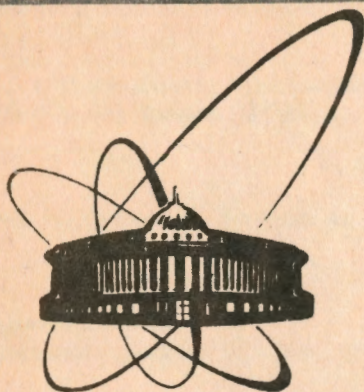


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ОБЪЕДИНЕННЫЙ
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A METHOD FOR OBTAINING
QUANTUM DOUBLES
FROM THE YANG-BAXTER R -MATRICES

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Метод получения квантовых дублей
из R -матриц Янга — Бакстера

В работах [1] и [2] был предложен метод, позволивший сопоставить каждому регулярному обратимому решению уравнений Янга — Бакстера без спектрального параметра некоторую квазитреугольную алгебру Хопфа. Мы развиваем этот подход и показываем, что эта алгебра Хопфа есть на самом деле квантовый дубль.

Работа выполнена в Лаборатории теоретической физики ОИЯИ.

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A Method for Obtaining Quantum Doubles
from the Yang-Baxter R -Matrices

We develop the approach of refs. [1] and [2] that enables one to associate a quasitriangular Hopf algebra to every regular invertible constant solution of the quantum Yang-Baxter equations. We show that such a Hopf algebra is actually a quantum double.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

1. It is well known [1] that any invertible constant matrix solution R of the quantum Yang-Baxter equation (QYBE)

$$R_{12}R_{13}R_{23} = R_{23}R_{13}R_{12} \quad (1)$$

naturally generates a bialgebra $A_R = \{1, t_{ij}\}$ defined by

$$R_{12}T_1T_2 = T_2T_1R_{12}, \quad \Delta(T_1) = T_1 \otimes T_1, \quad \varepsilon(T) = 1,$$

(generators t_{ij} form a matrix T , Δ is a coproduct and ε a counit) and also another bialgebra $U_R = \{1, l_{ij}^{\pm}, \bar{l}_{ij}^{\pm}\}$ with

$$R_{12}L_2^{\pm}L_1^{\pm} = L_1^{\pm}L_2^{\pm}R_{12}, \quad (2)$$

$$R_{12}L_2^+L_1^- = L_1^-L_2^+R_{12}, \quad (3)$$

$$\Delta(L_1^{\pm}) = L_1^{\pm} \otimes L_1^{\pm}, \quad \varepsilon(L^{\pm}) = 1, \quad (4)$$

which is paired to A_R . This pairing [1, 2] is established by the relations

$$\langle T_1, L_2^+ \rangle = R_{12}, \quad \langle T_1, L_2^- \rangle = R_{21}^{-1},$$

obeys the duality conditions

$$\langle \alpha\beta, a \rangle = \langle \alpha \otimes \beta, \Delta(a) \rangle, \quad \langle \Delta(\alpha), a \otimes b \rangle = \langle \alpha, ab \rangle,$$

and appears to be degenerate. With some additional effort (quotienting by appropriate null bi-ideals) these bialgebras can be made Hopf algebras \check{A}_R and \check{U}_R , dual to each other. Their antipodes are defined by

$$\langle T_1, S(L_2^+) \rangle = \langle S(T_1), L_2^+ \rangle = R_{12}^{-1},$$

$$\langle T_1, S(L_2^-) \rangle = \langle S(T_1), L_2^- \rangle = R_{21}.$$

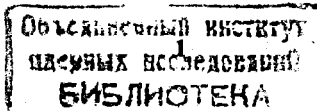
With essential use of this duality Majid [2] showed that in fact, with a certain reservation, \check{U}_R proves to be a quasitriangular Hopf algebra with the universal \mathcal{R} -matrix given by implicit formulas originated from $\langle T_1 \otimes T_2, \mathcal{R} \rangle = R_{12}$. By the way, Majid claims [2] that \check{U}_R is 'more or less' of the form of a quantum double. In the present note we argue that, modulo the same reservation, \check{U}_R is actually a quantum double.

2. Recall t(i.e. the duality with opposite coproduct and inverse antipode) means $\langle e$

$$^i, e_j \rangle = \delta_j^i \text{ and}$$

$$\langle \alpha\beta, a \rangle = \langle \alpha \otimes \beta, \Delta(a) \rangle, \quad \langle \Delta(\alpha), a \otimes b \rangle = \langle \alpha, ba \rangle, \quad (5)$$

$$\varepsilon(a) = \langle 1, a \rangle, \quad \varepsilon(\alpha) = \langle \alpha, 1 \rangle, \quad \langle S(\alpha), a \rangle = \langle \alpha, S^{-1}(a) \rangle,$$



where $a, b \in A$, $\alpha, \beta \in A^\circ$, and $\{e_j\}, \{e^i\}$ are the corresponding bases. To equip $A \otimes A^\circ$ with the Hopf algebra structure of the quantum double, one must define a very specific cross-multiplication recipe. If

$$e_i e_j = c_{ij}^k e_k, \quad \Delta(e_i) = f_i^{jk} (e_j \otimes e_k), \quad S(e^i) = \sigma_j^i e^j,$$

it reads

$$e^i e_j = \mathcal{O}_{jq}^{ip} e_p e^q, \quad \text{where } \mathcal{O}_{jq}^{ip} = c_{nq}^i c_{ts}^i \sigma_r^s f_j^{rl} f_l^{pn}.$$

In invariant form this looks like

$$\alpha a = \sum \sum \langle S(\alpha_{(1)}), a_{(1)} \rangle \langle \alpha_{(3)}, a_{(3)} \rangle a_{(2)} \alpha_{(2)}, \quad (6)$$

where

$$\Delta^2(\alpha) = \sum \alpha_{(1)} \otimes \alpha_{(2)} \otimes \alpha_{(3)}, \quad \Delta^2(a) = \sum a_{(1)} \otimes a_{(2)} \otimes a_{(3)}.$$

Here the usual notation for coproducts (cf. $\Delta(a) = \sum a_{(1)} \otimes a_{(2)}$) is used. The resulting Hopf algebra proves to be quasitriangular with the universal \mathcal{R} -matrix

$$\mathcal{R} = \sum_i (e_i \otimes 1) \otimes (1 \otimes e^i).$$

One easily finds that $gl_q(2)$ and other simple examples of quantum universal enveloping algebras are both the \hat{U}_R -type algebras and quantum doubles. Can it happen that \hat{U}_R would be a quantum double for any R ? Majid's approach based on the $\hat{U}_R \leftrightarrow \hat{A}_R$ duality does not readily answer this question. That is why we choose another way: not to use A_R at all. The key observation is that there exists an inherent antiduality between U_R^+ and U_R^- which is precisely of the form required for the quantum-double construction.

3. Let us define bialgebras $U_R^+ = \{1, l_{ij}^+\}$ and $U_R^- = \{1, l_{ij}^-\}$ by eqs. (2),(4). Note that the cross-multiplication relation (3) is not yet imposed, so U_R^+ and U_R^- are considered to be independent so far. However, the very natural pairing between them can be introduced. It is generated by

$$\langle L_1^-, L_2^+ \rangle = R_{12}^{-1}, \quad \langle L^-, 1 \rangle = \langle 1, L^+ \rangle = \langle 1, 1 \rangle = 1 \quad (7)$$

and in the general case looks like

$$\langle L_{i_1}^- \dots L_{i_m}^-, L_{j_1}^+ \dots L_{j_n}^+ \rangle = R_{i_1 j_n}^{-1} \dots R_{i_m j_1}^{-1}, \quad (8)$$

where the r.h.s. is a product of mn R^{-1} -matrices corresponding to all pairs of indices $i_q j_p$ with j -indices ordered from right to left. The consistency of (8) and (5) with (4) is evident, while the proof of the consistency with (2) reduces to manipulations like

$$\begin{aligned} \langle L_0^-, R_{12} L_2^+ L_1^+ - L_1^+ L_2^+ R_{12} \rangle &= \langle L_0^- \otimes L_0^-, R_{12} (L_1^+ \otimes L_2^+) - (L_2^+ \otimes L_1^+) R_{12} \rangle \\ &= R_{12} R_{01}^{-1} R_{02}^{-1} - R_{02}^{-1} R_{01}^{-1} R_{12} = 0 \end{aligned}$$

and repeated use of QYBE (1).

For general R , this pairing is degenerate. To remove the degeneracy, i.e. to transform pairing into antiduality, one should factor out appropriate bi-ideals [2]. In simple cases this procedure is explicitly carried out and works well. For general R it is of course not under our control. The situation is quite similar to [2]: we are to rely on that the factorization procedure is "soft" in a sense that it does not destroy the whole construction.

4. Keeping this in mind, we observe that, being antidual, U_R^\pm admit the Hopf algebra structure. Let us introduce an antipode S in U_R^- and an inverse antipode S^{-1} in U_R^+ by the relations

$$\langle S(L^-), 1 \rangle = \langle 1, S^{-1}(L^+) \rangle = 1, \quad \langle S(L_1^-), L_2^+ \rangle = \langle L_1^-, S^{-1}(L_2^+) \rangle = R_{12}, \quad (9)$$

extending them on the whole of U_R^+ (or U_R^-) as antihomomorphisms of algebras and coalgebras. The definition is correct due to

$$\begin{aligned} \langle m \circ (S \otimes id) \circ \Delta(L_1^-), L_2^+ \rangle &= \langle m \circ (S(L_1^-) \otimes L_1^-), L_2^+ \rangle \\ &= \langle S(L_1^-) \otimes L_1^-, L_2^+ \otimes L_2^+ \rangle = R_{12} R_{12}^{-1} = 1 = \langle \varepsilon(L_1^-), L_2^+ \rangle, \end{aligned}$$

$\langle S(L_0^-), R_{12} L_2^+ L_1^+ - L_1^+ L_2^+ R_{12} \rangle = \langle \Delta \circ S(L_0^-), R_{12} (L_1^+ \otimes L_2^+) - (L_2^+ \otimes L_1^+) R_{12} \rangle = \langle S(L_0^-) \otimes S(L_0^-), R_{12} (L_2^+ \otimes L_1^+) - (L_1^+ \otimes L_2^+) R_{12} \rangle = R_{12} R_{02} R_{01} - R_{01} R_{02} R_{12} = 0,$
 $\langle S(R_{12} L_2^- L_1^- - L_1^- L_2^- R_{12}), L_0^+ \rangle = \langle R_{12} S(L_1^-) S(L_2^-) - S(L_2^-) S(L_1^-) R_{12}, L_0^+ \rangle = \langle R_{12} (S(L_1^-) \otimes S(L_2^-)) - (S(L_2^-) \otimes S(L_1^-)) R_{12}, L_0^+ \otimes L_0^+ \rangle = R_{12} R_{10} R_{20} - R_{20} R_{10} R_{12} = 0,$
but there is no such a formula for $S(L^+)$ or $S^{-1}(L^-)$. Once again the factorization is hoped to be soft enough to allow the antipodes to be invertible.

If so, our bialgebras U_R^\pm become the mutually antidual Hopf algebras \hat{U}_R^\pm , and it is possible to define the multiplicative structure of the quantum double upon $\hat{U}_R^+ \otimes \hat{U}_R^-$. The cross-multiplication rule is deduced from (6):

$$\begin{aligned} L_1^- L_2^+ R_{12} &= \langle S(L_1^-), L_2^+ \rangle L_2^+ L_1^- \langle L_1^-, L_2^+ \rangle R_{12} \\ &= R_{12} L_2^+ L_1^- R_{12}^{-1} R_{12} = R_{12} L_2^+ L_1^-. \end{aligned} \quad (10)$$

Thus we regain eq.(3) as the quantum-double cross-multiplication condition!

Our conclusion is that R -matrices obeying QYBE generate the algebraic structures of quantum double in quite a natural way.

5. To illustrate the proposed scheme, consider the $sl_q(2)$ R -matrix

$$R_q = q^{-1/2} \begin{pmatrix} q & 0 & 0 & 0 \\ 0 & 1 & q - q^{-1} & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & q \end{pmatrix}, \quad R_q^{-1} = R_{q^{-1}}.$$

Here the bialgebras $U_{R_q}^\pm$ have 8 generators l_{ij}^\pm . The bi-ideals to be factored out are generated by the relations

$$l_{21}^- = 0, \quad l_{12}^+ = 0, \quad l_{11}^- l_{22}^- = l_{22}^- l_{11}^- = 1, \quad l_{11}^+ l_{22}^+ = l_{22}^+ l_{11}^+ = 1.$$

After factorization the number of independent generators is reduced to 4. We denote them X^\pm, H, H' (note that $H' \neq H$ so far):

$$L^+ = \begin{pmatrix} q^{H/2} & 0 \\ (q^{1/2} - q^{-3/2})X^+ & q^{-H/2} \end{pmatrix}, \quad L^- = \begin{pmatrix} q^{-H'/2} & (q^{-1/2} - q^{3/2})X^- \\ 0 & q^{H'/2} \end{pmatrix}.$$

The multiplication rules (inside each algebra), coproducts and antipodes are:

$$\begin{aligned} [H, X^+] &= 2X^+, & [H', X^-] &= -2X^-, \\ \Delta(H) &= H \otimes 1 + 1 \otimes H, & \Delta(H') &= H' \otimes 1 + 1 \otimes H', \\ \Delta(X^+) &= X^+ \otimes q^{H/2} + q^{-H/2} \otimes X^+, & \Delta(X^-) &= X^- \otimes q^{H'/2} + q^{-H'/2} \otimes X^-, \\ S(X^\pm) &= -q^{\pm 1} X^\pm, & S(H) &= -H, \quad S(H') = -H'. \end{aligned}$$

The quantum-double cross-multiplication rules (6) take the form

$$\begin{aligned} [H', X^+] &= 2X^+, & [H, X^-] &= -2X^-, & [H, H'] &= 0, \\ [X^+, X^-] &= (q^{(H+H')/2} - q^{-(H+H')/2}) / (q - q^{-1}). \end{aligned}$$

The identification $H' \equiv H$ leads to the ordinary $sl_q(2)$.

6. To give one more illustration, let us consider a bialgebra introduced in [5]. In a slightly simplified form it has generators $\{1, t_j^i, u_j^i, E_j, F^i\}$ which obey the following relations (here we prefer to display all the indices):

$$R_{mn}^{ij} t_p^m t_q^n = R_{pq}^{mn} t_n^i t_m^j, \quad E_p t_q^j = R_{pq}^{mn} t_n^j E_m, \quad (11)$$

$$\Delta(t_j^i) = t_k^i \otimes t_j^k, \quad \varepsilon(t_j^i) = \delta_j^i, \quad \Delta(E_j) = E_i \otimes t_j^i + 1 \otimes E_j, \quad \varepsilon(E_j) = 0, \quad (12)$$

$$R_{mn}^{ij} u_p^m u_q^n = R_{pq}^{mn} u_n^j u_m^i, \quad F^i u_p^j = R_{mn}^{ji} u_p^m F^n, \quad (13)$$

$$\Delta(u_j^i) = u_k^i \otimes u_j^k, \quad \varepsilon(u_j^i) = \delta_j^i, \quad \Delta(F^i) = F^i \otimes 1 + u_j^i \otimes F^j, \quad \varepsilon(F^i) = 0, \quad (14)$$

$$R_{mn}^{ij} u_p^m t_q^n = R_{pq}^{mn} t_n^j u_m^i, \quad E_j F^i - F^i E_j = t_j^i - u_j^i, \quad (15)$$

$$u_p^i E_q = R_{pq}^{mn} E_n u_m^i, \quad t_p^i F^j = R_{mn}^{ji} F^m t_p^n, \quad (16)$$

with R obeying QYBE (1). This is not a bialgebra of the form (2)-(4). Rather it is of the 'inhomogeneous quantum group' type [6]. Let us make sure that it is a quantum double as well.

Consider T, E -bialgebra (11),(12) and U, F -bialgebra (13),(14) firstly as being independent and fix nonzero pairings on the generators by

$$\langle u_j^i, t_q^p \rangle = R_{jq}^{ip}, \quad \langle u_j^i, 1 \rangle = \langle 1, t_j^i \rangle = \langle F^i, E_j \rangle = \delta_j^i, \quad (17)$$

extending them to the whole bialgebras with the help of (5). The definition is correct due to

$$\begin{aligned} \langle F^i, E_p t_q^j - R_{pq}^{mn} t_n^j E_m \rangle &= \langle F^i \otimes 1 + u_k^i \otimes F^k, t_q^j \otimes E_p - R_{pq}^{mn} (E_m \otimes t_n^j) \rangle \\ &= R_{kq}^{ij} \delta_p^k - \delta_m^i R_{pq}^{mn} \delta_n^j = 0, \\ \langle F^i u_p^j - R_{mn}^{ji} u_p^m F^n, E_q \rangle &= \langle F^i \otimes u_p^j - R_{mn}^{ji} (u_p^m \otimes F^n), E_k \otimes t_q^k + 1 \otimes E_q \rangle \\ &= \delta_k^i R_{pq}^{jk} - R_{mn}^{ji} \delta_p^m \delta_q^n = 0. \end{aligned}$$

After factoring out the corresponding null bi-ideals, we may define antipodes on the generators as follows:

$$\langle S(u_j^i), t_q^p \rangle = \langle u_j^i, S^{-1}(t_q^p) \rangle = (R^{-1})_{jq}^{ip},$$

$$\langle S(u_j^i), 1 \rangle = \langle 1, S^{-1}(t_j^i) \rangle = \delta_j^i, \quad S(F^i) = -S(u_j^i) F^j, \quad S(E_j) = -E_i S(t_j^i).$$

The proof of correctness is in complete analogy with the \check{U}_R -case.

Now a direct application of the recipe (6) exactly reproduces the cross-multiplication relations (15),(16). For example,

$$\begin{aligned} F^i E_j &= \langle S(F^i), E_m \rangle \langle 1, t_j^n \rangle t_m^n + \langle S(u_n^i), 1 \rangle \langle 1, t_j^m \rangle E_m F^n \\ &+ \langle S(u_n^i), 1 \rangle \langle F^m, E_j \rangle u_n^m = -t_j^i + E_j F^i + u_j^i, \end{aligned}$$

because of

$$\langle S(F^i), E_m \rangle = -\langle S(u_k^i) \otimes F^k, E_n \otimes t_m^n + 1 \otimes E_m \rangle = -\delta_m^i.$$

Therefore, bialgebras of the type (11)-(14) are also transformed into the quantum double using our method.

7. Consider at last a bialgebra [7] that is known to be related [8] to bicovariant differential calculus on quantum groups. Its coalgebra structure is given by (12), whereas the multiplication relations (11) are to be supplemented by

$$t_p^i E_q + f_{nm}^i t_p^m t_q^n = R_{pq}^{nm} E_m t_n^i + f_{pq}^n t_n^i, \quad (18)$$

$$E_i E_j - R_{ij}^{mn} E_n E_m = f_{ij}^m E_m, \quad (19)$$

f_{jk}^i being new structure constants. This bialgebra, unlike its ancestor (11),(12), exhibits the R -matrix-type representation

$$\mathbf{R}_{12} \mathbf{T}_1 \mathbf{T}_2 = \mathbf{T}_2 \mathbf{T}_1 \mathbf{R}_{12}, \quad \Delta(\mathbf{T}) = \mathbf{T} \otimes \mathbf{T}, \quad (20)$$

where, in terms of multi-indices like $I = \{0, i\}$,

$$\mathbf{T}_J^I = \begin{pmatrix} 1 & E_j \\ 0 & t_j^i \end{pmatrix}, \quad \mathbf{R}_{MN}^{IJ} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \delta_n^j & 0 & 0 \\ 0 & 0 & \delta_m^i & f_{mn}^i \\ 0 & 0 & 0 & R_{mn}^{ij} \end{pmatrix}.$$

Of course, \mathbf{R} must satisfy the QYBE (1) which now involves the structure constants f_{mn}^i as well as R_{mn}^{ij} . Note that, due to (18),(19), the bialgebra (11),(12) is not restored from (20) by mere setting $f_{mn}^i \equiv 0$.

Now let us try to develop a quantum double from the bialgebra (20). However, it seems to be quite uneasy task. A natural Ansatz for the candidate antidual bialgebra is

$$\mathbf{U}_J^I = \begin{pmatrix} 1 & 0 \\ F^i & u_j^i \end{pmatrix},$$

which causes the corresponding R -matrix to be

$$\bar{\mathbf{R}}_{MN}^{IJ} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \delta_n^j & 0 & 0 \\ 0 & 0 & \delta_m^i & 0 \\ 0 & \bar{f}_n^{ij} & 0 & R_{mn}^{ij} \end{pmatrix}$$

with different structure constants \bar{f} and another QYBE system involving R and \bar{f} . Now, attempting to fix a pairing in the form

$$\langle U_1, T_2 \rangle = Q_{12} \quad (21)$$

with a certain numerical matrix Q , we immediately arrive at the following general statement:

Let \mathbf{R} and $\bar{\mathbf{R}}$ be invertible solutions of QYBE. If there exists an invertible solution Q of the equations

$$\begin{aligned} Q_{12} Q_{13} R_{23} &= R_{23} Q_{13} Q_{12}, \\ \bar{R}_{12} Q_{13} Q_{23} &= Q_{23} Q_{13} \bar{R}_{12}, \end{aligned}$$

then (21) is a correct pairing between the T - and U -bialgebras generated by \mathbf{R} and $\bar{\mathbf{R}}$, respectively, and, assuming a proper quotienting procedure to be performed, the antipodes can be defined by the relations

$$\langle S(U_1), T_2 \rangle = \langle U_1, S^{-1}(T_2) \rangle = Q_{12}^{-1}$$

and the quantum-double structure can be established on the tensor product of these bialgebras by the cross-multiplication formula

$$Q_{12} U_1 T_2 = T_2 U_1 Q_{12}.$$

Whether such a program can really be carried through in interesting cases (e.g. for \mathbf{R} and $\bar{\mathbf{R}}$ given above) is the subject of further investigation.

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