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## LAGRANGIAN DISTRIBUTIONS

AND THEIR SINGULARITIES.
II. CALCULUS OF SINGULARITIES

AND APPLICATIONS

## 1. THE CASE WHEN THE RANK OF THE

## $\xi-$ PROJECTION OF $\Lambda$ IS CONSTANT*

Let $u \in \mathscr{F}^{\prime}\left(R^{n}\right)$ be a Lagrangian distribution, singspec $u \subset \Lambda$ and let the $\xi$-projection of $\Lambda$ has a constant rank, but is not necessarily one-to-one. Such is, for instance, the Schwartz kernel $\mathrm{K}(\mathrm{x}, \mathrm{y})$ of a pseudodifferential operator, where $(\mathrm{x}, \mathrm{y}) \in \boldsymbol{R}^{2 \mathrm{n}}$. In this case $\Lambda$ is the normal bundle of the diagonal $x=y$, so that the rank of the projection is not $2 n$ but $n$, hence the projection is not one-to-one. It is possible to carry out the procedure of the previous section for defining the symbol of $u$, but it is simpler in that case to use another symbol which also characterizes, the singularities, and which for the example of Schwartz kernel coincides with the usual symbol of pseudodifferential operators.

First of all let us prove that $\Lambda$ may be represented by

$$
\begin{align*}
& x_{I}=F_{1}\left(x_{\bar{I}}, \xi_{I}\right),  \tag{1.1}\\
& \xi_{-}=F_{2}\left(\xi_{I}\right),
\end{align*}
$$

where $\bar{I} \overline{\mathrm{I}}$ are subsets of $\{1,2, \ldots, n\}, \mathrm{I} \quad \overline{\mathrm{I}}=\boldsymbol{6}, \mathrm{I} \cup \overline{\mathrm{I}}=\{1, \ldots, \mathrm{n}\}$. Indeed there obviously exist representations of $\Lambda$ of the form

$$
\begin{align*}
& \mathrm{x}_{\mathrm{I}}=\Phi_{1}\left(\mathrm{x}_{\overline{\mathrm{I}}}, \xi_{\mathrm{I}}\right) \\
& \xi_{\overline{\mathrm{I}}}=\Phi_{2}\left(\mathrm{x}_{\overline{\mathrm{I}}}, \xi_{\mathrm{I}}\right) \tag{1.2}
\end{align*}
$$

Between all the representations of $\Lambda$ of the form (1.2) we choose one for which the number of elements of i, i.e., the number of $\xi$-variables in the right-hand part, is maximal. In that case $\Phi_{2}$ doesn ${ }^{\text {t }}$ depend on $x_{\bar{I}}$, i.e. $\partial \Phi_{2} / \partial x_{\bar{I}}=0$, because, if for instance, $j \in \bar{I}, \quad \xi_{j}=\Phi_{2, j}\left(x_{I}, \xi_{I}\right)$ and $\frac{\partial \Phi_{Q_{i}^{2}}}{\partial \mathrm{z}_{\mathrm{i}}} \neq 0$ for some $i \in \overline{\mathrm{I}}$, then, by the implicit function theorem, $x_{i}$ will be represented as a function of $\xi_{i}, \xi_{I I}, x_{1}\left\{_{i\}}\right.$ and substituting $x_{i}$ in (1.2) we get a representation of $\Lambda$ of the form (1.2), in
 tradicts the assumption of maximality of the number of $\xi$ variables in (1.2).

[^0]It is easy to see that the number of elements of $I$ in the representation (1.1) equals the rank $r$ of the $\xi$-projection of $\Lambda$.

Now consider the trace of $u$ on the subspace $x_{T}=x_{T}$. Such . trace exists since (see, ref. ${ }^{/ 5 /}$ ) the conormal bundle of the subspace $x_{\bar{I}}=x^{\circ}$ does not intersect $\Lambda$, as is easily seen from (1.1). We denote this trace by $u\left(0, x_{1}^{\circ}\right.$ ) The singular spectrum of $u\left(\circ, x^{\circ}\right)$ is contained in the Lagrangian manifold $\Lambda^{\prime} \subset R_{x}^{2 r}$ determined by the equations

$$
\begin{equation*}
\mathrm{x}_{\mathrm{I}}=\mathrm{F}_{1}\left(\mathrm{x}_{\overline{\mathrm{I}}}^{\circ}, \xi_{\mathrm{I}}\right) . \tag{1.3}
\end{equation*}
$$

Thus the $\xi_{I}$-projection of $\Lambda^{\prime}$ is one-to-one and $u\left(o, x_{\bar{I}}^{0}\right)$ is a primitive Lagrangian distribution. Take the symbol $\varrho\left[u\left(0, x_{\mathrm{I}}\right)\right]\left(\mathrm{x}_{\mathrm{I}}, \xi_{\mathrm{L}}\right)$ as defined in section $2^{\prime 9}$. Here $\left(\mathrm{x}_{\mathrm{I}}, \xi_{\mathrm{I}}\right) \in \Lambda^{\prime}$ and consequently (1.3) holds. We denote

$$
\mathscr{G} \cdot[u]\left(x_{\bar{I}}, \xi_{\mathrm{I}}\right)=\left[u\left(0, x_{\mathrm{I}}\right)\right]\left(\mathrm{F}_{1}\left(\mathrm{x}_{\overline{\mathrm{I}}}, \xi_{\mathrm{I}}\right), \xi_{\mathrm{I}}\right)
$$

and call it the modified symbol of $u$.
If $u$ is the kerne1 $K(x, y)$ of a pseudodifferential operator $\operatorname{Av}(\mathrm{x})=\int \mathrm{e}^{\mathrm{i}(\mathrm{x}-\mathrm{y}) \xi} \quad \mathrm{a}(\mathrm{x}, \xi) \mathrm{v}(\mathrm{y}) \mathrm{d} \xi \mathrm{dy}, \quad \mathrm{v} \in \mathrm{C}_{0}^{\infty}$,
then we can take $\mathrm{x}_{\mathrm{I}}=\mathrm{y}, \mathrm{x}_{\overline{\mathrm{I}}}=\mathrm{x}, \xi_{\mathrm{I}}=\xi, \xi_{\overline{\mathrm{I}}}=\eta$. The equations (1.1) are

$$
\mathrm{y}=\mathrm{x}, \quad \eta=-\xi
$$

and it is easy to calculate that

$$
\mathcal{G}^{\prime}[K](\mathrm{x}, \xi)=\mathbf{a}(\mathrm{x}, \xi),
$$

i.e., in this case the modified symbol of $K$ is the usual symbol of the pseudidifferential operator.

## 2. CALCULUS OF SINGULARITIES

For a distribution $u \in \mathscr{L}_{0}$ it is possible to calculate $A F^{l o c}[u]$ provided singspec $u$ and $G[\dot{u}]$ are given.
Theorem 2.1. If $u \in \mathscr{L}_{0}$, then the following formula holds.

$$
\begin{equation*}
\operatorname{asym} F[u \phi](\lambda \xi)=\mathrm{e}^{-\mathrm{i} \lambda g(\xi)} \left\lvert\, \sum_{|a|=0}^{\infty} \frac{1}{a!}\left(\frac{1}{\lambda}\right)^{|a|} \partial_{\xi}^{a} \mathcal{G}[u]\left(\mathrm{g}^{\prime}(\xi), \xi\right) \phi^{(a)}\left(\mathrm{g}^{\prime}(\xi)\right)\right., \phi \in \mathrm{C}_{0}^{\infty} . \tag{2,1}
\end{equation*}
$$

Proof Let $\left(x^{0}, \xi^{\circ}\right) \in \Lambda$. By the Taylor formula we have
$\phi(\mathrm{x})=\sum_{|a| \leq \mathrm{N}} \frac{1}{a!}\left(\mathrm{x}-\mathrm{x}^{0}\right)^{a} \phi^{(a)}\left(\mathrm{x}^{0}\right)+\sum_{|\beta|=\mathrm{N}+1} \frac{1}{\beta!}\left(\mathrm{x}-\mathrm{x}^{\mathrm{o}}\right)^{\beta}{ }_{\phi^{(\beta)}}\left(\mathrm{x}^{\left.\mathrm{o}+\left(\mathrm{x}-\mathrm{x}^{0}\right) \theta(\mathrm{x})\right) .}\right.$
Thus we get for $\phi \in 1_{0}^{\infty}$ ( $x^{0}$ )
$\mathrm{F}[\mathrm{u} \phi]\left(\lambda \xi^{\circ}\right)=\mathrm{F}[\psi \mathrm{u} \phi]\left(\lambda \xi^{\circ}\right)=\sum_{|a| \leq \mathrm{N}} \frac{1}{a!} \phi^{(a)}\left(\mathrm{x}^{\circ}\right) \mathrm{F}\left[\left(\mathrm{x}-\mathrm{x}^{\circ}\right)^{a} \psi \mathrm{u}\right]\left(\lambda \xi^{\circ}\right)+$

$$
+\underset{|a|=N+1}{ } \quad \frac{1}{a!} \cdot F\left[\left(x-x^{\circ}\right)^{a} \psi u \phi^{(a)}\left(x^{\circ}+\left(x-x^{\circ}\right) \theta(x)\right)\right]\left(\lambda \xi^{\circ}\right) .
$$

The last term may be represented in the form

$$
\begin{align*}
& F\left[\left(x-x^{0}\right)^{a} \psi u \phi^{(a)}\left(x^{0}+\left(x-x^{0}\right) \theta(x)\right)\right]\left(\lambda \xi^{0}\right)= \\
& =F\left[\left(x-x^{0}\right)^{a} \psi u^{\xi} * F\left[\phi^{(a)}\left(x^{0}+\left(x-x^{0}\right) \theta(x)\right)\right]\left(\lambda \xi^{0}\right)=\int \Phi_{a}\left(x^{0}, \lambda \xi^{0}-\xi\right) \Psi_{a}\left(x^{0}, \xi\right) d \xi\right. \tag{2.2}
\end{align*}
$$

where

$$
\begin{aligned}
& \Phi_{a}\left(x^{0}, \xi\right)=F\left[\phi^{(a)}\left(x^{0}+\left(x-x^{0}\right) \theta(x)\right)\right](\xi), \\
& \Psi_{a}\left(x^{0}, \xi\right)=F\left[\left(x-x^{0}\right)^{a} \psi u\right](\xi) .
\end{aligned}
$$

Using $(2.6)^{9}$ and standard considerations it may be proved that (2.2) is an asymptotic expansion indeed. According to $(2.3)^{9}$

$$
\begin{align*}
& \left.\mathrm{e}^{\left.\mathrm{i} \lambda g^{\circ}\right)} \mathrm{F}\left[\left(\mathrm{x}-\mathrm{x}^{\circ}\right)^{a} \psi \mathrm{u}\right]\left(\lambda \xi^{\circ}\right)=\left(\frac{\mathrm{i}}{\lambda}\right)^{|a|} \partial_{\xi^{\circ}}\left(\mathrm{e}^{\mathrm{i} \mathrm{\lambda g}\left(\xi^{\circ}\right)} \mathrm{Fl} u \psi\right]\left(\lambda \xi^{\circ}\right)\right) \in \\
& \in\left(\frac{\mathrm{i}}{\lambda}\right)^{|a|} \partial_{\xi^{\circ}}^{a} \mathcal{G}\left[\mathrm{w}\left(\mathrm{~g}^{\prime}\left(\xi^{\circ}\right), \xi^{\circ}\right)\right. \tag{2.3}
\end{align*}
$$

From (2.2) and (2.3) we get (2.1) and this completes the proof.

Let $u$ be a distribution from the class $\mathscr{L}_{0}$ of which we know the singular spectrum and the symbol. How can we calculate the symbol of Au, where $A$ is a pseudodifferential operator? In order to answer this question we resort to a somewhat formal reasoning.

Let $A$ be a pseudodifferential operator from the class $\Psi{ }_{\rho}^{\mathrm{m}} \delta\left(\right.$ see ref. $\left.{ }^{/ 5 /}\right), 1-\rho \leqq \delta<\rho$ with symbol $\mathrm{a}(\mathrm{x}, \xi)$. Take $\left(\rho^{\circ} \delta, \xi^{\circ}\right) \in \Lambda$, where $\Lambda$ is a Lagrangian manifold, $\Lambda \partial$ sing spec $u$ and let $\psi \in 1_{0}^{\infty}\left(x^{0}\right)$. Then
$F[\psi A u]\left(\lambda \xi^{\circ}\right)=\int e^{-\mathrm{t} \lambda \mathrm{x} \xi^{\circ}} \psi(\mathrm{x}) \mathrm{Au}(\mathrm{x}) \mathrm{dx}=$
$=\int \mathrm{e}^{-\lambda \lambda \cdot \xi^{\circ}} \psi(x)\left[\iint \mathrm{e}^{1(x-y) \cdot \eta} \mathrm{a}(\mathrm{x}, \eta) \mathrm{u}(\mathrm{y}) \mathrm{dyd} \eta\right] \mathrm{dx}=$
$\approx \int\left[\sum_{a} \frac{1}{a!} \iint e^{-1(y-x) \cdot \eta} \quad \partial_{x}^{a} a\left(x^{0}, \eta\right)\left(x-x^{0}\right)^{a} \mathrm{e}^{-\mathrm{i} \lambda_{x} \xi^{\circ}} \psi(x) d x d \eta\right] u(y) d y=$
$=(-1)^{\mathrm{n}} \sum_{a} \frac{1}{a!} \int A_{\alpha}^{o}\left[e^{-\lambda \lambda x \xi^{\circ}}\left(x-x^{0}\right) \psi(x)\right](y) u(y) d y$,
where $A_{a}^{\circ}$ is a pseudodifferential operator, whose symbol $\sigma\left[A_{a}^{\circ}\right](y, \eta)=\partial^{a} a\left(x^{\circ}, \eta\right)$ doesn't depend on $y$ and belongs to class $S_{\rho, 0}$ (see, ref. ${ }^{15 /)}$. We apply a wel1-known asymptotic expansion formula for pseudodifferential operators (see ref. ${ }^{15 /}$ ) and obtain

$$
\begin{equation*}
:\left.A_{\alpha}^{\circ}\left[\mathrm{e}^{-\lambda \mathrm{x} \xi^{\circ}}\left(\mathrm{x}-\mathrm{x}^{0}\right) \psi(\mathrm{x})\right](\mathrm{y}) \approx \mathrm{e}^{-\mathrm{i} \lambda \mathrm{y} \xi^{\circ}} \sum_{\beta} \frac{1}{\beta!} \partial_{\eta}^{\beta} \partial_{\mathrm{x}}^{\alpha} \mathrm{a}\left(\mathrm{x}^{\circ},-\lambda \xi^{\circ}\right) \mathrm{D}_{\mathrm{z}}^{\beta}\left[\left(\mathrm{z}-\mathrm{k}^{\circ}\right)^{\alpha} \psi(\mathrm{z})\right]\right|_{\mathrm{z}=\mathrm{y}} \tag{2.5}
\end{equation*}
$$

By the Leibniz formula

$$
\begin{align*}
& \mathrm{D}_{\mathrm{z}}^{\beta}\left[\left(\mathrm{z}-\mathrm{x}^{0}\right)^{\alpha} \psi(\mathrm{z})\right]=\sum_{\gamma \leqq \beta}\binom{\beta}{\gamma} \mathrm{D}_{\mathrm{z}}^{\beta-\gamma}\left(\mathrm{z}-\mathrm{x}^{0}\right)^{\alpha} \mathrm{D}_{\mathrm{z}}^{\gamma} \psi(\mathrm{z})=  \tag{2.6}\\
& =\Sigma_{\gamma \leqq \beta}\binom{\beta}{\gamma} \frac{\alpha!}{(a-\beta+\gamma)!}\left(\mathrm{z}-\mathrm{x}^{0}\right)^{\alpha-\beta+\gamma} \mathrm{D}_{\mathrm{z}}^{\gamma} \psi(\mathrm{z}) .
\end{align*}
$$

We substitute (2.6) in (2.5) and then (2.5) in (2.4)

$$
\begin{aligned}
& \mathrm{F}[\psi \mathrm{Au}]\left(\lambda \xi^{\circ}\right)=(-1)^{\mathrm{n}} \sum_{\alpha} \frac{1}{\alpha!} \sum_{\beta} \frac{1}{\beta!} \partial_{\eta}^{\beta} \partial_{\mathrm{x}}^{a} \mathrm{a}\left(\mathrm{x}^{\mathrm{o}},-\lambda \xi^{0}\right) \times \\
& \times \sum_{\gamma \leq} \frac{a!\beta!}{(\alpha-\beta+\gamma)!(\beta-\gamma)!\gamma!} \int \mathrm{e}^{-\mathrm{i} \lambda \mathrm{y} \xi^{\mathbf{o}}}\left(\mathrm{y}-\mathrm{x}^{0}\right)^{\alpha-\beta+\gamma} \mathrm{D}_{\mathrm{y}}^{\gamma} \psi(\mathrm{y}), \mathrm{u}(\mathrm{y}) \mathrm{dy}= \\
& =(-1)^{\mathrm{n}} \sum_{\alpha} \sum_{\beta} \sum_{\gamma \leqq} \frac{1}{\beta(\alpha-\beta+\gamma)!(\beta-\gamma)!\gamma!} \partial_{\eta}^{\beta} \partial_{\mathrm{x}}^{\alpha} \mathrm{a}\left(\mathrm{x}^{\mathrm{o}},-\lambda \xi^{0}\right) \times \\
& \times \int \mathrm{e}^{-\mathrm{i} \lambda \mathrm{y} \zeta^{\infty}}\left(\mathrm{y}-\mathrm{x}^{\circ}\right)^{a-\beta+\gamma} \mathrm{u}(\mathrm{y}) \mathrm{D}_{\mathrm{y}}^{\gamma} \psi(\mathrm{y}) \mathrm{dy} .
\end{aligned}
$$

For $\gamma>0$ the 1 ast integral equals zero asymptotically when $\lambda \rightarrow \infty$ since $\mathrm{D}_{y}^{\gamma} \psi(\mathrm{y})=0$ in a neighbourhood of $\mathrm{x}^{\circ}$. For $\gamma=0$ we get in view of $(2,4)^{9}$

$$
\begin{align*}
& \int \mathrm{e}^{-\mathrm{i} \lambda y \xi^{\circ}}\left(\mathrm{y}-\mathrm{x}^{0}\right)^{\alpha-\beta} u(y) \psi(y) \mathrm{dy} \approx \mathrm{e}^{-i \lambda x^{\circ} \xi^{\circ}} \mathrm{s}\left[\left(\mathrm{y}-\mathrm{x}^{\circ}\right)^{\alpha-\beta} \mathrm{u}(\mathrm{y})\right]\left(\mathrm{x}^{\circ}, \xi^{\circ} ; \lambda\right) \approx \\
& =\left(\frac { 1 } { \lambda } | ^ { | \alpha - \beta | } \partial _ { \xi ^ { \circ } } ^ { \alpha - \beta } \mathrm { s } \left[\mathrm{u}\left(\mathrm{~g}^{\prime}\left(\xi^{\circ}\right), \xi^{\circ} ; \lambda\right)\right.\right. \tag{2.8}
\end{align*}
$$

Substituting (2.8) into (2.7) we get

$$
s[A u]\left(x^{\circ}, \xi^{\circ} ; \lambda\right) \approx e^{i \lambda x^{\circ}, \xi^{\circ}} F[\psi A u]\left(\lambda \xi^{\circ}\right)=
$$

$$
\begin{equation*}
\Rightarrow(-1)^{\mathrm{n}} \sum_{\alpha} \sum_{\beta} \frac{1}{(\alpha-\beta)!\beta!}\left(\frac{\mathrm{i}}{\lambda}\right)^{|\alpha-\beta|} \partial_{\eta}^{\beta} \partial_{\mathrm{x}}^{\alpha} \mathrm{a}\left(\mathrm{x}^{0},-\lambda \xi^{\circ}\right) \partial_{\xi^{\circ}}^{\alpha-\beta} \mathrm{s}[\mathrm{u}]\left(\mathrm{g}^{\circ}\left(\xi^{\circ}\right), \xi^{\circ} ; \lambda\right) \tag{2.9}
\end{equation*}
$$

We won't justify the above reasoning, but only state the final result.

## 3. HINTS ON APPLICATIONS

We will not consider any of the sophisticated and important applications of "microlocal" methods in different fields of mathematics and physics, which are the subject of an overwhelming torrent of papers during the last years. In this section we only want to give a rough idea of how the previous concepts might be used for exploring the singularities of solutions to partial differential equations.

A most important problem in analysis is to determine the singularities of the kernels of different operators connected with differential equation problems - first of all the kernels of the differential operators themselves and of their inverses, i.e., the singularities of their fundamental solutions. We already saw that the singularities of the kernel of a pseudodifferential operator are characterized by the symbol of the operator: What about the inverse operators? Consider an example.

Let $P$ be a hyperbolic differential operator of second order. We are interested in the singularities of the Green function of the Cauchy problem for $P$. If we denote $x=\left(x_{0}, x_{1}, \ldots, x_{n}\right)=$ $=\left(x_{0}, x^{\prime}\right)$,
then the Green function $G\left(x ; y^{\prime}\right)$ satisfies

$$
\begin{align*}
& P_{x} G\left(x ; y^{\prime}\right)=0,  \tag{3.1}\\
& \left.G\left(x ; y^{\prime}\right)\right|_{x_{0}}=0=0, \\
& \left.\frac{\partial G}{\partial x_{0}}\left(x ; y^{\prime}\right)\right|_{x_{0}}=0=\delta\left(x^{\prime}-y^{\prime}\right) .
\end{align*}
$$

We want to determine the singularity of $G\left(x ; y^{\prime}\right)$ considered as a distribution with respect to $y^{\prime}$ depending on the parameter $x$.

Suffice it to calculate the local Fourier asymptotic asym $\mathrm{Fy}^{\prime}\left[\mathrm{G}\left(\mathrm{x} ; \mathrm{y}^{\prime}\right) \phi\left(\mathrm{y}^{\prime}\right)\right]\left(\lambda \eta^{\prime}\right)$. We multiply by $\phi\left(\mathrm{y}^{\prime}\right)$ in
and carry out the Fourier transform with respect to $\mathrm{y}^{\prime}$

$$
\begin{align*}
& P_{z} F_{y^{\prime}}[G \phi]\left(\lambda \eta^{\prime}\right)=0, \\
& \left.F_{y^{\prime}}[G \phi]\left(\lambda \eta^{\prime}\right)\right|_{x_{0}}=0=0,  \tag{3.2}\\
& \left.\frac{\partial}{\partial x_{0}} F_{y^{\prime}}[G \phi]\left(\lambda \eta^{\prime}\right)\right|_{x_{0}}=0=e^{-1 \lambda x^{\prime} \eta^{\prime}} \phi\left(x^{\prime}\right) .
\end{align*}
$$

If asym.F. $[$ ' $\mathrm{G} \phi]\left(\lambda \eta^{\prime}\right)$ may be differentiated with respect to $x$ then (3.2) implies

$$
\begin{align*}
& P_{x} \operatorname{asym} F_{y},[G \phi]\left(\lambda \eta^{\prime}\right) \approx 0, \\
& \left.\quad \operatorname{asym} F_{y^{\prime}} \cdot[\mathrm{G} \phi]\left(\lambda \eta^{\prime}\right)\right|_{x_{0}}=0=0,  \tag{3.3}\\
& \left.\frac{\partial}{\partial x^{0}} \operatorname{asym}_{y^{\prime}} \cdot[\mathrm{G} \phi](\lambda \eta)\right|_{x_{0}}=0 \approx e^{-i \lambda x^{\prime} \eta^{\prime}} \phi\left(x^{\prime}\right) .
\end{align*}
$$

But (3.3) means that asym $\mathrm{F}_{\mathrm{y}} \cdot[\mathrm{C} \phi]\left(\lambda \eta^{\prime}\right)$ is an asymptotic solution of the oscillatory Cauchy problem. We state this result in a slightly more general form.

Theorem 3.1. Let $P$ be a differential operator and $G(x, y)$ be the Green function of the Cauchy problem for P. If
asym $F_{y}[G(z ; y) \phi(y)]\left(\lambda \eta^{\prime}\right)$ is differentiable with respect to $x$ then it is an asymptotic solution of an oscillatory Cauchy problem.

This is the reason why asymptotic solutions characterize the singularities of the Green function. Obviously if we know the singularities of the Green function we can deduce information about the singularities of the solutions. Using the results of the previous section we can develop a technique for calculation of the singularities of solutions.

Let $P$ be a differential operator and let $u \in \mathscr{L}_{0}$ be a solution of the $\mathrm{Pu}=0$. What relations will satisfy the singularities of $u$, i.e., the singular spectrum and the symbol of $u$ ?

Theorem 3.2. Let $P=|a|<\mathrm{m}{ }^{2}(\mathrm{x}) \mathrm{D}_{1}^{a}$ be a differential operator.Denote

$$
P_{0}(x, \xi)=\mid \sum_{\mid=\mathrm{m}} a(\mathrm{x}) \xi^{a}, \quad \mathrm{P}_{1}(\mathrm{x}, \xi)=\sum_{|a|=m+1} a_{a}(\mathrm{x}) \xi^{a} .
$$

Let $u \in \mathscr{L}_{0}$, singspec $u \subset \Lambda=\left\{(x, \xi): x=g^{\circ}(\xi)\right\}$. Suppose that the symbol of $u$ has the asymptotic expansion
$s[u](x, \xi ; \lambda)=\frac{1}{\lambda^{\nu}} \sum_{k=0}^{\infty} \frac{r_{k}(x, \xi)}{\lambda^{k}}, \quad(x, \xi) \in \Lambda$.

2. $r_{0}, r_{1}, \ldots$ satisfy the following relations

$$
\begin{align*}
& H_{P_{0}} r_{0}(x, \xi)+\left(P_{1}-\frac{1}{2 i} \frac{\partial^{2} P_{0}}{\partial x_{j} \partial \xi_{j}}\right) r_{0}(x, \xi)=0,  \tag{3.4}\\
& H_{P_{0}} r_{1}(x, \xi)+\left(P-\frac{1}{2 i} \frac{\partial^{2} P_{0}}{\partial x_{j} \partial \xi_{j}}\right) r_{1}(x, \xi)+B r_{2}(x, \xi)=0,
\end{align*}
$$

where $H_{P_{0}}$ is the Hamiltonian field of $P_{0}$ and $B$ is a differential operator of second order.

The proof of this theorem is a direct application of the results of the previous section.

## APPENDIX 1

## ASYMPTOTICS

We recall here some well known notions and facts about asymptotics. The proofs may be found in refs. ${ }^{14,7 /}$.

Let $f$ and $g$ be two functions $R_{+} \rightarrow V$, where $V$ is a topological vector space. We say that $f$ and $g$ are $V$-asymptotically equal at infinity (and note $f \underset{\lambda \rightarrow \infty}{v} g$ ) iff for every $N>0$ we have

$$
\lim \lambda^{N}(f(\lambda)-g(\lambda))=0
$$

This is obviously an equivalence relation in the set of functions $R \underset{+}{\square} V$ and we call the equivalence classes $V$-asymptotics (or $\hat{V}$-asymptotic classes). Denote by $G(V)$ the set of all V -asymptotics. We note by
asym f
the ${ }^{\lambda} \overrightarrow{c l}^{\infty}{ }^{\infty}$ ss to which $f$ belongs and call it $V$-asymptotic of $f$. Let the topology of $V$ is determined by a set of functionals $\left\{F_{i} \|_{1 \in 1}\right.$ in the sense that
$\mathrm{u}_{\mathrm{n}} \xrightarrow{\stackrel{u}{u}} \Longleftrightarrow F_{\mathrm{i}}\left(\mathrm{u}_{\mathrm{n}}-\mathrm{u}\right) \xrightarrow{R^{1}} 0, \quad i \in I$.

$$
f: R_{+} \rightarrow V, g: R_{+} \rightarrow V \quad \text { we say that }
$$

asymf \& asymg
iff

$$
F_{i}(f(\lambda)) \underset{\lambda \rightarrow \infty}{=} o\left(F_{i}(g(\lambda))\right), \quad i \in I
$$

We define operations with asymptotics: $a_{1}+a_{2}, a_{1} a_{2}$, etc.
A sequence of asymptotics $\left\{a_{n}=\operatorname{asym} \phi_{n}\right\}, n=1,2, \ldots$ is called asymptotic scale iff $a_{n+1}\left\langle a_{n}, n=1,2, \ldots\right.$ - and for each $N$ there is an $n$ such that $a_{n} \prec$ asym $(\lambda-N)$

Example: $a_{1}=\operatorname{asym} \lambda^{m}, a_{2}=\operatorname{asym} \lambda^{m-\rho} \quad, \quad a_{3}=\operatorname{asym} \lambda^{m-2 \rho}, \ldots$
If $\left\{a_{n}\right\}$ is an asymptotic scale, then
$a_{1}+a_{2}+\ldots$
is called an asymptotic series.
Let $\left\{a_{n}\right\}$ be an asymptotic scale and $a$ be an asymptotic such that
$a-\sum_{n=1}^{N} a_{n}\left\langle a_{N} \quad\right.$ for each $N$.
In that case we say that $a$ is the sum of the asymptotic series $a_{1}+a_{2}+\ldots$ or that $a_{1}+a_{2}+\ldots$ is an asymptotic expansion of $a$ and note
$a \approx a_{1}+a_{2}+\ldots$.
Theorem 1. Any asymptotic series has a sum and this sum is unique.

Theorem 2. Let $a_{1}, a_{2}, \ldots$ be an asymptotic scale. An asymptotic a may have not more than one asymptotic expansion of the form
$a \approx c_{1} a_{1}+c_{2} a_{2}+\ldots$,
where $c_{1,} c_{2}, \ldots$ are numbers.

## APPENDIX 2

## LAGRANGIAN MANIFOLDS

Some elementary facts concerning Lagrangian manifolds will be enumerated. The proofs may be found in refs. ${ }^{1,3,4,5,8}$,

Consider $R^{2 \mathrm{n}}=R_{x}^{\mathrm{n}} \times R_{\xi}^{\mathrm{g}}$ provided with the differentiai 2form
$\omega=d \xi_{1^{\wedge}} \wedge \mathrm{dx}_{1}+\ldots+\mathrm{d} \xi_{n} \wedge d x_{*}$,
which is non-degenerated.
An $n$-dimensional manifold $\Lambda \subset R^{2 n}$ is called Lagrangian manifold iff $\left.\omega\right|_{\Lambda}=0$.

A Lagrangian manifold $\Lambda$ may be represented locally by a single function of $n$ variables. For instance if $\Lambda$ projects one-to-one on the $\xi$-space, then there exist a function $g\left(\xi_{1}, \ldots, \xi_{\mathrm{n}}\right)$ so that $\Lambda$ is represented by the equations

$$
\begin{align*}
& x_{1}=\frac{\partial g}{\partial \xi},  \tag{1}\\
& 1
\end{align*}
$$

$x_{n}=\frac{\partial g}{\partial \xi_{n}}$.
In the general case the following theorem holds:

Theorem 1. Let $\Lambda$ be a Lagrangian manifold and $\left(x^{\circ}, \xi^{\circ}\right) \in \Lambda$.There exists a neighbourhood $U \subset R^{2 n}$ of $\left(x^{\circ}, \xi^{\circ}\right)$ and a function. $g\left(\xi_{i}, \ldots, \xi_{i_{k}} X_{j_{1}}, \ldots, x_{j_{n-k}}\right)$ so that $\Lambda \cap U$ is represented by the

## equations

$$
\mathbf{x}_{i_{1}}=\frac{\partial g}{\partial \xi_{i_{1}}}
$$

$$
\begin{equation*}
\mathbf{x}_{\mathbf{i}_{k}}=\frac{\partial \mathrm{g}}{\partial \xi_{\mathbf{i}_{k}}} \tag{2}
\end{equation*}
$$

$$
\xi_{j_{1}}=-\frac{\partial g}{\partial x_{j_{1}}}
$$

$$
\xi_{j_{n-k}}=-\frac{\partial g}{\partial x_{j_{n-k}}}
$$

where ( $i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{n-k}$ )
is a permutation of $(1,2, \ldots, n)$.
The next theorem enables us to reduce the general case (1) to the special case (1) by a suitable change of local coordinates.

Theorem 2. Let $\Lambda$ be a Lagrangian manifold and ( $x^{0}, \xi^{\circ}$ ) $\in \Lambda$. There exists a neighbourhood $U \subset R^{2 n}$ of $\left(x^{\circ}, \xi^{\circ}\right)$ and a change
of the $x$-coordinates $x=\kappa(y)$, so that if we perform in $U$ the canonical change

## $\mathrm{x}=\kappa(\mathrm{y})$,

$\mathrm{D} \kappa^{*} \xi=\eta$,
then $\Lambda \cap U$ has a one-to-one projection on the $\eta$-space.
There is a close connection between Lagrangian manifolds and Hamiltonian systems in classical mechanics. Roughly speaking any Lagrangian manifold may be foliated into Hamiltonian paths.

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[^0]:    * This paper is the continuation of ref. ${ }^{/ 9 /}$ and (k) ${ }^{8}$ means formula (k) from ref. ${ }^{1 / \%}$.

