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PARTIAL RESTORATION OF THE CHIRAL SYMMETRY IN THE GENERALIZED SKYRME MODEL

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In the low-energy region, where the hadronic mass spectrum is formed, quantum chromodynamics is characterized by spontaneous breaking of the chiral and conformal symmetry. This leads to generation of quark and gluon condensates. The effective lagrangian describing strong interactions in this region must express these pecularities of QCD in terms of the effective degrees of freedom. The latter are bosonic fields. In this approach the observable mesons are treated as small fluctuations above chirally nonsymmetric vacuum while the topologically nontrivial solitons correspond to baryons.

The aim of our work is to investigate some static nucleon properties in the framework of the model obtained by the simultaneous conformal and chiral bosonization method¹¹. The model we consider is a generalization of the well-known Skyrme model ²/₂ which takes into account the scalar field G(x) interacting with the chiral field U(x).

The model is defined by following expression:

$$\begin{aligned} \exists_{eff}(U,\sigma) &= \frac{F_{\pi}^{2}}{4} \cdot e^{-2\sigma} \operatorname{Tr} \partial_{\mu} U \partial^{\mu} U^{\dagger} + N_{f} \frac{F_{\pi}^{2}}{4} e^{-2\sigma} (\partial_{\mu} \sigma)^{2} \\ &+ \frac{1}{128\pi^{2}} \operatorname{Tr} \left[\partial_{\mu} U \cdot U^{\dagger}, \partial_{\nu} U \cdot U^{\dagger} \right]^{2} \\ &- C_{g} \frac{N_{f}}{48} \left[e^{-4\sigma} + \frac{4}{\epsilon} \left(1 - e^{-\epsilon\sigma} \right) \right] , \end{aligned}$$

where $\int_{\mathcal{T}}$ is the pion decay constant (93 MeV) and $N_{\mathcal{F}}$ is the number of flavours. The magnitude of the gluon condensate parameter $C_{\mathcal{G}} = \langle G_{\mu\nu}^2 \rangle$ lies within the interval of admissible values: $[(300-400) \text{ LeV}]^4$. The first and second terms are kinetic energies of the chiral and scalar fields, the third corresponds to the Skyrme selfinteraction. The effective potential for a scalar field given by the last term of the Lagrangian is the extrapolation of the low-energy potential^{/3/} into the high energy region. In this extrapolation, the one-loop approximation to the Gell-Mann-Low QCD β -function is used. The parameter $\hat{\mathcal{E}}$ depends on the flavour number: $\hat{\mathcal{E}} = 8N_{\mathcal{F}}/(33-2N_{\mathcal{F}})$.

Passing to baryonic sector we have to make some assumption about the form of the static chiral and scalar fields. In particular,

for the chiral field we choose the Skyrme-Witten ansatz $U(x) = = \exp \{i\vec{T}\vec{n}, F(r)\}$ (where $\vec{n} = \vec{r}/r$) and propose spherical symmetry for the scalar field $\sigma(x)$. Then for the mass functional we have

$$M = M_2 + M_4 + V.$$

(2)

Here

$$M_{2} = 4\pi (F_{\pi}/e) \cdot \int_{0}^{\infty} dx \left\{ \frac{N_{4}}{4} x^{2} (\rho')^{2} + \rho^{2} \left[\frac{x^{2} (F')^{2}}{2} + \sin^{2} F \right] \right\}, \quad (3)$$

$$M_{4} = 4\pi (F_{\pi}/e) \cdot \int_{0}^{\infty} dx \left\{ \frac{\sin^{4} F}{2x^{2}} + \sin^{2} F \cdot (F')^{2} \right\}, \quad (4)$$

$$V = 4\pi (F_{\pi}/e) \cdot Deff \cdot \int_{0}^{\infty} dx \cdot x^{2} \left[\rho^{4} - 1 + \frac{4}{\epsilon} (1 - \rho^{\epsilon}) \right], \quad (4)$$

$$\rho = e^{-G(x)}. \quad (5)$$

In eqs. (3)-(5), the dimensionless variable $x = F_{\pi} \cdot e \cdot r$ includes the Skyrme parameter e equal to 2π in accordance with eq.(1). For the potential factor D_{eff} we have $D_{eff} = C_g N_f / 48 e^2 F_{\pi}^4$. The mass functional leads to the following system of equations for F(x) and $\rho(x)$:

$$F''[g^{2}x^{2}+2Sin^{2}F]+2F'x[xgg'+g^{2}]+Sin(2F)(F')^{2}$$

$$-g^{2}Sin(2F)-\frac{Sin^{2}F}{x^{2}}Sin(2F)=0,$$

$$\frac{N_{4}}{2}x[xg''+2g']-2g[\frac{x^{2}(F')^{2}}{2}+Sin^{2}F]-4D_{eff}[g^{2}-g^{E-1}]x^{2}=0.$$
(6)

where prime corresponds to the derivative with respect to \mathfrak{L} . According to the virial theorem, the solutions should satisfy the following condition: $\mathcal{M}_{4} - \mathcal{M}_{2} - 3V=0$, that we use to control the accuracy of the numerical calculations. For small \mathfrak{L} we have $\mathbf{F}(\mathbf{x}) \sim \mathcal{T} \mathcal{N} - d\mathbf{x}$, $\mathcal{P}(\mathbf{x}) \sim \mathcal{P}(0) + \beta \mathfrak{X}^{2}$ and for asymptotically large \mathbf{x} : $F(\mathbf{x}) \sim \alpha / x^{2}$, $\mathcal{P}(\infty) \sim 4 - b / \mathbf{x}^{6}$. There are nontrivial relations between the numerical coefficients \mathcal{L} and β , \mathcal{A} and b. The boundary conditions ensure a finite mass functional for a given value of the topological charge = N. We may quantize the rotational degrees of freedom by means of collective variable method $^{\prime 4\prime}$. As a result, we get the following expression for the nucleon mass $M_{\rm A\prime}$

$$\mathcal{M}_{N} = \mathcal{M} + \frac{3}{81} ,$$

where the rotational moment of inertia is now

$$I = \frac{8\pi}{3} \left(F_{\pi} e^{3} \right)^{-1} \int dx \sin^{2} F \left[\rho^{2} x^{2} + (F')^{2} x^{2} + \sin^{2} F \right].$$

Some numerical results are given in the table, where also the mean square root radius of the baryon charge distribution $\langle r_b^2 \rangle^{1/2}$ is shown.

Table. Our numerical results in the generalized Skyrme model. In the last column the results obtained in the original Skyrme model are given for the same values of F_{-} , e_{-} .

	-		π	
	(300 MeV) ⁴	(300 MeV) ⁴	Skyrme	
N _f	3	2	2	
F _T	. 93 Me∛	93 MeV	93 MeV	
e.	250	2 X	251	
ρ(0)	0.29	0.22	1	
M	867 MeV	827 MeV	1087 MeV	
$\langle r_{\mu}^{2} \rangle^{\gamma_{2}}$	0.37 Fm	0.38 Fm	0.36 Fm	
MN	1 0 72 MeV	1033 MeV	126 0 MeV	

One can see a partial restoration of the chiral symmetry which appears as the big deviation of $\rho(0)$ from its asymptotical value 1. (See the Figure).





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We should like to point out that the classical and rotational component of the nucleon mass are much smaller here as compared to the original Skyrme model.

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Received by Publishing Department on July 18, 1988. Николаев В.А., Новожилов В.Ю., Ткачев О.Г. Е4-88-536 Частичное восстановление киральной симметрии в обобщенной модели Скирма

Рассмотрена обобщенная модель Скирма, которая учитывает основные характерные свойства КХД /образование кваркового и глюонного конденсатов, благодаря нарушению киральной и конформной симметрий/. Показано, что эффективный потенциал для скалярных мезонов, полученный методом одновременной киральной и конформной бозонизации в КХД, существенно определяет статические свойства топологических солитонов. Теоретические значения барионных наблюдаемых значительно улучшаются при экспериментальном значении пионной постоянной F_{π} в сравнении с оригинальной моделью Скирма.

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Nikolaev V.A., Novozhilov V.Yu., Tkachev O.G. E4-88-536 Partial Restoration of the Chiral Symmetry in the Generalized Skyrme Model

The improved Skyrme model that takes into account the main features of QCD (the formation of quark and gluon condensates due to the breaking of chiral and conformal symmetries) is considered. The effective potential for scalar mesons that was derived within the framework of the joint chiral and conformal bosonization method in QCD erucially defines the static properties of the topological solitons. The baryon observables improve considerably for the experimental value of the pion decay constant F_{π} in comparison with the original Skyrme model results.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

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