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ON FRACTAL STRUCTURE OF QUANTUM GRAVITY AND RELIC RADIATION ANISOTROPY



#### 1 Introduction

It seems that a fractal structure is an intimate property of the Universe. The superclusters (large clusters of galaxies containing up to hundreds of thousands of galaxies) of size about 50 Mpc are separated by almost void space: the mean distance between two superclusters is about 100 Mpc. Clusters of galaxies (the typical cluster size is about 5 Mpc) containing hundreds of galaxies are, in their turn, separated by voids about few Mpc. This fractal hierarchy can be easily traced up to subnuclear scales ( $10^{-13}$  cm.). Quantitatively, the large-scale fractal structure of the Universe can be described in terms of the mass interior to the spherical volume of certain radius r. The typical dependence, measured by observing 21 cm hydrogen emmission of gas clouds moving around the galaxy is

$$\mathcal{M}(r) \propto r^{\alpha}, \qquad \alpha \approx 1.$$
 (1)

whereas a luminous mass associated with the light would supply only  $r^{-1/2}$ . It is commonly accepted that an additional mass in the form of non-luminous dark matter [1]. Since  $\alpha < 3$  in the power law (1), we have a typical mass distribution on a fractal set embedded in D=3 space.

On the other hand, one of the most important recent developments in gravity theory was related to the fractal based regularization of quantum gravity [3](See also review [8]). In view of this one may believe that a fractal structure is a fundumental property of physical space-time itself.

In this note we interpret the COBE satellite data on the anisotropy of cosmic microwave background radiation (CMBR) as a possible manifestation of fractal structure of the Universe.

## 2 On the discrete symmetry in quantum gravity

A regularization of two-dimensional quantum gravity, made by V. G. Knizhnik, A. M. Polyakov and A. B. Zamolodchikov (KPZ), comes from the fact that continuum formulation [2] and the dynamical triangulation [4] are equivalent. On the basis of the Polyakov regularization procedure [5], where the position of the surface in the embedding space  $X_{\mu}$  and the internal surface geometry  $(g_{ab})$  are treated as independent fields, one can construct a Nambu-like action

$$S[X_{\mu}, g_{ab}] =$$

$$= \frac{1}{2} \int_{M} g_{ab} \frac{\partial X_{\mu}}{\partial \xi_{a}} \frac{\partial X_{\mu}}{\partial \xi_{b}} \sqrt{\det g} \ d^{2}\xi + \beta \int_{M} \sqrt{\det g} \ d^{2}\xi + \qquad (2)$$

+fermion terms

where  $\xi = (\xi_1, \xi_2)$  is the parametrization of the manifold M defined by function  $X_{\mu} = X_{\mu}(\xi)$ . This or a similar Nambu-Goto action usually stands in the string functional integral taken with respect to both independent fields X and g.

In the dynamical triangulation [4] of 2D quantum gravity, the path integral over internal metric  $g_{ab}$  is replaced by summation of all the different types of surface configurations with given number of triangles. For the sake of preserving reparametrization invariance after discretization [4], the topology of the manifold

M is usually specified as the sphere  $S^2$  [4, 8]. The partition function takes the form

$$Z(A) = \int_{M} \mathcal{D}X \mathcal{D}g \exp(-S), \qquad (3)$$

or its discrete counterpart [8]

$$Z_{reg}(A) = \sum_{G} Z_m(G) \delta_{Na^2, A}, \qquad (4)$$

where A is the total area, N is the number of equilateral triangles and  $a^2$  is the area of an triangle. The matter part of the partition function  $Z_m(G)$  comes from the fermion term of the KPZ lagrangian

$$\mathcal{L} = \bar{\phi} v^{\alpha a} \gamma^a \partial_\alpha \phi, \tag{5}$$

where  $v_{\alpha a}$  are ordinary "zweibeins". (See [8] and references therein for the details concerning its discrete counterpart.)

Formally substituting functional integral (3) by its discrete counterpart (4) we need to sum over all possible triangulations of  $S^2$ . Practically, we are to impose some additional conditions to avoid summation over singular triangulations, i. e. triangulations which include links with coinciding ends. Referring the reader to [4, 9] for detailed study of triangulations and fractal properties of related partition functions, we shall concentrate on some of their properties significant for the phenomenological applications.

First. The triangulation procedure can be extended to an  $S^n$  sphere embedded in Euclidean space  $\mathbb{R}^q, q > n$ : sphere  $S^n$ , as a boundary of (n + 1)-dimensional simplex, can be divided into *n*-dimensional simplices [10].

Second. From the conformal invariance standpoint, of all the subdivisions of  $S^n$  the subdivision into equilateral simplices is preferable.

Third. The whole partition function (3) is related to the physical object, which is isotropic (in the sense of having no preferable direction on  $S^n$ ), but may have a discrete symmetry

group, and hence, have certain distinguished correlation angles. For example, if we sum over all possible triangulations of  $S^2$ using equilateral triangles, the correlations of any observables depending on matter fields increase at angles  $0, \frac{2\pi}{3}, \frac{4\pi}{3}$  because of Z<sub>3</sub> symmetry group. Similarly, the correlations should increase at tetrahedron group angles when  $S^3$  is considered.

Fourth. The two-dimensional quantum gravity can be regarded as only the simplest case of extended object physics. However, when reducing the physics from arbitrary *n*-dimensional space to (n-1) dimensions we restrict  $S^n$  triangulation with *n*dimensional simplices to  $S^{n-1}$  triangulation with (n-1) - dimensional ones, because an (n-1)-dimensional simplex is a boundary of an *n*-dimensional one. Thus, for the case of equilateral symplexes we should always have Z<sub>3</sub>-symmetry in D = 2, or tetrahedron symmetry in D = 3.

### 3 Discrete symmetry as a possible source of relic radiation anisotropy

Let us consider the data [11] on relic radiation anisotropy. The relic microwave radiation  $(T = 2.73^{\circ}K)$  was not significantly affected by the late-stage processes in the Universe, that is why its amplitudes depend mostly on the early Universe parameters. It is worth to note, that the large scale anisotropy of relic radiation, found in COBE and RELICT-1 experiments, has a rather small value  $\frac{\Delta T}{T} \sim 10^{-5}$ , but a high confidence level — up to 90%, including systematic errors [11, 14].

Of course, the first aim of the observers in both COBE and RELICT experiments was to measure the dipole and quadrupole components of microwave background [11] and to test the existence of anomalous signal distinguishing from mean background [14]. Besides that, basing on the COBE experiment data, the autocorrelation function

$$C(\alpha) = \langle \Delta T(\theta) \Delta T(\theta + \alpha) \rangle \tag{6}$$

has been obtained. Here  $\alpha$  is the angle separation and  $\theta$  is an angular coordinate on certain two-dimensional plane.

Qualitatively, the behavior of the relic signal autocorrelation function (See fig.2, reprinted at the end of our paper, and fig.3 from [11]) is the following: it has a sharp maximum up to 1000  $(\mu K)^2$  at  $\alpha = 0$ , it has another maximum (up to 200  $(\mu K)^2$ ) localized at  $\alpha$  close to 120 degrees, and it has two minimums at 60° and 180°. (The data cover the interval  $0 < \alpha < 180^{\circ}$ .). The behavior of the autocorrelation functions is just the same for the data obtained at frequencies 53 GHz and 90 GHz [11].

Correlation function (6) has been studied in [13] in connection with present cosmological models. In particular, an attempt has been made to compare the COBE data with certain Dark Matter (DM) models. This comparison does not suit well. For instance, the relic density anisotropy given by Holtzman model [15] increases monotonously with  $\alpha$  increasing from 60 to 180 degrees [13].

Taking into account all the arguments mentioned above, we interpret the regularities of autocorrelation function (6) behavior, as a manifestation of Z<sub>3</sub>-symmetry. The presense of Z<sub>3</sub>-symmetry does not imply n preferable directions in space here, instead we have a preferable separation angle. It should be mentioned that in COBE theoretical study [13] the best line fit for autocorrelation function (6) was taken in the form

$$C(\alpha) = A + B\cos\alpha + C_M^0 \exp\left[-\frac{\alpha^2}{2\sigma^2}\right],$$
 (7)

though the locations of autocorrelation function maximums at  $0^{\circ}$  and  $120^{\circ}$  and minimums at  $60^{\circ}$  and  $180^{\circ}$  suggest more direct parametrization

$$C(\alpha) = A + B\cos 3\alpha + C_M^0 \exp\left[-\frac{\alpha^2}{2\sigma^2}\right]$$
(8)

## 4 Conclusion

The data on relic radiation anisotropy obtained by both RELICT and COBE groups are worth further deep investigation. Nonetheless, even the results already obtained from data processing seem to be in good agreement with the hypothesis of discrete symmetry of space-time arising in fractal quantum gravity. Other cosmological data, e.g. mass distribution, also do not contradict either possible fractal structure of the Universe. It might be argued, that both the tetrahedron symmetry, if found, and the fractal structure of the visible Universe, can be regarded as an argument for the existence of cosmic strings [16]. Indeed, cosmic strings, as topological defects which could be formed at a phase transition in the early Universe, can have a number of cosmological applications. In particular, they can form a network with a fractal structure having tetrahedron symmetry [17]. Though the question, why cosmic strings must form the tetrahedron structure still remains.

Therefore, some new tests for possible discrete symmetry can be proposed. The simplest among them are: (i) to test  $\cos n\alpha$ , n > 1 in (7) for other  $Z_n$  groups, (ii) to use COBE and RELICT data to search for the tetrahedron or other essentially three-dimensional space symmetry groups.

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Correlation functions  $C(\alpha)$ , at various Gallactic lattitude cuts for the 53MHz map. (Reprinted from [11], fig.2).

#### References

- R.Sancisi and T.S. van Albada in *Dark Matter in the Universe*, eds J.Kormendy and G.Knapp (Reidel, Dordrecht, 1987) p.67
- [2] A.M.Polyakov, Mod. Phys. Lett A2(1987)893
- [3] V.G.Knizhnik , A.M.Polyakov and A.B.Zamolodchikov, Modern. Phys. Lett. A3(1988)819
- [4] D.V.Boulatov, V.A.Kazakov, I.K.Kostov and A.A.Migdal, Nucl. Phys. B275(1986)641
- [5] A.M.Polyakov, Phys. Lett. B103(1981)207
- [6] E.Brezin and V.A.Kazakov, Phys. Lett. B236(1990)114
   M.Douglas and S.Shenker, Nucl. Phys. B335(1990)635
   D.J.Gross and A.A.Migdal, Phys. Rev. Lett. 64(1990)127
- [7] M.E.Agishtein and A.A.Migdal. Int. J. Mod. Phys. C1(1990)167; Nucl. Phys. B350(1991)690
- [8] N.Kawamoto, Institute for Nuclear Study of Tokyo, Preprint INS-Rep. 972, Apr. 1993
- [9] N.Kawamoto, V.A.Kazakov, Y. Saeki and Y.Watabiki, Phys. Rev. Lett. 68(1992)2113
- [10] A.Fomenko, D.Fuchs, A Course in Homotopy Topology. (Chapter 1), Moscow, Nauka, 1989
- [11] G.Smoot et al., Astrophys. J. 396(1992)L1
- [12] C.L.Bennet et al., Astrophys. J. 396(1992)L7
- [13] E.L. Wright *et al.*, Astrophys. J. 396(1992)L13

- [14] I.A.Strukov, A.A. Brukhanov, D.P. Skulachev and M.V. Sazhin, Phys.Lett. B315(1993)198
- [15] J.A.Holtzman, Astrophys. J. Suppl. 71(1989)1
- [16] T.W.B. Kibble, J.Phys. A9(1976)1387
- [17] M.Aryal, A.E. Everett and A.Vilenkin, Phys. Rev. D34(1986)434

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