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## THREE-QUASIPARTICLE EXCITATIONS IN ${ }^{205}$ At

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Investigations of many-particle excitations of the type (j) ${ }^{n}$ are of theoretical interest, especially with regard to the determination of the residual interaction parameters. Recently the $\left(h_{9 / 2}\right)^{3}$ multiplet has been investigated in ${ }_{211}^{20} \mathrm{At}_{126}$, in which nucleus 3 protons are outside the double magic ${ }_{82}^{208} \mathbf{P b}_{126}$ core. This multiplet is fed by the $21 / 2^{\vec{~}}$ isomeric state, which is the highest spin state of the same configuration $/ 1 /$.

Such high spin states are preferably excited in heavy-ion compound nuclear reactions. At the U-300 heavy-ion cyclotron of the Laboratory for Nuclear Reactions, JINR Dubna we started, therefore, pulsed beam experiments for the investigation of similar isomeric states. By use of the reaction ${ }^{197} \mathrm{Au}\left({ }^{12} \mathrm{C}, 4 \mathrm{n}\right){ }_{85}^{205} \mathrm{At}_{120}$ an isomeric state with the half-life $T 1 / 2=110 \pm 25 \mathrm{~ns}$ was found, which also can be ascribed (cf. below) to the $21 / 2^{-}$state of the $\left(\mathrm{h}_{9 / 2}\right)^{3}$ configuration. Basing on two-dimensional time and energy spectra, the intensities of the prompt and delayed gamma transitions and the measured excitation functions we concluded the level scheme shown in fig. 1. A detailed description of the experiment will be published in a forthcoming paper $/ 2 /$.

In fig. 2 our level scheme is compared with those of the neighbouring even isotones $N=120$, namely $\quad{ }^{202} \mathbf{P b} / 3 /$ and
${ }^{204} \mathbf{P}_{0} / 4 /$. A similar but clearer comparison is possible ${ }^{/ 1 /}$ for the neighbouring closed neutron shell isotones ${ }^{210} \mathbf{P o}$ and ${ }^{211} \mathrm{~A}_{\mathrm{t}}$ with $\mathbf{N}=126$, which is also drawn in fig. 2. The level schemes of ${ }^{210} \mathrm{Po}$ and ${ }^{211} \mathrm{At}$ can be described by the pure $\mathrm{p}\left(\mathrm{h}_{\mathrm{g} / 2}\right)^{2}$ and $p\left(h_{9 / 2}\right)^{3}$ configurations.

In case of the $N=120$ isotones ${ }^{204} \mathbf{P o}_{0}$ and ${ }^{205} \mathrm{At}$ the situation is somewhat different. Going away from the closed neutron shell we find an increasing collective character of the first low-spin levels $I^{\pi}=2^{+}$and $4^{+}$in ${ }^{204} \mathbf{P o}$, for the two-neutron quasiparticle states are increasingly admixed. For example the $2^{+}$and $4^{+}$states of the $\left(\mathrm{f}_{0 / 2}\right)^{2}$ neutron configuration in ${ }^{202} \mathbf{P b}$ are placed at 960 keV and 1384 keV , respectively $/ 3 /$. On the other hand, the $I^{\pi}=6^{+}$and $8^{+}$levels in ${ }^{204} \mathbf{P o}$ remain fairly pure proton states of the $p\left(h_{9 / 2}\right)^{2}$ configuration, for the corresponding neutron configurations in ${ }^{202} \mathbf{P b}$ which give such high spin values, have considerably higher energies and should, therefore, be of little influence.

It seems to $u s$, that the $I^{\pi}=21 / 2^{-}$- level in ${ }^{205} \mathrm{At}$ corresponds to the $I^{\pi}=6^{+}$level in ${ }^{204} P$ ond may be regarded as the highest spin state with $I=I_{\text {max }}$ of the $\left(h_{g / 2}\right)^{3}$ proton configuration, similar to the situation in ${ }^{211} \mathrm{At}$. The excitation energies of the $I^{\pi}=21 / 2^{-}$levels in ${ }^{208} \mathrm{At}$ and ${ }^{211} \mathrm{At}$ are nearly the same ( 1492 keV and 1416 keV , respectively) ancı somewhat lower than the $I^{\pi}=6^{+}$levels in ${ }^{204}{ }^{4} \mathbf{P o}$ and ${ }^{210} \mathbf{P o}$ (1624 keV and 1472 keV , respectively). Deviating from the ${ }^{211} \mathrm{At}_{\mathrm{t}}$ deexcitation we observed two gamma transitions of 125 keV and 191 keV of nearly equal intensity, outgoing from the isomeric $21 / 2^{-}$level. The intensity balance of the whole decay scheme may only be fulfilled if we assume an E2 multipole order for both the 125 keV and 191 keV transitions. The $\mathbf{B}(E 2)$ value of the 125 keV transition (table)
does not deviate much from the theoretical value $\mathbf{B}(\mathbf{E} 2)=61 e^{2} \mathrm{fm}^{4}$
 transition gives a very low $\mathbf{B}(\mathbf{E} 2)$ value. We, therefore, assume the existence of a second $17 / 2^{-}$Ievel not belonging to the $\left(h_{g / 2}\right)^{3}$ configuration. Corresponding to the coupling rules ${ }^{/ 6 /}$ such a $17 / 2^{-}$ state may be formed by coupling the $\left(\mathrm{h}_{9 / 2}\right)_{9 / 2}^{3}$ proton state to the second $4^{+}$state in ${ }^{202} \mathbf{P b} \quad / 3 /$ with the structure $n_{1}\left(f_{5 / 2}\right) n_{2}\left(p_{3 / 2}\right)$. The 191 keV transition from the $21 / 2^{-}$level to this $17 / 2^{-}$level at 1301 probably goes via an $\left(\mathrm{h}_{\mathrm{g} / 2}\right)^{3}$ admixture of the $17 / 2^{-}$state. From the above mentioned cause we expect, that the states with spins $17 / 2^{-}, 15 / 2^{-}, 13 / 2^{-}$and $11 / 2^{-}$are of complicated structure. They correspond to the $4^{+}$and $2^{+}$levels in ${ }^{202} \mathbf{P b}$ and ${ }^{204} \mathbf{P o .}$ The assumption, that these states are not pure proton states is supported by the facts, that no E 2 transition from the 17/2- level at 1367 keV to the $13 / 2^{-}$level at 663 keV could be found, and that the 332 keV gamma transition to the $15 / 2^{-}$level is of mainly M1 character, which assumption results in a better intensity balance than a pure E2 transition.

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Table

| E (keV) | $I_{\gamma}$ | (relativ) | $a_{\text {tot }}$ | branching <br> ratio $\beta$ | B (E 2) <br> $\mathbf{e}^{2} \mathrm{fm}^{4}$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| 125 | $55 \pm 5$ | 2.78 | $0.37 \pm 0.07$ | $31 \pm 9$ |  |
| 191 | $49 \pm 4$ | 0.58 | $2.71 \pm 0.49$ | $3.4 \pm 1.4$ |  |

Fig. 1. Level scheme of 205 m At.



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9 / 2-\frac{0_{85}^{211} \mathrm{At}_{128}}{0_{84}^{20} \mathrm{Po}_{126}}
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Fig. 2. Comparison of many-particle levels in the $N=120$ and $\mathbf{N}=126$ isotones.

