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Yu.A.Alexandrov

## ON POLARIZABILITY OF THE NEUTRON AND ITS MEAN SQUARE RADIUS

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From ref.<sup>/1/</sup> follows the most exact for the present value for the mean square radius of charge distribution in the neutron, i.e.,

$$\langle r_{\rm E}^2 \rangle = 6 \left( \frac{\partial G_{\rm E}}{\partial q^2} \right)_{q^2=0} = 3 \frac{\hbar^2}{{\rm Me}^2} a_{\rm ne}, \qquad (1)$$

where  $(\partial G_{\rm E}/\partial q^2)_q 2_{=0}$  is the slope of the neutron form factor at zero square of the transferred momentum, and  $a_{\rm ne}$  is the amplitude of neutron-electron interaction. In ref. /1/ the value of  $a_{\rm ne} = (-1.378\pm0.018) \, 10^{-3}$  fm corresponding to  $(\frac{\partial G_{\rm E}}{\partial q^2})_q 2_{=0} =$  $= -0.0199\pm0.0003$  fm<sup>2</sup> was obtained from the comparison of experimental data on coherent scattering amplitudes  $b_{\rm coh}$  of neutrons on Bismuth and Lead measured using the neutron gravity reffractometer with those on scattering cross sections of neutrons on the nuclei of the same elements at an energy of 1.26 and 5.19 eV. This value of  $(\frac{\partial G_{\rm E}}{\partial q^2})_q 2_{=0}$  is usually used for the analysis of experimental data on scattering of high energy electrons by deuterons (see, for example, ref. '2').

The purpose of the present paper is to draw attention to the necessity of taking into account the effect of the electrical polarizability of the neutron on the value of  $a_{ne}$  during the analysis of experimental data obtained in ref.<sup>/1/</sup>. As it will be shown below this effect causes a considerable increase of the error in determining the  $a_{ne}$  value.

Let us use the formalism conventionally used for the problems of scattering of a particle on a sum of short-range and longrange potentials (see, for example, ref. '3'). The neutron scattering amplitude  $f_t$  may be expressed through the phases of nuclear  $\delta_\ell$ , neutron-electron  $\eta_\ell$  and polarization  $\xi_\ell$  scattering. At low energies, when only the s-scattering is important, one has

$$f_{t} = \frac{1}{2ik} (S_{0} - 1) \exp[2i(\eta_{0} + \xi_{0})] + \frac{1}{k} \sum_{\ell} (2\ell + 1) \sin(\eta_{\ell} + \xi_{\ell}) \exp[i(\eta_{\ell} + \xi_{\ell}) P_{\ell}(\cos\theta)],$$
(2)

where

$$S_0 = (1 - i\Sigma \frac{\Gamma_n}{\Delta E + i\Gamma/2}) \exp(2i\delta_0).$$

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Within the Born approximation and in the absence of resonances (their effect on the cross sections of Bismuth and Lead at energies 1.26 and 5.19 eV is weak and accounted for in ref.<sup>(1)</sup>) we have:

$$\operatorname{Ref}_{t}(0) = \frac{1}{k} \sin \delta_{0} \cos \left[ \delta_{0} + 2(\eta_{0} + \xi_{0}) \right] + z a_{ne} + a = b_{coh}$$
(3)

 $\operatorname{Im} f_{t}(0) = \frac{1}{k} \sin \delta_{0} \sin \left[ \delta_{0} + 2(\eta_{0} + \xi_{0}) \right] = \frac{k \sigma_{s}}{4\pi}, \qquad (4)$ 

where

$$a = \frac{M_{\alpha}}{2R} \left(\frac{ze}{h}\right)^2 qR \left(\frac{\sin qR}{(qR)^2} + \frac{\cos qR}{qR} + \sin qR\right)$$

is the polarization scattering amplitude<sup>14</sup>,  $\alpha$  is the coefficient of electrical polarizability of the neutron,  $q = \frac{2}{5} \sin \theta / 2$ 

$$\eta_0 = \frac{a_{ne}k}{2} \int_0^{\pi} f(\frac{\sin\theta}{\lambda}) \sin\theta d\theta \quad \xi_0 = \frac{Ma}{R} \left(\frac{ze}{n}\right)^2 k(1 - \frac{\pi}{3}kR)$$

 $f(\frac{\sin \theta}{2})$  is the atomic form-factor.

The relationships (3) and (4) contain the unknown values of  $\delta_0$ ,  $a_{ne}$ , and a. They can be found by substituting into (3) and (4) of experimental values from ref.<sup>/1/</sup> (at two values of  $\sigma_s$ ). The results are listed in table 1. It shows that the data from ref.<sup>/1/</sup> allows one to find  $a_{ne}$ , but with a large error.

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ane	(-1.5+0.7)×10 <sup>-3</sup> fm	$(-1.4+0.8) \times 10^{-3} \text{ fm}$
a	$(12+12) \times 10^{-3}$ fm <sup>3</sup>	$(6+6) \times 10^{-3}$ fm <sup>3</sup>

The authors of ref.<sup>/5/</sup> obtained a value of  $a_{ne} = (-1.33 + .0.03) \times 10^{-3}$  fm (with the account for the Schwinger scattering)\* One may try to find *a* using this value of  $a_{ne}$  in the analysis of data from ref.<sup>/1/</sup>. It appears to be the complex value, which might indicate to some contradiction between the data on  $b_{coh}$  and  $\sigma_s$  (ref.<sup>/1/</sup>) and those on  $a_{ne}$  (ref.<sup>/5/</sup>). The minimum real value of  $a \simeq 9 \cdot 10^{-4}$  fm<sup>3</sup> corresponds to the value of  $a_{ne} \simeq -1.37 \cdot 10^{-3}$  fm. The value of a increases with increasing  $|a_{ne}|$ . Table 2 gives the results of data procession (ref.<sup>/1/</sup>) under assumption that  $a_{ne}$  is determined by the Foldy interaction (ref.<sup>/8/</sup>) giving

$$a_{ne} = \mu_n \frac{e^2}{2Mc^2} = -1.468 \times 10^{-3} \text{ fm}.$$

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$$a_{ne} = -1.468 \times 10^{\circ} \text{ fm}$$
  
Bismuth Lead  
eV 5.19 eV 1.26 eV 5.19

-3

 $\frac{a (12+2) \times 10^{-3} \text{ fm}^{8} (9+2) \times 10^{-3} \text{ fm}^{3} (12+2) \times 10^{-8} \text{ fm}^{3} (13+2) \times 1+ \text{ fm}^{3}}{a} = (11+1) \times 10^{-8} \text{ fm}^{8}$ 

Let us note in conclusion that the above information about  $\alpha$  is not in contradiction with the value of  $\alpha = (9 \pm 5) \times 10^{-3} \text{ fm}^3$ obtained in ref.<sup>77</sup> in the analysis of slow neutron scattering on <sup>186</sup>W.

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Table 1

1.26

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<sup>\*</sup> The results of other works on  $a_{ne}$  contain even larger errors, and some of them should be corrected for the effect of the polarizability of the neutron.

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Обращается внимание на ческого формфактора нейтро на атомах свинца и висмута нейтрона.	а то, что при полу на из данных по ра необходимо учитые	чении величины Ссеянию медленн ать влияние пол	наклона электри ых нейтронов яризуемости
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The investigation has Physics, JINR.	been performed a	t the Laboratory	y of Neutron
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