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**THE CLASSICAL  
WEYL-VON NEUMANN THEOREM  
IN ALGEBRAS  
OF UNBOUNDED OPERATORS**

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The classical Weyl-von Neumann theorem states the following:

Theorem 1/1/: Let  $H$  be a selfadjoint operator. Then for any  $\varepsilon > 0$  there is an  $A = A^* \in \mathcal{J}_2(\mathcal{H})$  (Hilbert-Schmidt operators) such that  $\|A\|_2 < \varepsilon$  and  $H+A$  has pure point spectrum, i.e., there exist a real sequence  $(\lambda_n)$ , an orthonormal basis  $(\varphi_n) \subset \mathcal{D}(H)$  with  $(H+A)\varphi_n = \lambda_n \varphi_n$ .

In this paper we give two variants (Propositions 1 and 3) of this result in algebras of unbounded operators. The proofs are adaptations of the classical one /1/. We give the details of the proof of Prop.1, because we think that the ideas can be used to extend the results further. It is also possible to give a short proof of Propositions 1 and 3 which uses the statement of the classical theorem. We give such a proof for Prop.3, then Prop.1 is obtained as a corollary.

We start with some necessary notions and notations /2/, /4/. Let  $\mathcal{D}$  be a dense linear manifold in a separable Hilbert space  $\mathcal{H}$ . By  $\mathcal{L}^*(\mathcal{D})$  we denote the following set of linear operators:  $\mathcal{L}^*(\mathcal{D}) = \{A: A\mathcal{D} \subset \mathcal{D}; A^*\mathcal{D} \subset \mathcal{D}\}$ . With the usual operations and the involution  $A \rightarrow A^* = A^*|_{\mathcal{D}}$  becomes a  $\kappa$ -algebra. It can be shown that for  $1 \leq p \leq \infty$  the sets

$$\mathcal{J}_p(\mathcal{D}) = \{A \in \mathcal{L}^*(\mathcal{D}): \overline{BAC} \in \mathcal{J}_p(\mathcal{H}) \text{ for all } B, C \in \mathcal{L}^*(\mathcal{D})\}$$

are two-sided  $\kappa$ -ideals in  $\mathcal{L}^*(\mathcal{D})$ . Here  $\overline{\phantom{x}}$  denotes the closure of the operator  $D$  and  $\mathcal{J}_p(\mathcal{H})$  are the usual ideals in  $\mathcal{B}(\mathcal{H})$ . We need also the two-sided  $\kappa$ -ideal  $\mathcal{B}(\mathcal{D}) = \{A \in \mathcal{L}^*(\mathcal{D}): BAC \text{ bounded for all } B, C \in \mathcal{L}^*(\mathcal{D})\}$ .

An important class of domains  $\mathcal{D}$  is given by  $\mathcal{D} = \mathcal{D}^\infty(H) = \bigcap_{n=0}^{\infty} \mathcal{D}(H^n)$ , where  $H=H^*$  is an (unbounded) selfadjoint operator, without loss of generality  $H \geq 1$ . In this case  $\mathcal{J}_p(\mathcal{D})$  and  $\mathcal{B}(\mathcal{D})$  can be characterized as follows /4/:  $\mathcal{J}_p(\mathcal{D}) = \{A \in \mathcal{L}^*(\mathcal{D}): H^n A, H^{n+1} A^* \in \mathcal{J}_p(\mathcal{H}) \text{ for all } n \in \mathbb{N}\} = \{A \in \mathcal{L}^*(\mathcal{D}): A H^n, A^* H^{n+1} \in \mathcal{J}_p(\mathcal{H}) \text{ for all } n \in \mathbb{N}\}$ . and for  $\mathcal{B}(\mathcal{D})$  analogously. Now we state the result:

Proposition 1: Let  $\mathcal{D} = \mathcal{D}^\infty(H)$ ,  $H=H^*$ . Then for any  $\delta > 0$  there is an  $A = A^* \in \mathcal{J}_2(\mathcal{D})$  such that  $\|A\|_2 < \delta$  and  $H+A$  has pure point spectrum.

The proof of this proposition is based on an adaptation of the corresponding lemma in the classical case /1/.

Lemma 2: Let  $\mathcal{D} = \mathcal{D}^\infty(H)$ ,  $H=H^*$ . Then for any  $\varphi \in \mathcal{D}$ ,  $\varepsilon > 0$ ,  $m \in \mathbb{N}$  there are a finite dimensional projection  $P \in \mathcal{L}^*(\mathcal{D})$  and an  $Y=Y^* \in \mathcal{L}^*(\mathcal{D})$  such that

- i)  $\|(1-P)\varphi\| < \varepsilon$
- ii)  $\|H^m Y\|_2 < \varepsilon$
- iii)  $H+Y$  is reduced by  $P\mathcal{H}$ .

Remark that in what follows we will not indicate the closure  $\overline{\phantom{x}}$  for



bounded  $\mathcal{D}$  if there cannot arise confusions. Statement iii) means that  $\mathcal{D}_0 = P\mathcal{R} \cap \mathcal{D}$  is dense in  $P\mathcal{R}$ ,  $(H+Y)\mathcal{D}_0 \subset \mathcal{D}_0$  and  $\overline{(H+Y)\mathcal{D}_0}$  is a selfadjoint operator in  $\mathcal{R}_0 = \mathcal{D}_0$ , or what is the same,  $P$  commutes with  $H+Y$ .

Proof of Lemma 2: Let  $a > 0$  be so large that

$$(1) \quad \| [1 - E(I_a)] \varphi \| < \varepsilon$$

where  $I_a = (-a, a]$  and  $E(\cdot) \in \mathcal{L}^+(\mathcal{D})$  are the spectral projections corresponding to  $H$ . For each  $n \in \mathbb{N}$  put  $E_k = E(\left(\frac{2k-n-2}{n} a, \frac{2k-n}{n} a\right])$ ,  $k=1, \dots, n$ . Remark that  $E_j E_k = \delta_{jk} E_k$ ,  $E_k \in \mathcal{L}^+(\mathcal{D})$  and consequently  $\mathcal{R}_k = E_k \mathcal{R} \subset \mathcal{D}$ . Put  $\varphi_k = E_k \varphi$ ,  $\psi_k = \varphi_k / \|\varphi_k\|$  if  $\varphi_k \neq 0$  and  $\psi_k = 0$  otherwise. Thus,  $(\psi_k)$  is an orthonormal system satisfying

$$(2) \quad \sum_{k=1}^n \|\varphi_k\| \psi_k = \sum_{k=1}^n \varphi_k = E(I_a) \varphi$$

and the  $\psi_k$  are "almost-eigenvalues" of  $H$  in the sense that for

$$\lambda_k = \frac{2k-n-1}{n} a$$

$$(3) \quad \| (H - \lambda_k) \psi_k \| = \| (H - \frac{2k-n-1}{n} a) \psi_k \| \leq a/n.$$

Let  $P$  denote the orthoprojection on the space spanned by  $\psi_1, \dots, \psi_n$ , then

$$(4) \quad P \in \mathcal{L}^+(\mathcal{D}), \dim P \leq n$$

$$(5) \quad (1-P)\psi_k = 0 \text{ implies } \|(1-P)H\psi_k\| = \|(1-P)(H - \lambda_k)\psi_k\| \leq a/n \text{ and together with (2): } (1-P)E(I_a)\varphi = 0.$$

$$(6) \text{ It is } (1-P)H\psi_k \in \mathcal{R}_k, \text{ hence } \langle (1-P)H\psi_j, (1-P)H\psi_k \rangle = 0 \text{ for } j \neq k.$$

From (5) the first statement of the Lemma follows:  $\|(1-P)\varphi\| =$

$$= \|(1-P)(\varphi - E(I_a)\varphi)\| \leq \|(1-P)\varphi\| + \|(1-E(I_a))\varphi\| < \varepsilon.$$

Put  $-Y = (1-P)HP + ((1-P)HP)^\dagger$ , then from  $H = PHP + (1-P)H(1-P) + (1-P)HP + ((1-P)HP)^\dagger$  one gets  $H+Y = PHP + (1-P)H(1-P)$ . Clearly, this operator is reduced by  $P\mathcal{R}$ , i.e., iii) is valid. Next let us estimate  $\|H^m Y\|_2$ . This is done by estimating  $H^m(1-P)HP$  and  $H^m P H(1-P)$  separately. Using  $H^m(1-P)H\psi_k \in \mathcal{R}_k$  it follows that  $\|H^m(1-P)HP\psi\|^2 = \|\sum_k \langle \psi_k, \psi \rangle H^m(1-P)H\psi_k\|^2 = \sum_k |\langle \psi_k, \psi \rangle|^2 \|H^m(1-P)H\psi_k\|^2 \leq \|\psi\|^2 a^{2m} a^2 / n^2$  for all  $\psi \in \mathcal{R}$ . Here we used (6),  $(1-P)H\psi_k \in E(I_a)\mathcal{R}$  and  $H^m$  restricted to  $E(I_a)\mathcal{R}$  is bounded by  $a^m$ . Consequently  $\|H^m(1-P)HP\| \leq a^{m+1}/n$  and therefore  $\|H^m(1-P)HP\|_2 \leq a^{m+1}/n^{1/2}$  because the dimension of  $H^m(1-P)HP$  is smaller than  $n$ . To estimate  $H^m P H(1-P)$  one considers the adjoint operator, proceeds similarly and obtains the same estimation. Thus,  $\|H^m Y\|_2 \leq 2a^{m+1}/n^{1/2}$ . Because  $n$  can be chosen independent of  $a$ , ii) is proved.

q.e.d.

Proof of Proposition 1: Let  $(\varphi_n) \subset \mathcal{D}$  be a countable set, dense in  $\mathcal{R}$ . Apply Lemma 2 successively:

1.  $\varphi = \varphi_n$ ,  $\varepsilon = \delta/2$ ,  $m = 0$ . One obtains  $P_1, Y_1 \in \mathcal{L}^+(\mathcal{D})$  with i)-iii) as in Lemma 2.

2. Apply the Lemma to that part of  $H+Y_1$  which lies in  $(1-P_1)\mathcal{R} \cap \mathcal{D} = (1-P_1)\mathcal{D}$  with  $\varphi = (1-P_1)\varphi_n$ ,  $\varepsilon = \delta/4$ ,  $m = 1$ . This gives  $P_2, Y_2$  which we continue to whole  $\mathcal{D}$  by 0 and denote the resulting operators by the same symbol. Then  $H+Y_1+Y_2$  is reduced by  $P_1\mathcal{R}$  and  $P_2\mathcal{R}$ . Continuing this procedure one gets two sequences of operators in  $\mathcal{L}^+(\mathcal{D})$ :  $P_1, \dots$ , and  $Y_1, Y_2, \dots$  with  $\|(1-P_1 - \dots - P_k)\varphi_n\| \leq \delta/2^k$  and  $\|H^{k-1}Y_k\|_2 \leq \delta/2^k$  for all  $k$ . Put  $A_n = \sum_{k=1}^n Y_k$  and  $A = \sum_k Y_k$ . To prove that  $A = A^* \in \mathcal{J}_2(\mathcal{D})$  we show that  $H^m A \in \mathcal{J}_2(\mathcal{R})$  for all  $m \in \mathbb{N}$ . For  $r, m$  the estimation

$\|H^m \sum_{k=r}^{\infty} Y_k\|_2 \leq \sum_{k=r}^{\infty} \|H^k Y_k\|_2 \leq \sum_k \delta/2^k$  implies that  $(H^m A_n)$  is a Cauchy sequence in  $\mathcal{J}_2(\mathcal{R})$ . Now it is standard to see that  $H^m A_n \xrightarrow{u} H^m A \in \mathcal{J}_2(\mathcal{R})$ . The estimation above also gives  $\|A\|_2 < \delta$ . It remains to prove that  $H+A$  has pure point spectrum. As in the classical case one sees that  $\sum P_k = 1$  and  $P_k \mathcal{R}$  reduces  $H+A$ . Thus the finite dimensionality of each  $P_k \mathcal{R}$  and  $\sum P_k = 1$  gives the desired result that  $H+A$  has a complete system of eigenvectors.

q.e.d.

The following Proposition gives one of the possible generalizations of Proposition 1.

Proposition 3: Let  $\mathcal{D} = \mathcal{D}^-(T)$ ,  $H \in \mathcal{L}^+(\mathcal{D})$  so that the following conditions are fulfilled

- i)  $H$  is essentially selfadjoint
- ii) The spectral projections  $E_\lambda$  corresponding to  $\overline{H}$  are in  $\mathcal{L}^+(\mathcal{D})$ .
- iii)  $E(I)\mathcal{R} \subset \mathcal{D}$  for any bounded interval  $I \subset \mathbb{R}$ .

Then for any  $\delta > 0$  there is an  $A = A^* \in \mathcal{J}_2(\mathcal{D})$  with  $\|A\|_2 < \delta$  and  $H+A$  has pure point spectrum.

Proof: Put  $E_n = E((n, n+1])$ ,  $\mathcal{R}_n = E_n \mathcal{R}$ ,  $H_n = H|_{\mathcal{R}_n}$ ,  $\|T^m\|_{\mathcal{R}_n, \mathcal{R}_n} = a_{nm}$  for all  $n \in \mathbb{Z}$ ,  $m \in \mathbb{N}$ . Since we may suppose  $T \geq 1$ , it follows that  $a_{nm} \leq a_{n, m+1} \leq \dots$  for all  $m, n$ . Choose  $(\delta_n)$  so that  $\delta_n > 0$  and  $\sum \delta_n < \delta$ ,  $\sum a_{nm} \delta_n < \infty$ . On  $\mathcal{R}_n$  construct  $A_n = A_n^*$  according to the classical Weyl-von Neumann theorem so that  $\|A_n\|_2 < \delta_n$  and  $H_n + A_n$  has pure point spectrum as an operator on  $\mathcal{R}_n$ . Put  $A = \sum_n A_n$ . Then  $\|T^m A\|_2 \leq \sum_n \|T^m A_n\|_2 \leq \sum_n a_{nm} \delta_n = \sum_{n \leq -1} a_{nm} \delta_n + \sum_{n \geq 0} a_{nm} \delta_n \leq \sum_{n \leq -1} a_{nm} \delta_n + \sum_{n \geq m} a_{nm} \delta_n < \infty$ . For  $m=0$  it follows that  $\|A\|_2 < \delta$ . Clearly,  $H+A$  has pure point spectrum.

q.e.d.

As in the classical case one can get stronger statements about  $A$ . This is formulated as a corollary to Propositions 1 and 3.

Corollary 4: Suppose that the assumptions of Propositions 1 or 3 are fulfilled. Then for any  $p > 1$   $A$  can be chosen so that  $A \in \mathcal{J}_p(\mathcal{D})$ .

Proof: As in the classical case it is possible to choose  $A \in \mathcal{J}_q(\mathcal{K})$ ,  $q > 1$  arbitrary. From the polar decomposition  $A = |A|V$  it follows that  $|A| \in \mathcal{J}_q(\mathcal{K})$ . Moreover, considerations similar to those in the proof of Proposition 3 give  $\|T^m|A|\| = \|T^mA\| < \infty$ , i.e.  $|A| \in \mathcal{B}(\mathcal{H})$  and consequently  $|A|^\alpha \in \mathcal{B}(\mathcal{H})$  for all  $\alpha > 0$  [4]. Now let  $p > 1$  be given. Choose  $0 < \alpha < 1$  and  $q$  so that  $p = q/(1 - \alpha)$ . Construct  $A$  so that  $A \in \mathcal{J}_q(\mathcal{K})$ , then  $|A| \in \mathcal{J}_q(\mathcal{K})$  and  $|A|^{1-\alpha} \in \mathcal{J}_p(\mathcal{K})$ . In the decomposition  $T^mA = (T^m|A|^\alpha)(|A|^{1-\alpha}V)$  the first factor is in  $\mathcal{B}(\mathcal{K})$ , the second in  $\mathcal{J}_p(\mathcal{K})$ , hence  $T^mA \in \mathcal{J}_p(\mathcal{K})$ . This means  $A \in \mathcal{J}_p(\mathcal{H})$ .  
q.e.d.

Remark 5: The condition iii) of Proposition 3 seems to be very strong. It is of course fulfilled if  $T$  and  $H$  commute strongly. But there is also a large class of operators  $H$  non-commuting with  $T$  such that iii) holds. Such operators can be constructed as follows. Let  $T = \int_{\mu} dF_{\mu}$ ,  $[1, \infty) = \bigcup_n I_n$ ,  $I_n$  bounded and pairwise disjoint intervals,  $\mathcal{K}_n = F(I_n)\mathcal{K}$ . Then for appropriately chosen operators  $H_n = H_n^*$  in  $\mathcal{K}_n$  the operator  $H = \sum_n H_n$  satisfies the assumptions of Proposition 3. That means,  $H$  has a kind of block-decomposition with respect to  $T$ . It seems that condition iii) cannot be dropped completely.

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Классическая теорема Вейля-фон Неймана  
в алгебре неограниченных операторов

Дается вариант классической теоремы Вейля-фон Неймана в алгебрах неограниченных операторов с областью определения вида  $\mathcal{D} = \mathcal{D}^\infty(H)$ ,  $H = H^*$  - неограниченный самосопряженный оператор.

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The Classical Weyl-von Neumann Theorem  
in Algebras of Unbounded Operators

There is given a variant of the classical Weyl-von Neumann theorem in algebras of unbounded operator defined on domains  $\mathcal{D} = \mathcal{D}^\infty(H)$ , where  $H = H^*$  is an unbounded selfadjoint operator.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

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