ 23 /$. As standard sources ${ }^{138} \mathrm{P} \mathrm{Pr}$ and ${ }^{140} \mathrm{Pr}$ with well known end point energies ${ }^{/ 24,25 /}$ were applied. As an example, the beta-spectrum obtained in coincidence with the 180.8 keV gamma-ray is shown in fig. 6.

The $\beta^{+}$-ray spectra obtained in coincidence with the more intense gamma-lines were consistent and yielded the maximum energy for the betadecay of $3.45 \pm 0.10 \mathrm{MeV}$, which is in good agreement with the semi-empirical mass formula given in ref. ${ }^{/ 26 /}$. Special attempts to reveal the $\beta^{+}$component associated with the isomeric state /11/ have not been successful. This result together with the data of sect. 3 indicates that the 10 and 16 min activities reported in ref. /11/ in this experiment have not been observed.

## 7. Half-Life Measurements of Low-Lying Excited States in ${ }^{123,125,127,129} \mathrm{Cs}$

The half-lives of low-lying excited states in ${ }^{123-127} \mathrm{C}_{\mathrm{s}}$ isotopes were measured using a NE 102 (3.4 cm $\phi$ x $x 1.5 \mathrm{~cm}$ ) plastic scintillator (mounted on a FEU-36 photomultiplier tube) as the start detector and a 27 cc $\mathrm{Ge}(\mathrm{Li})$ detector to obtain the delayed coincidence gammaspectra. These spectra were measured using a slightly modified (one of the ADC units was fed by the output signals from a time-to-pulse-height convertor) two para-

meter analysis equipment used earlier in the $\gamma-\gamma$ experiments (sect. 5). The time resolution of the coincidence circuit was 8 ns FWHM with a gate set at 100 keV in the stop detector.The results of the $\beta-\gamma$ coincidence experiments are shown in fig 7.

In the case of ${ }^{129} \mathrm{C}_{\mathrm{s}}$ aNE102( $1 \mathrm{~cm} \phi \times 0.1 \mathrm{~cm}$ ) plastic scintillator (mounted on a EMI9524S photomultiplier tube) was used as a stop detector and the 27 cc $\mathrm{Ge}(\mathrm{Li})$ detector as a start one. The ${ }^{129} \mathrm{Ba}$ sources used in this experiment were prepared by implantation of the barium ions into the scintillator material in the electromagnetic mass-separator. The prompt and delayed $\gamma$ ray spectra gated by the $6.54 \mathrm{M}+\mathrm{N} \ldots$ conversion electron line and $L X$ radiation recorded in the plastic scintillator are shown in fig. 8.

The final results of these two experiments are given in table 3.
8. The Decay Scheme of $\quad{ }^{127} \mathrm{Ba}$

In the $\beta-\gamma$ coincidence experiment only one gamma transition of 66.3 keV was observed in the delayed coincidence gamma-ray spectra (fig. 7). This implies the existence of an excited state of the same energy. Furthermore, the $\gamma-\gamma$ coincidence measurements (fig. 5) clearly indicate that the 66.3 keV transition is in coincidence with the $\gamma$-rays of 72.8 and 114.8 keV . Since the 114.8 keV $\gamma$-transition is not observed in coincidence with the 72.8 keV one, it follows that there are two levels at 139.0 and 181.0 keV . These levels are supported by the cross-over transition of 139.0 and 180.8 keV observed in the single gamma-ray spectra. The three well established levels at $66.3,139.1$ and 181.0 keV served as a main basis for the further construction of the decay scheme shown in fig. 9.

The spin and parity of the gound state of ${ }^{127} \mathrm{Cs}$ has been found ${ }^{31-33}$ / to be $1 / 2^{++}$. By using the intensity balance of gamma transitions depopulating the $11 / \mathbf{2}^{-1}$ isomeric state in ${ }^{127} \mathrm{Cs}$ the spins and parities of the first and second excited states were estimated in ref. $/ 8 /$ to be

$$
\varepsilon \text { əтqе山 }
$$

Half-lives of some excited states in odd Cs isotopes

| Isotope | ${ }^{123} \mathrm{Cs}$ | ${ }^{125} \mathrm{Cs}$ | ${ }^{127} \mathrm{Cs}$ | ${ }^{129} \mathrm{Cs}$ |
| :--- | :---: | :---: | :---: | :---: |
| $\mathrm{E}_{\gamma}(\mathrm{keV})$ | 94.5 | 85.4 | 66.3 | 6.54 |
| $\mathrm{~T}_{1 / 2}(\mathrm{~ns})$ | $9 \pm 3$ | $14.5 \pm 1.5$ | $24.5 \pm 3$ | $72 \pm 6$ |



Fig. 7. Prompt and delayed coincidence $\gamma$-ray spectra (obtained by adding the counts of 26 measurements) of the short lived Ba fraction gated by $\beta^{+}$radiation recorded in the plastic scintillator with an integral gate set above 500 keV.


Fig. 8. Prompt and delayed coincidence $\gamma$-ray spectra of ${ }^{129} \mathrm{Ba}$ obtained with a gating window (A) set on the $6.54 \mathrm{keV} \mathrm{M}+\mathrm{N} . .$. conversion electron line and $L X-$
radiation. In the $\gamma$-ray spectra gated by $K X$ radiation (window B) no $\gamma$-lines were observed in the delayed spectra.

$5 / 2^{+}$and $3 / 2^{+}$respectively. These values are compatible with our experimental data.

The high intensity of the beta component to the ground state ( $\log \mathrm{ft}=5.2$ ) and no population (in experimental error limits) of the $66.3 \mathrm{keV} 5 / 2^{+}$state in the beta decay made it possible to conclude that the ground state spin value of ${ }^{127} \mathrm{Ba}$ is $1 / 2^{+}$. Taking into account the beta branch $(\log \mathrm{ft}=5.4)$ to the 181 keV level and the multipolarities of the 114.7 and 180.8 keV gamma transitions (table 2) the spin and parity of this state are likely to be $3 / 2^{+}$. The $1 / 2$ or $3 / 2$ spin values for higher levels have been postulated on the basis of the systematics ${ }^{/ 34 /}$ of the $\log \mathrm{ft}$ values in the neighbouring nuclei.

## 9. Discussion

Figure 10 shows the systematics of some low-lying levels of odd-mass Cs isotopes $/ 27-30$ / relative to the lowest $5 / 2^{+}$state. As it is seen from this figure the first $1 / 2^{+}$state is systematically lowered with decreasing mass number A , and starting from $\mathrm{A}=129$ it becomes the ground state of the lighter Cs isotopes. The same tendency can be observed for the low-lying $3 / 2^{+}$states. The general trend mentioned above leads to a level sequence in the light Cs isotopes the properties of which cannot be described by the single particle shell model. To explain the experimental data of the ground state of ${ }^{125} \mathrm{Cs}$ and ${ }^{127} \mathrm{Cs}$ in the framework of the intermediate coupling version of the unified model $/ 35,36 /$ adjustable parameters indicating deformable cores have to be used. This seems to suggest that the studied nuclei should be treated as deformed ones in agreement with the theoretical calculations of refs $/ 2,3 /$.

In the decay of ${ }^{12}{ }^{7} \mathrm{Ba}$ the two strong allowed $\beta^{+}$-transitions to the ground and third excited states are observed (fig. 9). Assuming that these two levels belong to the same rotational band with $K=1 / 2$ the theoretical ratio /37/ of ft -values for these transitions was estimated (about 2) to be in agreement with the experimental one. An ana-


Fig. 10. Systemstics of the low-lying levels of odd-mass $\mathrm{Cis}_{\mathrm{s}}$ isotopes $/ 27-30 /$ relative to the first $5 / 2^{+}$states.
lysis of the level energy, intensity and $B$ (E2) transition probability (table 4) led the authors to the conclusion that the $66.3 \mathrm{keV} 5 / 2^{+}$state cannot betreated as a further member of the rotational family of interest. Thus, a rotational nature can be assigned only to the $181 \mathrm{keV} 3 / 2^{+}$ state. This conclusion seems to be supported by the small amount of Ml radiation in the 181 keV gamma-transition (table 2) and the energy consistency with the first $2^{+}$ excited state in the equivalent even-even nucleus ${ }^{126} \mathrm{Xe}^{/ 38 /}$ (assuming the decoupling parameter to be small).

According to the quasi-particle Nilsson model ${ }^{1 / 3 /}$ the $1 / 2^{-1}[420\}$ orbit (for $0>0$ ) can be assigned to the unpaired proton of the lighter Cs isotopes. To explain the nature of the $11 / 2^{-}$isomeric state in ${ }^{127} \mathrm{Cs}$, it has been suggested $/ 8 /$ that this state and two others $\left(7 / 2^{+}\right.$at 272.5 keV and $5 / 2^{+}$at 66.3 keV ) via which the isomeric level is depopulated by gamma-transitions, should be described by the Nilsson orbitals $11 / 2^{-}[505], 7 / 2^{+}[413]$ and $5 / 2^{+}[413]$ for the oblate shape. Hence, the $5 / 2^{+}$first excited level would be a candidate of the intriguing shape isomerism. According to the theoretical calculations and the systematics given in fig. 10 one could expect that the $5 / 2^{+}$states in the neighbouring odd $C_{5}$. nuclei should reveal similar properties. The experimental B(E2) values confirmed the same nature of these levels but instead of the expected retardation an enhancement has been observed (table 4). Furthermore, according to the data of refs. ${ }^{/ 8,9 /}$ the 385.5 keV E3 transition, connecting the $11 / 2^{-}$and $5 / 2^{+}$states in ${ }^{127} \mathrm{Cs}$, has the theoretical single particle speed. These facts seem to indicate that the ground and all excited states (including the $11 / 2^{-}$level) in ${ }^{27}{ }^{7} \mathrm{C}$ s observed in ref. $/ 8 /$ have the same sign of deformation.

To summarize, the experimental data obtained in the present work did not reveal any consistency with the shape isomerism predictions ${ }^{3 /}$. The same deformation sign is to be assigned to ${ }^{127} \mathrm{Cs}$ and to the ground state of ${ }^{127} \mathrm{Ba}$. The ground state of ${ }^{127} \mathrm{Cs}$ can be better described by the prolate shape of the nucleus than by the oblate one. The high energy part of the excited states observed in
$\mp$ ƏТqе山


| Isotope | Transition | $\mathrm{E}_{\boldsymbol{y}}(\mathrm{keV})$ | $B(E 2)$ | $\begin{aligned} & \left(e^{2} \times 10\right. \\ & \exp \end{aligned}$ | $\frac{B(E 2)}{B(E 2)}$ | $\frac{\exp }{s \cdot p .}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{123} C_{8}$ | $\left(5 / 2^{+}\right) \cdots\left(1 / 2^{+}\right)$ | $94 \cdot 5$ | 0.25 | $\pm 0.08$ | . $5 \pm$ | - 2 |
| ${ }^{125} \mathrm{C8}$ | $5 / 2^{+} \rightarrow 1 / 2^{+}$ | 85.4 | 0.20 | $\pm 0.02$ | $1 \pm$ | 1.8 |
| ${ }^{127} \mathrm{Cs}$ | $5 / 2^{+} \rightarrow 1 / 2^{+}$ | 66.3 | 0.18 | $\pm 0.02$ | $.9 \pm$ | . 9 |
| ${ }^{129} \mathrm{Cs}$ | $5 / 2^{+} \rightarrow 1 / 2^{+}$ | 6.54 | 0.16 | $\pm 0.02$ | . $4 \pm$ | 1-2 |
| ${ }^{131_{\mathrm{Cs}}}$ | $1 / 2^{+} \rightarrow 5 / 2^{+}$ | 123.7 | 0.30 | $\pm 0.019)$ | $2 \pm$ | . 2 |



Fig. 11. Comparison of the experiment and theoretical intrinsic levels calculated $/ 39 /$ using the Saxon-Woods potential for equilibrium deformation of ${ }^{127} \mathrm{Cs}$.
${ }^{127} \mathrm{Cs}$ cannot, however, be explained in the framework of the simple Nilsson model because the number of $1 / 2,3 / 2$ states observed in the experiment exceeds the number of expected ones, as can be seen in fig. 11 where the intrinsic states in ${ }^{127} \mathrm{Cs}$ calculated with Saxon-Woods potential $/ 39$ / are compared with the experiment.

It is worth noting that the unified model in the intermediate coupling version/35,36/ and the pairing + quadrupole forces model $/ 40 /$ are also able to explain the experimental data of the ground states of odd Cs isotopes. The energy systematics of the low-lying levels shown in fig. 10 seems to continue the general trend of those given in ref. $/ 40$ / for heavier Cs nuclei.

The authors would like to thank Prof. K.Ya.Gromov for his interest and support, Drs. D.J.Salamov and T.Kozlowski for helpful discussions, Dr. M.Gasior for supplying a low-noise scintillation counter, M.Jachim for carr ying out the isotope separations, Dr. W.D.Fromm for the computer processing of the gamma-spectra and M.Honusek and W.Habenicht for their help during the experiments.

Note added in proof: During the course of publication of the present paper, a new work on the ${ }^{127} \mathrm{Ba}$ decay has been published by Pathak and Preiss (Phys. Rev., Cll, 1762 (1975)). Their results are in agreement with the data presented here.

## References

1. K.Sheline, T.Sikkeland., R.N, Chandra. Phys.Rev. Lett., 7, 446 (1961).
2. K.Kumar, M.Baranger. Phys.Rev.Lett., 12, 73 (1964).
3. D.A.Arseniev, A.Sobiczewski, V.G.Soloviev. Nucl. Phys., A126, 15 (1969).
4. J.E.Clarkson, R.M.Diamond, F.S.Stephens. J.Perlam. Nucl. Phys., A93, 272 (1967).
5. D.Ward, R.M.Diamond, F.S.Stephens. Nucl.Phys., A117, 309 (1968).
6. J.Bergstrom, C.J.Herrlander, A.Kerek, A.Luukko. Nucl.Phys., A123, 99 (1969).
7. W.Kutschera, W.Dehnhardt, O.C.Kistner, P.Kump, B. Povh, H.J.Sann. Phys.Rev., C5, 1958 (1972).
8. T.W.Conlon. Nucl.Phys., A161, 289 (1971).
9. T.W.Conlon. Nucl.Phys., A213, 445 (1973).
10. K.F.Alexander, W.Neubert, H.Rotter, S.Chojnacki, Ch.Droste, T.Morek. Nucl.Phys., A133, 77 (1969).
11. J.M.D'Auria, H.Bakhru, J.L.Preiss. Phys.Rev., 172, 1176 (1968).
12. R.Arlt, G.Beyer, H.-G.Ortlepp, H.Tyrroff, E.Herrmann, H.Haupt, A.Jasinski. Conf. Nucl. Spectr. Nucl. Theory, Dubna, 1973, p. 98.
13. R.Arlt, G.Beyer, H.Haupt, E.Herrmann, A.Jasinski, G.Musiol, W.Neubert, H.-G.Ortlepp, H.Tyrroff. Proc. Int. Conf. Nucl. Phys., Münich, 1973, p. 694.
14. H.-G.Ortlepp, H.Haupt, A.Jasinski, M.Jachim. Conf. Nucl. Spectr. Nucl. Theory, Kharkov, 1974, p. 87.
15. G.Beyer, H.-G.Ortlepp, A.Jasinski, M.Jachim. Conf. Nucl. Spectr. Nucl. Theory, Leningrad 1975, p. 90.
16. H.-U.Siebert, H.-G.Ortlepp, A.Jasinski, M.Jachim. Conf. Nucl.Spectr. Nucl. Theory, Dubna, 1975, p. 100.
17. Z.Malek, G.Pfepper. JINR, 12-4013, Dubna, 1968.
18. G.Musiol, V.I.Raiko, H.Tyrroff. JINR, P6-4487, Dubna, 1969.
19. G.Beyer, E.Herrmann, A.Piotrowski, V.I.Raiko, H.Tyrroff. Nucl. Instr. and Meth., 96, 437 (1971).
20. A.Latuszynski, K.Zuber, A.Potempa, W.Zuk. JINR, 6-7469, Dubna, 1973.
21. R.S.Hager, E.C.Selzer. Nucl.Data , A4, 397 (1968).
22. K.Andert, R.Arlt, M.Honusek, H.-U.Siebert,
A.I.Kalinin, S.V.Medved, G.Musiol, H.-G.Ortlepp,
A.N.Sinaev, W.Habenicht, H.Strusny. JINR, P6-8564, Dubna, 1975.
23. F.Charoenkwan. Nucl.Instr. Meth., 34, 93 (1965).
24. L.N.Abesalashvili, Zh.Zhelev, V.G.Kalinnikov, Ya.Liptak, U.Nazarov, Ya.Urbanets. JINR, P6-3348, Dubna, 1967.
25. V.S.Butsev, Zh.Zhelev, V.G.Kalinnikov, A.V.Kadiav-
tseva, Ya.Liptak, F.Molnar, U.Nazarov, Ya.Urbanets JINR, P6-3541, Dubna, 1967.
26. N.Zeldes, A.Grill, ASimievic. Mat. Fys. Skr. Dan. Vid. Selsk., 3, No. 5 (1967).
27. H.W.Taylor, B.Sing, F.S.Trato, J.D.King. Nucl. Phys., A179, 417 (1972).
28. K.Ishi, T.Aoki, S.Kageyama. J.Phys.Soc.Japan, 34, 285 (1973).
29. C.M.Lederer, J.M.Hollander, J.Perlman. Table of Isotopes (John Wiley and Sons, Inc. New York, 1967).
30. R.Arlt, A.Jasinski, W.Neubert, H.-G.Ortlepp. JINR, E6-7762, Dubna, 1974; Acta Phys. Pol., B6, 433 (1975).
31. O.B.Dabousi, M.H.Prior, H.A.Shugart. Phys.Rev., C3, 1326 (1971).
32. A.Spalek, I.Rezanka, J.Frana, A.Mastalka. Zeitsch. Phys., 204, 129 (1967).
33. S.Jha, N.F.Peak, W.J.Knox, E.C.May. Phys.Rev., C6, 2193 (1972).
34. A.G. de Pinko, J.V.Goldstein, J.M.F.Veronymo. An. Acad.Brasil. Ciene., 43, 63 (1971).
35. J.V.Goldstein, A.G.De Pinho. Can.J.Phys., 49, 1794 (1971); Phys.Rev., C4, 653 (1971).
36. S.G.Nilsson, Kgl. Dansk. Vid. Selsk. Mat. Fys. Medd., 29, No. 16 (1955).
37. B.S.Reehal, R.A.Sorensen. Phys.Rev., C2, 819 (1970).
38. D.J.Salamov. Prog. Thes. XXV Conf. on Nucl. Spectr. \& Nucl. Struct., Leningrad, 1975, p. 198.
39. L.S.Kisslinger, R.A.Sorensen. Rev.Mod.Phys., 35, 853 (1963).
40. J.Fechner, A.Hammerfahr, A.Kluge, S.K.Sen, H.Toschincki, J.Voss, P.Weight. Nucl.Phys., A130, 545 (1969).

Received by Publishing Department on October 22, 1975.

