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V.A.Kuzmin, V.G.Soloviev

EFFECT OF THE PARTICLE-PARTICLE INTERACTION ON THE GAMOW-TELLER β^+ -DECAY IN SPHERICAL NUCLEI

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Кузьмин В.А., Соловьев В.Г. Е4-87-831 Влияние взаимодействия в канале частица-частица на гамов-теллеровский *В*⁺-распад в сферических ядрах

В приближении хаотических фаз с учетом остаточных сепарабельных взаимодействий в каналах частица-частица и частица-дырка вычислены ft-величины гамов-теллеровских β^{+} -распадов ядер ¹⁵² Yb, ¹⁵⁰ Er, ^{148,146} Dy и ⁹⁶ Pd. При фиксированном отношении констант частично-частичного и частично-дырочного взаимодействий получено согласие с экспериментально определенными значениями ft для $|G_A/G_V| = 1$. Показано, что из расчетов вероятностей β -распадов ядер, удаленных от полосы стабильности, нельзя сделать вывод о величине перенормировки аксиально-векторной константы G_A слабого взаимодействия.

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In RPA approximation, taking account of residual separable interactions in the particle-hole and particleparticle channels, ft-values for the Gamow-Teller β^{+} -decays of ¹⁵² Yb, ¹⁵⁰ Er, ^{148,146} Dy and ⁹⁶Pd are calculated. At a fixed constant ratio of the particle-particle and particle-hole interaction agreement with experimentally determined ft-values for $|G_A/G_V| = 1$ is obtained. It is shown that from the calculations of probabilities of β -decays of nuclei far from stability one cannot conclude on the renormalisation of the axial-vector constant G_A of the weak interaction.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

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Advances in describing some characteristics of collective lowlying states within the interacting boson model^{/1/}, that takes into account particle-particle interactions, entailed studies of the role of particle-particle interactions. Taking interactions in the particlehole and particle-particle channels into account simultaneously allows one to fit the description of the two-neutrino double ß-decay with the experiment^{/2/}. This letter is devoted to the description of the Gamow-Teller ß⁺-decay of even spherical nuclei in the RPA taking into account interactions between quasiparticles in the particle-particle and particle-hole channels.

The QPNM Hamiltonian³ contains particle-hole and particle-particle interactions between quasiparticles. From the general formulae of the model one can easily derive equations for the energies and wave functions of the Gamow-Teller 1⁺ states in the RPA with particle-particle and particle-hole charge-exchange interactions with the constants G_{I}^{00} and \mathcal{B}_{I}^{01} , respectively. The creation operator of the charge-exchange phonon is written in the form

 $\Omega_{i}^{+} = \sum_{i=l_{n}} \left\{ \Psi_{jpj_{n}}^{i} A^{(j_{p}j_{n};10)} + \Psi_{jpj_{n}}^{i} A^{(j_{p}j_{n};10)} \right\},$

where

 $A^{\dagger}(j_{p}j_{n};10) = \sum \langle j_{p}m_{p}j_{n}m_{n}|10\rangle d^{\dagger}_{j_{p}}m_{p}d^{\dagger}_{j_{n}}m_{n}$

Here d_{jm} is the quasiparticle creation operator, $j_{p}m_{p}(j_{n}m_{n})$ are quantum numbers of proton (neutron) single-particle states, $\ell = 1, 2, 3, ...$ is the root number of the relevant secular equation.

The matrix element of the \mathfrak{B}^+ -decay of the ground state of the doubly even nucleus with the wave function $\Psi_{\mathbf{0}}$ into the one-phonon 1⁺state of the doubly odd nucleus with the wave function $\Omega_{\mathbf{1}}^+ \Psi_{\mathbf{0}}^-$ is

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 $(\Psi_{0}^{*}\Omega_{i}H_{\beta}^{i}\Psi_{0}) = \sum_{i_{0}} \langle j_{n}\|\Gamma_{\beta}\|j_{p}\rangle(\Psi_{j_{0}j_{n}}^{i_{0}}J_{j_{0}}^{p}u_{j_{n}} + \varphi_{j_{0}j_{n}}^{i_{0}}u_{j_{0}}^{p}J_{j_{n}}^{n}),$

where $\langle in | f_{\alpha} | i_{\beta} \rangle$ is the single-particle Gamow-Teller matrix element, \mathcal{U}_{j} and \mathcal{V}_{i} are the coefficients of the Bogolubov canonical transformation. The inclusion of the particle-particle interaction alongside with the particle-hole interaction leads to the increase of \mathcal{U}_{j} for the lowest states. Since the signs of the functions \mathcal{U}_{j} and \mathcal{U}_{j} are opposite, the increase in the absolute value of the particle-particle interaction constant leads to the suppression of the \mathbb{B}^+ transition probabilities.

The results of calculations of $\log \hat{ft}$ for the β^+ transitions to the lowest 1^+ states of neutron-deficit nuclei calculated in the RPA and the relevant experimental sum values of log ft, defined as

 $(\widetilde{\mathfrak{ft}})^{-1} = \sum_{\kappa} (\mathfrak{ft})_{\kappa}^{-1}$

of refs. $^{5-9/}$ are shown in the table. The same single-particle energies and wave functions of the Saxon-Woods potential and the pairing constants were used in the calculations as in ref. $^{4/}$. The value of $|\mathfrak{A}_{j}^{0}, \mathcal{A}_{j}|$ is 1.5 times as large as the value used in ref. $^{4/}$ when studying nuclei of the stability zone. The constants of the particle-particle interaction $\mathcal{G}_{j}^{0} = -\theta_{0} 2 \, \mathfrak{A}_{j}^{0}$ were used in the calculations. With increasing constant one can suppress twice the strength of β -transitions to the low-lying states without violating the applicability conditions for the RPA.

Table. Values of $\log \frac{2}{ft}$ for B^+ -transitions from $O_{g.s.}^+$ to 1^+ states

		1 × 1 × 1		10 A 19 A
ß ⁺ -transition	log ft (experimenț)	$\frac{\log \widetilde{ft} (calculus)}{G_{I}^{ol} - 0.2 \varkappa_{I}^{ol}}$	$G_{i}^{\prime \prime}=0$	· •
$152_{\rm Yb} - 152_{\rm Tm}$	3.4	3.5	3.1	
$150_{\rm Er} - 150_{\rm Ho}$	3.6	3.5	3.2	· ·
148 Dy - 148 Tb		3.7	3.4	
$\frac{146}{Dy} - \frac{146}{Tb}$	3.8	3.8	3.3	
⁹⁶ Pd - ⁹⁶ Rh	3.3	3.3	3.0	×

In more detail we consider the results of calculations of β^+ -decay in ¹⁴⁸Dy. In comparison with the independent particle model the

inclusion of the particle-hole interaction results in a 2.65 decrease of the total β^+ transition strength. The inclusion of the particle-particle interaction gives an additional twice suppression. At $\mathcal{C}_{I}^{eq} = -\Omega_{2}I \mathscr{L}_{I}^{eq}$ we get log ft=3.9 that coincides with the experimental value. With increasing \mathcal{C}_{I}^{eq} the total β -transition strength decreases on retention of the relevant sum rule.

We have calculated the matrix elements of the two-neutrino double B-decay for ^{128,130}Te at $G_{f}^{0\ell} = -0.2 \, \varkappa_{f}^{0\ell}$ that are in agreement with ref.^{/2/} and the experimental value. The necessity of taking into account the interaction in the particle-particle channel to suppress the strength of the two-neutrino double β -decay is confirmed in ref.^{/10/}.

In describing ß-decays of nuclei far from stability, the axialvector constant G_A of a weak interaction is renormalised. In order to obtain an agreement with the experimental ft-values, it has been assumed in refs.^{5,6/} that $|G_A/G_V| = 0.6-0.8$ and in ref.^{11/} that $|G_A/G_V| =$ = 0.7-1.0. Our calculations are performed with $|G_A/G_V| = 1$. The ft-values, that are in agreement with the experiment, have been obtained for $|G_A/G_V| = 1.25$ with increasing G_4^{04} by 3%. Therefore, one cannot conclude on the renormalisation of the axial-vector constant G_A of the weak interaction from the calculations of β^+ -decays of nuclei far from stability.

References:

- Arima A., Iachello F., Phys.Rev.Lett., 1975, <u>35</u>, 1069; Ann.Phys. 1976, <u>99</u>, 253.
- 2. Vogel P., Zirnbauer M.R., Phys.Rev.Lett., 1986, 57, 3148.
- 3. Soloviev V.G. Prog.Part.Nucl.Physics, 1987, 19, 107.
- 4. Kuzmin V.A., Soloviev V.G., J.Phys.G: Nucl.Phys. 1984, 10, 1507.
- 5. Алхазов Г.Д., Ганбаатар И., Громов К.И. и др. Ядерная физика, 1984, <u>40</u>, 554.

Gromov K.Ya. In: "Intern.Symp. IN-BELM Nuclear Spectroscopy, ed. Dambradi Z., Fenyes T., Akademiai Kiado, 1984, Budapest, p. 269.

- 6. Алхазов Г.Д., Быков А.А., Витман В.Д. и др. Ядерная физика, 1985, <u>42</u>, 1313.
- 7. Rykaczewski K., Grant I.S., Kirchner R. et al. Z.Phys.A, 1985, <u>322</u>, 263.

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 Kleinheinz P., Zuber K., Canci C. et al., Phys.Rev.Lett., 1985, <u>55</u>, 2664.

9. Zuber K., Liang C.F., Paris P. et al. Z.Phys.A, 1987, 327, 357.

10. Civitarese O., Faessler A., Tomoda T. Phys.Lett.B, 1987, 194, 11.

 Alkhazov G.D., Artamonov S.A., Isakov V.I. et al. Preprint LNPI, N. 1305, 1987, Leningrad.

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