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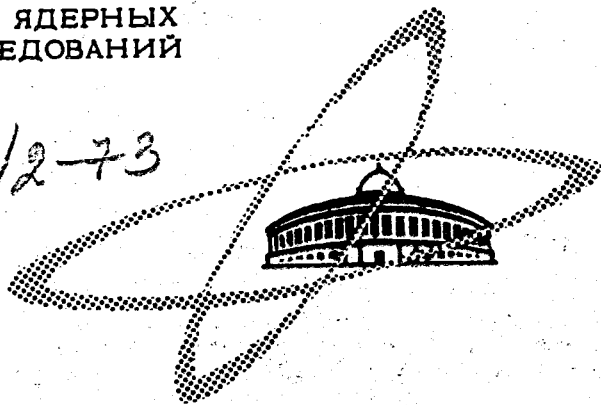
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ЛАБОРАТОРИЯ ЯДЕРНЫХ РЕАКЦИЙ

BETA DECAY OF ²²F

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BETA DECAY OF ^{22}F

Submitted to Nuclear Physics

Объединенный институт
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БИБЛИОТЕКА

1. Introduction

In the last years the interest in the properties of nuclei far from the region of beta stability has strongly increased. However, the experimental information about the neutron-rich isotopes of light elements is still poor, and therefore, new measurements of masses and half-lives and the studies of decay properties of these nuclei are very useful.

In 1965 Vaughn et al.^{/1/} reported on the production of ^{22}F in the reaction $^{22}\text{Ne}(n,p)^{22}\text{F}$. On the basis of the decay scheme and the measured beta end-point energy of 11.2 ± 0.6 MeV they determined a mass excess $\Delta = 4.5 \pm 0.6$ MeV for ^{22}F .

In 1969 Stokes and Young^{/2/} deduced from the kinematics of the reaction $^{22}\text{Ne}(t, ^3\text{He})^{22}\text{F}$ a mass excess value of 2.828 ± 0.030 MeV.

In the present work we have investigated the beta decay of ^{22}F and the following gamma de-excitation. New data on the decay scheme of ^{22}F are obtained.

2. Experimental Procedure

The experiment has been performed using the mass separator EMSONHIB^{/3/}. The separator operates on-line with the 310-cm heavy ion cyclotron of the JINR at Dubna.

The ^{22}F nuclei were produced in the two-nucleon-transfer reaction $^{181}\text{Ta}(^{22}\text{Ne}, ^{22}\text{F})$. The ^{22}Ne ion beam, accelerated to 174 MeV, bombarded a 20 μm thick metallic tantalum target. Separated nuclear reaction products were collected on a 100 μm thick aluminium foil. Beta and gamma detectors were placed directly behind the collecting foil of the mass separator.

Beta particles were detected in a $dE/dx-E$ scintillation spectrometer consisting of a plastic scintillator (0.5 mm thick, 35 mm in diameter) and a 70 x 70 - mm stilbene crystal. The energy calibration was performed with ^{207}Bi , ^{144}Pr , ^{23}Ne and ^{20}F electron emitters^{/4/} under exactly the same conditions as those for the measurement of the ^{22}F spectrum (short-lived nuclei of ^{23}Ne and ^{20}F were also produced in the $^{181}\text{Ta}+^{22}\text{Ne}$ multinucleon-transfer reaction). The operation of the detecting system up to 10.4 MeV was checked by beta particles from the ^{16}N decay.

The gamma spectrum was measured by a 30-cm³ Ge(Li) detector. The gamma rays^{/5/} from a ^{226}Ra source and the background gamma lines from the reaction $^{56}\text{Fe}(n, \gamma)$ (ref. 6) were used for the calibration.

Ions with mass number 22 were collected on the catcher of the mass separator. As shown by Artukh et al.^{/7/}, the cross section for the formation of ^{22}O , ^{22}Na and ^{22}Mg are much smaller than that for ^{22}F . Of all produced nuclei with $A = 22$, the radioactivity of ^{22}F was measured preferentially. Due to the large cross section for the one-nucleon-transfer reaction $^{181}\text{Ta}(^{22}\text{Ne}, ^{21}\text{F})$, a small amount of ^{21}F was separated in form of HF molecular ions. Background spectra were measured without interrupting the target irradiation and with the switched-off mass separator.

The half-life measurement was performed under the pulsed operating conditions of the cyclotron and the separator. In this measurement only gamma quanta of $E_\gamma > 0.7$ MeV were recorded in order to eliminate the gamma rays of ^{21}F . Beta particles were absorbed in 6 mm of lead.

3. Experimental Results and Discussion

The beta spectra of ^{22}F and the background are shown in the upper part of fig. 1. In the beta spectrum we do not observe the intense (33%) beta transition of 11.2 ± 0.6 MeV reported by Vaughn et al.^{/1/}

In the lower part of fig. 1 the Fermi-Kurie plots of ^{20}F and ^{22}F are shown. Since our measurements were carried out at a small distance between the scintillation detector and the collecting foil, the Fermi-Kurie plots are not straight lines. Figure 1 shows that the difference between the end-point energies of ^{20}F and ^{22}F is very small. The maximum beta energy $E_{\beta} = 5.419 \pm 0.013$ MeV (99.8%) of ^{20}F is well known^{/8/}. Therefore, we are able to determine the maximum beta energy of ^{22}F with high accuracy. From a least square fit to the Fermi-Kurie plots we derived the ^{22}F maximum beta energy $E_{\beta} = 5.5 \pm 0.1$ MeV.

The measured gamma spectrum is shown in fig. 2. All gamma transitions fit into the known level scheme^{/9/} of ^{22}Ne with the exception of the 1.90-MeV line. The gamma intensities obtained agree with the branching ratios reported by Kutschera et al.^{/10/}

The decay curve of ^{22}F is shown in fig. 3. The half-life, determined by means of a least square fit procedure, is 4.2 ± 0.1 sec.

The proposed decay scheme of ^{22}F is shown in fig. 4. The beta decay of ^{22}F leads to the excitation of the 5.52-MeV (4^+), 6.35-MeV (6^+) and 7.54-MeV levels in ^{22}Ne . The beta branching ratios given in table 1 were determined under the assumption that all the states of ^{22}Ne excited by the beta decay of ^{22}F de-excite through the 1.28-MeV level. The missing intensities of about 16% in the total sum of beta branching intensities may be due to beta transitions to higher excited states of ^{22}Ne de-exciting by high-energy gamma transitions for which the Ge(Li) detector efficiency is too small. On the other hand, the excessive intensity of the 1.28-MeV transition cannot be caused by the contribution of the decay of long-lived ^{22}Na , because the background spectrum (fig. 2) does not contain its 1.28-MeV line. The background spectrum was measured immediately after the end of the gamma measurement.

All the values of $\log ft$ (table 2) calculated using the tables of Dzhelepov et al.^{/11/} are within the empirical range for allowed transitions. This result and the proposed decay scheme allow us to assign a spin and parity value of 5^+ to the ground state of ^{22}F . Consequently, the 5.64-MeV level of ^{22}Ne with a spin and parity of 2^+ or 3^+ can not be excited by an allowed beta transition. According to our gamma measurement we suggest the existence of a 7.54-MeV level which decays to the 5.64-MeV level by the 1.90-MeV gamma transition.

The proposed decay scheme and the measured maximum beta energy of $E_{\beta} = 5.5 \pm 0.1$ MeV lead to the value of $Q_{\beta} = 11.0 \pm 0.1$ MeV.

On the basis of the ^{22}Ne mass excess $\Delta = -8.025$ MeV (ref. 12) we obtained for the mass excess of ^{22}F the quantity $\Delta = 3.0 \pm 0.1$ MeV. This result agrees fairly well with the value 2.828 ± 0.030 MeV reported by Stokes and Young ^{12/1}.

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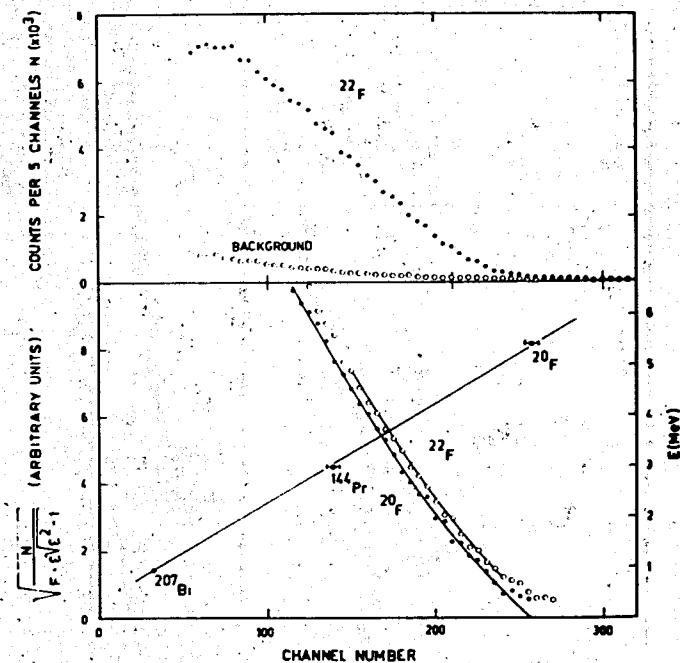


Fig. 1. Spectra of the ^{22}F beta decay and the background (upper part). Fermi-Kurie plots of the ^{22}F and ^{20}F beta spectra with the energy calibration (lower part).

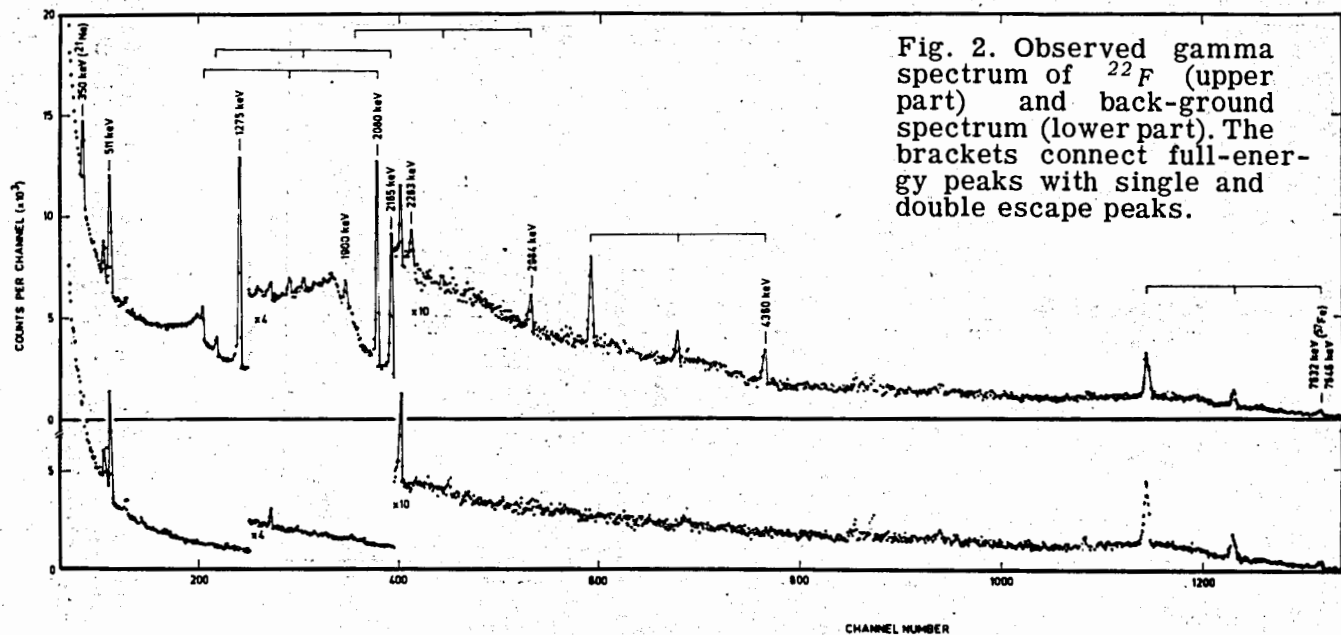


Fig. 2. Observed gamma spectrum of ^{22}F (upper part) and back-ground spectrum (lower part). The brackets connect full-energy peaks with single and double escape peaks.

Fig. 3. Decay curve of ^{22}F .

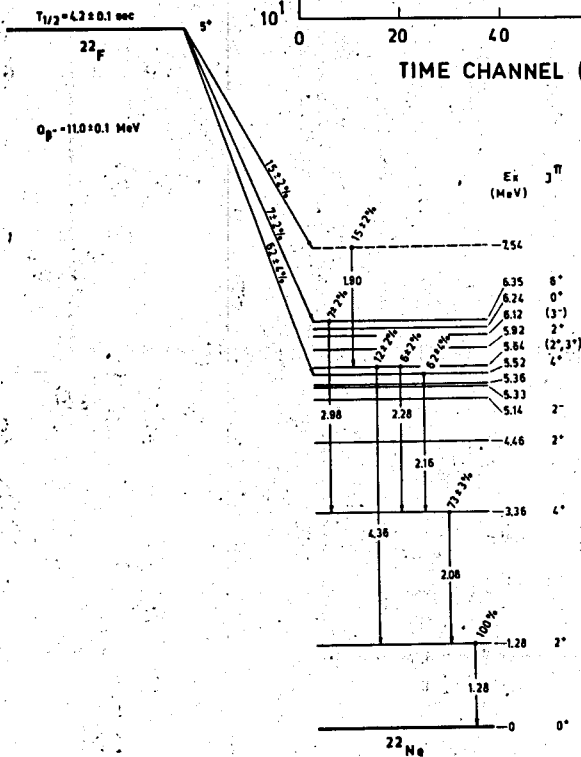
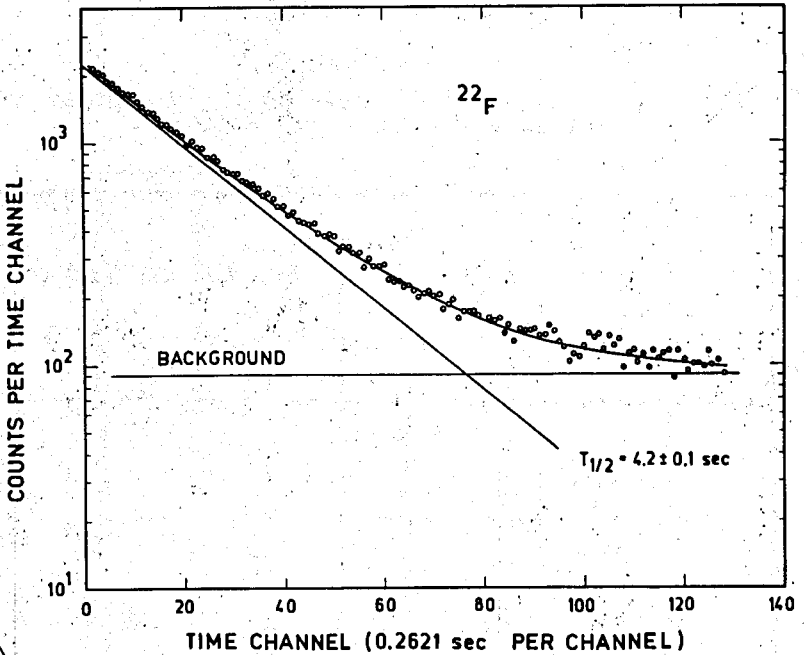


Fig. 4. Decay scheme of ^{22}F .

Table 1
 Energies and relative intensities of gamma transitions
 following the beta decay of ^{22}F

E_{γ} (keV)	Relative intensities (%)
1275 \pm 5	100
1900 \pm 5	15 \pm 2
2080 \pm 5	73 \pm 3
2165 \pm 5	62 \pm 4
2283 \pm 5	6 \pm 2
2984 \pm 5	7 \pm 2
4360 \pm 5	12 \pm 2

Table 2.
Characteristics of the ^{22}F beta decay

^{22}Ne level (MeV)	J^π	Beta end-point energy (MeV)	Beta branching ratio (%)	log ft
5.52	4^+	5.5 ± 0.1	62 ± 4	4.8
6.35	6^+	4.6 ± 0.1	7 ± 2	5.4
7.54	-	3.5 ± 0.1	15 ± 2	4.5