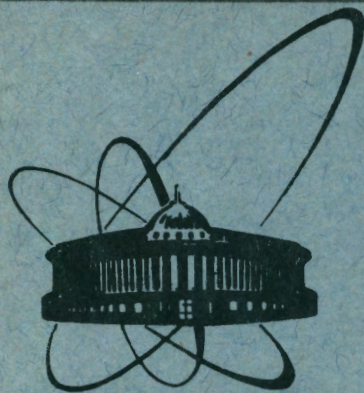


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$H^+H^-$  INTERACTION UP TO HIGHER ORDERS OF  
PERTURBATION THEORY IN THE MODEL  
WITH TWO HIGGS DOUBLETS  
(COUNTER TERMS)

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The paper is devoted to a calculation of the counterterms needed for the renormalization of the  $H^+H^-$  amplitude performed on the mass shell and in the unitary gauge. This work is the continuation of our previous paper <sup>1/1</sup> where we presented the general expressions for  $\Gamma^{Ren.}$  and  $i\mathcal{M}^{Ren.}$

$$\Gamma_{H^+H^-A}^{c.t.}(p^2, q^2, t) = -ie(\mathcal{Z}_{H^+} - 1), \quad (1)$$

$$\Gamma_{H^+H^+Z}^{c.t.}(p^2, q^2, t) = -\frac{ig(2R-1)}{2\sqrt{R}} \left[ (\mathcal{Z}_{H^+} - 1) + \frac{1}{2}(\mathcal{Z}_Z - 1) + \frac{\delta q}{q} + \frac{2R+1}{2(2R-1)} \frac{\delta R}{R} \right] - ie \cdot \left[ (\mathcal{Z}_{H^+} - 1) + \frac{1}{2}(\mathcal{Z}_M - 1) - \frac{1}{2}(\mathcal{Z}_A - 1) \right] \quad (2)$$

$$\text{and } \Gamma_{H^+H^+H_{1,2}^0}^{c.t.}(p^2, q^2, t) = -igM_W \left[ \left\{ \frac{\cos(\beta-\alpha)}{\sin(\beta-\alpha)} \right\} \mp \frac{1}{2R} \right] \cdot \cos 2\beta \left\{ \frac{\cos(\beta+\alpha)}{\sin(\beta+\alpha)} \right\} \cdot \left[ (\mathcal{Z}_{H^+} - 1) + \frac{1}{2}(\mathcal{Z}_{H_{1,2}^0} - 1) + \frac{\delta q}{q} + \frac{1}{2} \frac{\delta M_W^2}{M_W^2} \right] - igM_W \left[ \mp \frac{1}{2R} \cos 2\beta \left\{ \frac{\cos(\beta+\alpha)}{\sin(\beta+\alpha)} \right\} \right] \cdot \frac{\delta R}{R} - igM_W \left[ \left\{ \frac{\sin(\beta-\alpha)}{\cos(\beta-\alpha)} \right\} \mp \frac{1}{2R} \cos 2\beta \left\{ \frac{\sin(\beta+\alpha)}{\cos(\beta+\alpha)} \right\} \right] \cdot \left[ (\mathcal{Z}_{H^+} - 1) + \frac{1}{2}(\mathcal{Z}_{M_{2,M_1}} - 1) + \frac{\delta q}{q} + \frac{1}{2} \frac{\delta M_W^2}{M_W^2} \right] - igM_W \left[ \mp \frac{1}{2R} \cos 2\beta \left\{ \frac{\sin(\beta+\alpha)}{\cos(\beta+\alpha)} \right\} \right] \frac{\delta R}{R} \quad (3)$$

Besides

$$i\mathcal{M}^{\text{c.t.}} = -\frac{g^2 \cos^2 2\beta}{2 \cos^2 \theta_W} \left[ 2 \frac{\delta g}{g} + 2(\mathcal{Z}_{H^\pm} - 1) - \frac{\delta R}{R} \right]^{(4)}$$

So we need to find the next counter terms:  $\delta M_W^2/M_W^2$ ,  $\delta M_{H^\pm}^2/M_{H^\pm}^2$ ,

$$\delta M_{\mathcal{Z}}^2/M_{\mathcal{Z}}^2, \mathcal{Z}_A - 1, \mathcal{Z}_Z - 1, \mathcal{Z}_M - 1, \mathcal{Z}_{H_1^0} - 1, \mathcal{Z}_{H_2^0} - 1, \mathcal{Z}_{H^\pm} - 1, \mathcal{Z}_{M_1} - 1, \mathcal{Z}_{M_2} - 1.$$

The mass counter terms are

$$\begin{aligned} \frac{\delta M_W^2}{M_W^2} &= \frac{ig^2}{16\pi^2} \left\{ \left[ \frac{34}{3} - \frac{3}{2} \frac{1}{R} - \frac{1}{3} M_f - r_+ - r_1 - r_2 + \frac{1}{M_W^2} \times \right. \right. \\ &\times \left. \left. \text{Tr } m_i^2 \right] P + \text{finite part of SM with the exception of Higgs} \right. \\ \text{terms} &- \frac{5}{18} + \frac{r_+}{2} + \frac{r_1(1 + \sin^2(\beta - \alpha)) + r_2(1 + \cos^2(\beta - \alpha))}{4} - \\ &- \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{12} - \frac{r_1^2 \cos 2(\beta - \alpha) - r_2^2 \cos 2(\beta - \alpha) - r_3^2}{12} - \\ &- r_+ \frac{r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha) - r_3}{6} + \left( \frac{r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha) - r_3}{12} - \right. \\ &- \left. \frac{1}{2} \right) r_+ \ln r_+ + \left( \frac{r_+ - r_1 - 1}{12} \sin^2(\beta - \alpha) + \frac{r_1}{12} \cos^2(\beta - \alpha) - \frac{1}{4} \right) \times \\ &\times r_1 \ln r_1 + \left( \frac{r_+ - r_2 - 1}{12} \cos^2(\beta - \alpha) + \frac{r_2}{12} \sin^2(\beta - \alpha) - \frac{1}{4} \right) r_2 \ln r_2 + \\ &+ \left( \frac{r_3 - r_+ + 1}{12} - \frac{1}{4} \right) r_3 \ln r_3 - \cos^2(\beta - \alpha) \left( 1 - \frac{1}{3} r_1 + \frac{1}{12} r_1^2 \right) \\ &\times \bar{I}_0(-M_W^2, M_W^2, M_{H_1^0}^2) - \sin^2(\beta - \alpha) \left( 1 - \frac{1}{3} r_2 + \frac{1}{12} r_2^2 \right) \bar{I}_0(-M_W^2, \end{aligned}$$

$$M_{W, H_2^0}^2) + \frac{(r_+ - r_1 - 1)^2 - 4r_1}{12} \sin^2(\beta - \alpha) \bar{I}_0(-M_W^2, M_{H^\pm}^2, M_{H_1^0}^2) +$$

$$+ \cos^2(\beta - \alpha) \bar{I}_0(-M_W^2, M_{H^\pm}^2, M_{H_2^0}^2) \frac{(r_+ - r_2 - 1)^2 - 4r_2}{12} + \frac{1}{3} r_3 \cdot$$

$$\times \bar{I}_0(-M_W^2, M_{H^\pm}^2, M_{H_3^0}^2) \} \quad (5)$$

and

$$\frac{\delta M_{\bar{Z}}^2}{M_{\bar{Z}}^2} = \frac{ig^2}{16\pi^2} \left\{ \left[ -\frac{5}{3} + \frac{40}{3} R - \frac{13}{6} \frac{1}{R} - \frac{1}{3} \left( 2 - \frac{1}{R} \right) N_f - \right. \right.$$

$$\left. - \frac{8}{3} \frac{(1-R)^2}{R} \text{Tr } Q_i^2 + \frac{1}{M_W^2} \text{Tr } m_i^2 \right] P + \text{finite part of SM}$$

with the exception of Higgs terms  $-\frac{2}{9} + \frac{2}{9} R - \frac{1}{6} \frac{1}{R} -$

$$- R \left( \frac{r_1^2 + r_2^2 + r_3^2}{12} - r_3 \frac{r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha)}{6} \right) +$$

$$+ \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{6} + r_+ \ln r_+ \left( -\frac{1}{3} + \frac{4}{3} R - \frac{4}{3} R^2 \right) -$$

$$- \frac{\ln R}{R} \left( \frac{1}{6} - R \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{12} \right) + r_1 \ln r_1 \left( \frac{\sin^2(\beta - \alpha)}{12} - \right.$$

$$\left. - \frac{1}{4} \right) + r_2 \ln r_2 \left( \frac{\cos^2(\beta - \alpha)}{12} - \frac{1}{4} \right) + \frac{r_1^2 R}{12} \ln r_1 + \frac{r_2^2 R}{12} \ln r_2 -$$

$$- r_3 \ln r_3 \left( \frac{1}{6} - \frac{r_3 R}{12} + R \frac{r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha)}{12} \right) - r_3 R \cdot$$

$$\times \frac{r_1 \sin^2(\beta - \alpha) \ln r_1 + r_2 \cos^2(\beta - \alpha) \ln r_2}{12} + \left( \frac{1}{3} - \frac{1}{12} \frac{1}{R} - \frac{1}{3} R + \right.$$

$$+ \frac{1}{3} r_+ - \frac{4}{3} r_+ R + \frac{4}{3} r_+ R^2) \bar{I}_0(-M_{\bar{Z}}^2, M_{H^\pm}^2, M_{H^\pm}^2) + \sin^2(\beta - \alpha) \cdot$$

$$\cdot \left( \frac{1}{3} r_1 - \frac{(r_3 R - r_1 R - 1)^2}{12 R} \right) \bar{I}_0(-M_{\bar{Z}}^2, M_{H_1^0}^2, M_{H_3^0}^2) + \cos^2(\beta - \alpha) \cdot$$

$$\cdot \left( \frac{1}{3} r_2 - \frac{(r_3 R - r_2 R - 1)^2}{12 R} \right) \bar{I}_0(-M_{\bar{Z}}^2, M_{H_2^0}^2, M_{H_3^0}^2) - \cos^2(\beta - \alpha) \cdot$$

$$\cdot \left( \frac{1}{R} - \frac{1}{3} r_1 + \frac{1}{12} r_1^2 R \right) \bar{I}_0(-M_{\bar{Z}}^2, M_{\bar{Z}}^2, M_{H_1^0}^2) - \sin^2(\beta - \alpha) \left( \frac{1}{R} - \right.$$

$$\left. - \frac{1}{3} r_2 + \frac{1}{12} r_2^2 R \right) \bar{I}_0(-M_{\bar{Z}}^2, M_{\bar{Z}}^2, M_{H_2^0}^2) \} \quad (6)$$

The counter terms for a renormalization of the Higgs ( $H^\pm$ ) self-energy are

$$\frac{\delta M_{H^\pm}^2}{M_{H^\pm}^2} = \frac{ig^2}{16\pi^2} \left\{ \left[ \frac{1}{R} \left( \frac{3}{2} - \cos^2 2\beta \right) - (r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha)) - \right. \right.$$

$$\left. - r_3 \right] - \frac{1}{r_+} \left( 1 + \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{2} - \frac{r_1^2 \sin^2(\beta - \alpha) + r_2^2 \cos^2(\beta - \alpha) - r_3^2}{2} \right) \cdot$$

$$+ \frac{1}{r_+ R} \left( 6 + 2 \cos^2 2\beta + \frac{(r_1 - r_2) \cos 2\beta \cos 2\alpha + r_3 \cos^2 2\beta}{4} \right) - \frac{1}{r_+ R^2} \cdot$$

$$\cdot \left( \frac{3}{2} + \frac{1}{2} \cos^2 2\beta \right) - \text{Tr} \frac{m_u^2 \text{tg}^2 \beta + m_d^2 \text{tg}^2 \beta}{M_W^2} + \frac{2}{M_W^2 M_{H^\pm}^2} \text{Tr} (m_u^2 m_d^2) +$$

$$+ (m_u^2 + m_d^2) (m_u^2 \text{tg}^2 \beta + m_d^2 \text{tg}^2 \beta) \left. \right] P + \frac{1}{R} \left( \frac{1}{4} + \frac{\cos^2 2\beta}{2} \right) +$$

$$+ \frac{2}{r_+} \left( 1 - \frac{1}{R} + \frac{1}{4R^2} \right) \ln R - \frac{1}{R} \left( \frac{1}{4} + \frac{\cos^2 2\beta}{2} \right) \ln r_+ +$$

$$+ \frac{1}{r_+} \left( \frac{1}{2} + \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{4} - \frac{r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha) - r_3}{2} \right) \cdot$$

$$\begin{aligned}
& - \frac{r_1 \ln r_1 \cos 2(\beta-\alpha) - r_2 \ln r_2 \cos 2(\beta-\alpha) + r_3 \ln r_3}{4} - \frac{1}{r_+ R} \\
& \times \left( \frac{(r_1 - r_2) \cos 2\beta \cos 2\alpha + r_3 \cos^2 2\beta}{8} - \frac{(r_1 \ln r_1 - r_2 \ln r_2) \cos 2\beta \cos 2\alpha}{8} \right. \\
& - \frac{r_3 \ln r_3 \cos^2 2\beta}{8} \left. \right) - \frac{1}{M_W^2} \text{Tr} \left( \frac{m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta}{4} + \frac{m_u^2 m_d^2}{M_{H^\pm}^2} + \right. \\
& + \frac{m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta}{2M_{H^\pm}^2} \left( m_u^2 \ln \frac{m_u^2}{M_W^2} + m_d^2 \ln \frac{m_d^2}{M_W^2} \right) \left. \right) + 4(1-R) \\
& \times \bar{I}_0(-M_{H^\pm}^2, 0, M_{H^\pm}^2) - \left( 4 - 4R - \frac{1}{R} + \frac{1}{r_+} - \frac{1}{r_+ R} + \frac{1}{4r_+ R^2} \right) \\
& \times \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) + \left( \frac{(r_1 - r_+ - 1)^2}{r_+} - 4 \right) \frac{\sin^2(\beta-\alpha)}{4} \bar{I}_0(-M_{H^\pm}^2, \\
& M_W^2, M_{H^\pm}^2) + \left( \frac{(r_2 - r_+ - 1)^2}{r_+} - 4 \right) \frac{\cos^2(\beta-\alpha)}{4} \bar{I}_0(-M_{H^\pm}^2, M_W^2, M_{H^\pm}^2) \\
& - \left( \frac{1}{r_+} - 1 \right) \bar{I}_0(-M_{H^\pm}^2, M_W^2, M_{H^\pm}^2) - \frac{1}{r_+} \left( \cos(\beta-\alpha) - \frac{\cos 2\beta \cos(\beta+\alpha)}{2R} \right) \\
& \times \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) - \frac{1}{r_+} \left( \sin(\beta-\alpha) + \frac{\cos 2\beta \sin(\beta+\alpha)}{2R} \right) \bar{I}_0(-M_{H^\pm}^2, \\
& M_{H^\pm}^2, M_{H^\pm}^2) - \frac{1}{2M_W^2} \text{Tr} \left[ m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta - \frac{4m_u^2 m_d^2}{M_{H^\pm}^2} - \right. \\
& \left. - \frac{(m_u^2 + m_d^2)(m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta)}{M_{H^\pm}^2} \right] \bar{I}_0(-M_{H^\pm}^2, m_u^2, m_d^2) \left. \right\}. \quad (7)
\end{aligned}$$

and

$$\mathcal{Z}_{H^\pm} - 1 = \frac{ig^2}{16\pi^2} \left\{ \left[ 2 - 2R - \frac{3}{2} \frac{1}{R} + r_1 \sin^2(\beta-\alpha) + r_2 \cos^2(\beta-\alpha) \right. \right.$$

$$\begin{aligned}
& - r_3 + \frac{1}{M_W^2} \text{Tr}(m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta) \left. \right] P + 4(1-R) P_{IR} - 5 + 5R \\
& + \frac{1}{4R} + \left( R - 1 + \frac{1}{4R} \right) \ln R + 2(1-R) \ln r_+ + \frac{1}{4M_W^2} \\
& \times \text{Tr}(m_u^2 \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta) - 2(1-R) \bar{I}_0(-M_{H^\pm}^2, 0, M_{H^\pm}^2) - \\
& - 2\left(R - 1 + \frac{1}{4R}\right) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) + \frac{\sin^2(\beta-\alpha)}{2} (1 - r_+ + \\
& + r_1) \bar{I}_0(-M_{H^\pm}^2, M_W^2, M_{H^\pm}^2) + \frac{\cos^2(\beta-\alpha)}{2} (1 - r_+ + r_2) \bar{I}_0(-M_{H^\pm}^2, \\
& M_W^2, M_{H^\pm}^2) - \frac{1 - r_+ + r_3}{2} \bar{I}_0(-M_{H^\pm}^2, M_W^2, M_{H^\pm}^2) + \frac{1}{2M_W^2} \text{Tr}(m_u^2 \\
& \times \text{ctg}^2 \beta + m_d^2 \text{tg}^2 \beta) \bar{I}_0(-M_{H^\pm}^2, m_u^2, m_d^2) + \left( 4R - 4 + \frac{1}{R} - \frac{1}{r_+} + \right. \\
& + \frac{1}{r_+ R} - \frac{1}{4r_+ R^2} \left. \right) M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} + \frac{\sin^2(\beta-\alpha)}{4} \times \\
& \times \left( \frac{(r_1 - r_+ - 1)^2}{r_+} - 4 \right) M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_W^2, M_{H^\pm}^2)}{\partial q^2} + \frac{\cos^2(\beta-\alpha)}{4} \times \\
& \times \left( \frac{(r_2 - r_+ - 1)^2}{r_+} - 4 \right) M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_W^2, M_{H^\pm}^2)}{\partial q^2} - \left( \frac{1}{r_+} - 1 \right) \\
& \times M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_W^2, M_{H^\pm}^2)}{\partial q^2} - \frac{1}{r_+} \left( \cos(\beta-\alpha) - \frac{1}{2R} \cos 2\beta \times \right. \\
& \times \cos(\beta+\alpha) \left. \right)^2 M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} - \frac{1}{r_+} \left( \sin(\beta-\alpha) + \right. \\
& + \frac{1}{2R} \cos 2\beta \sin(\beta+\alpha) \left. \right)^2 M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} -
\end{aligned}$$

$$-\frac{r_+}{2} \text{Tr} \left[ m_u^2 c t_\gamma^2 \beta + m_d^2 t_\gamma^2 \beta - \frac{(m_u^2 + m_d^2)(m_u^2 c t_\gamma^2 \beta + m_d^2 t_\gamma^2 \beta)}{M_{H^\pm}^2} \right. \\ \left. - \frac{4m_u^2 m_d^2}{M_{H^\pm}^2} \right] \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, m_u^2, m_d^2)}{\partial q^2} \quad (8)$$

The other counter terms are

$$\mathcal{Z}_A - 1 = \frac{ie^2}{16\pi^2} \left\{ \left[ -\frac{40}{3} + \frac{8}{3} \text{Tr} Q_i^2 \right] P + \frac{1}{3} \ln r_+ + \frac{2}{3} (1 + \text{Tr} Q_i^2) + \frac{4}{3} \text{Tr} Q_i^2 \ln \frac{m_i^2}{M_W^2} \right\} \quad (9)$$

$$\mathcal{Z}_Z - 1 = \frac{ig^2}{16\pi^2} \left\{ \left[ 4 - \frac{40}{3} R + \frac{1}{R} + \frac{1}{3} (2 - \frac{1}{R}) N_f + \frac{8}{3} \frac{(1-R)^2}{R} \text{Tr} Q_i^2 \right] P + \text{finite part of SM with the exception} \right.$$

of Higgs terms

$$+ \frac{2}{9} - \frac{2}{9} R - \frac{1}{4} \frac{1}{R} - \frac{R}{12} (r_1^2 + r_2^2 + r_3^2) -$$

$$\frac{1}{12} \ln R \left( \frac{1}{R} - r_1 \cos^2(\beta - \alpha) - r_2 \sin^2(\beta - \alpha) \right) + \frac{r_1 \cos^2(\beta - \alpha) + r_2 \sin^2(\beta - \alpha)}{6}$$

$$- \frac{r_1 \ln r_1 \cos^2(\beta - \alpha) + r_2 \ln r_2 \sin^2(\beta - \alpha)}{12} + \frac{R}{12} (r_1^2 \ln r_1 + r_2^2 \ln r_2 +$$

$$+ r_3^2 \ln r_3) + \frac{r_3 R}{6} (r_1 \sin^2(\beta - \alpha) + r_2 \cos^2(\beta - \alpha)) - \frac{r_3 R}{12} (r_1 \ln r_1$$

$$\times \sin^2(\beta - \alpha) + r_2 \ln r_2 \cos^2(\beta - \alpha)) - \frac{r_3 R \ln r_3}{12} (r_1 \sin^2(\beta - \alpha) +$$

$$+ r_2 \cos^2(\beta - \alpha)) + \frac{1}{12} (4R - 4 + \frac{1}{R}) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) -$$

$$- \frac{\sin^2(\beta - \alpha)}{12} \left( \frac{1}{R} - R(r_3 - r_1)^2 \right) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) - \frac{\cos^2(\beta - \alpha)}{12}$$

$$\times \left( \frac{1}{R} - R(r_3 - r_2)^2 \right) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) + \frac{r_1}{6} \left( 1 - \frac{r_1 R}{2} \right) \cos^2(\beta - \alpha) \\ \times \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) + \frac{r_2}{6} \left( 1 - \frac{r_2 R}{2} \right) \sin^2(\beta - \alpha) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) \\ + \frac{1}{3} \left( 1 - R + \frac{r_+}{2} - \frac{1}{4R} - 2r_+ R - 2r_+ R^2 \right) M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} \\ + \frac{\sin^2(\beta - \alpha)}{3} \left[ r_1 - \frac{R}{4} (r_3 - r_1 - \frac{1}{R})^2 \right] M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} - \\ + \frac{\cos^2(\beta - \alpha)}{3} \left[ r_2 - \frac{R}{4} (r_3 - r_2 - \frac{1}{R})^2 \right] M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} - \\ - \frac{\cos^2(\beta - \alpha)}{R} \left[ 1 - \frac{r_1 R}{3} + \frac{r_1^2 R^2}{12} \right] M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} - \\ - \frac{\sin^2(\beta - \alpha)}{R} \left[ 1 - \frac{r_2 R}{3} + \frac{r_2^2 R^2}{12} \right] M_{H^\pm}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} \quad (10)$$

The counter term  $\sqrt{\mathcal{Z}_M}$  is calculated from the ZA - mixing diagrams.

$$\sqrt{\mathcal{Z}_M} = \frac{ieg}{16\pi^2} R^{-1/2} \left\{ \left[ \frac{44}{3} R - \frac{8}{3} - \frac{1}{6} \frac{1}{R} - \frac{1}{3} N_f + \right. \right.$$

$$\left. + \frac{8}{3} (1 - R) \text{Tr} Q_i^2 \right] P + \text{finite part of SM with the exception} \quad (11)$$

of Higgs terms

$$+ \frac{1}{9} - \frac{2}{9} R - \frac{2}{3} R r_+ \ln r_+ (1 - 2R) -$$

$$- \left( \frac{1}{6} - \frac{1}{3} R + \frac{2}{3} r_+ R - \frac{4}{3} r_+ R^2 \right) \bar{I}_0(-M_{H^\pm}^2, M_{H^\pm}^2, M_{H^\pm}^2) \left. \right\}$$

For neutral Higgs particles we have

$$\begin{aligned}
 \sum_{H_{1,2}^0} -1 &= \frac{ig^2}{16\pi^2} \left\{ \left[ \frac{\sin^2(\beta-\alpha)}{\cos^2(\beta-\alpha)} \right] \times \left( \frac{3}{2} - \frac{3}{2R} + r_+ - r_3 \right) - \right. \\
 &\left. \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} \left( 3 + \frac{3}{2R} - \frac{3}{2} r_{1,2} \right) + \frac{1}{4M_W^2} \text{Tr} \left( \frac{m_u^2}{\sin^2\beta} \times \right. \right. \\
 &\left. \left. \left\{ \frac{\sin^2\alpha}{\cos^2\alpha} \right\} + \frac{m_d^2}{\cos^2\beta} \left\{ \frac{\cos^2\alpha}{\sin^2\alpha} \right\} \right) \right] P - \frac{1}{4} + \frac{3}{4} \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} + \\
 &+ \frac{1}{4R} (1 + \ln R) + \frac{1}{16M_W^2} \text{Tr} \left( \frac{m_u^2}{\sin^2\beta} \left\{ \frac{\sin^2\alpha}{\cos^2\alpha} \right\} + \right. \\
 &+ \left. \frac{m_d^2}{\cos^2\beta} \left\{ \frac{\cos^2\alpha}{\sin^2\alpha} \right\} \right) + \frac{1}{2} \left\{ \frac{\sin^2(\beta-\alpha)}{\cos^2(\beta-\alpha)} \right\} (r_+ - r_{1,2} + 1) \bar{I}_0(-M_{H_{1,2}^0}^2, \\
 &M_W^2, M_{H^\pm}^2) - \frac{1}{2} \left\{ \frac{\sin^2(\beta-\alpha)}{\cos^2(\beta-\alpha)} \right\} (r_3 - r_{1,2} + \frac{1}{R}) \bar{I}_0(-M_{H_{1,2}^0}^2, \\
 &M_{\tilde{X}}^2, M_{H_3^0}^2) - \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} \left( 1 - \frac{1}{2} r_{1,2} \right) \bar{I}_0(-M_{H_{1,2}^0}^2, M_W^2, M_W^2) - \\
 &- \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} \frac{1}{2R} \left( 1 - \frac{1}{2} r_{1,2} R \right) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{\tilde{X}}^2, M_{\tilde{X}}^2) + \frac{1}{8M_W^2} \\
 &\times \text{Tr} \left[ \frac{m_u^2}{\sin^2\beta} \left\{ \frac{\sin^2\alpha}{\cos^2\alpha} \right\} \bar{I}_0(-M_{H_{1,2}^0}^2, m_u^2, m_u^2) \right] + \frac{1}{8M_W^2} \text{Tr} \left[ \frac{m_d^2}{\cos^2\beta} \times \right. \\
 &\left. \left\{ \frac{\cos^2\alpha}{\sin^2\alpha} \right\} \bar{I}_0(-M_{H_{1,2}^0}^2, m_d^2, m_d^2) \right] + \left( \frac{(r_+ - r_{1,2} - 1)^2}{4r_{1,2}} - 1 \right) \times M_{H_{1,2}^0}^2 \\
 &\times \left\{ \frac{\sin^2(\beta-\alpha)}{\cos^2(\beta-\alpha)} \right\} \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_W^2, M_{H^\pm}^2)}{\partial q^2} - \left( \frac{(r_3 - r_{1,2} - 1/R)^2}{4r_{1,2}} \right)
 \end{aligned}$$

$$\begin{aligned}
 &- \frac{1}{R} \left\{ \frac{\sin^2(\beta-\alpha)}{\cos^2(\beta-\alpha)} \right\} M_{H_{1,2}^0}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{\tilde{X}}^2, M_{\tilde{X}}^2)}{\partial q^2} + \left( 1 - \right. \\
 &- \left. \frac{r_{1,2}}{4} - 3 \frac{1}{r_{1,2}} \right) M_{H_{1,2}^0}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_W^2, M_W^2)}{\partial q^2} \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} \\
 &+ \left( \frac{1}{2R} - \frac{r_{1,2}}{8} - \frac{3}{2} \frac{1}{R^2 r_{1,2}} \right) \left\{ \frac{\cos^2(\beta-\alpha)}{\sin^2(\beta-\alpha)} \right\} M_{H_{1,2}^0}^2 \times \\
 &\times \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{\tilde{X}}^2, M_{\tilde{X}}^2)}{\partial q^2} - \frac{1}{r_{1,2}} \left[ \left\{ \frac{\cos(\beta-\alpha)}{\sin(\beta-\alpha)} \right\} \mp \frac{\cos 2\beta}{2R} \right. \\
 &\times \left. \left\{ \frac{\cos(\beta+\alpha)}{\sin(\beta+\alpha)} \right\} \right]^2 M_{H_{1,2}^0}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{H^\pm}^2, M_{H^\pm}^2)}{\partial q^2} - \frac{9}{8} \frac{1}{R^2 r_{1,2}} \times \\
 &\times \cos^2 2\alpha \left\{ \frac{\cos^2(\beta+\alpha)}{\sin^2(\beta+\alpha)} \right\} M_{H_{1,2}^0}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2)}{\partial q^2} - \\
 &- \frac{1}{8} \frac{1}{R^2 r_{1,2}} \left[ 2 \sin 2\alpha \left\{ \frac{\sin(\beta+\alpha)}{\cos(\beta+\alpha)} \right\} \mp \cos 2\alpha \left\{ \frac{\cos(\beta+\alpha)}{\sin(\beta+\alpha)} \right\} \right]^2 \times M_{H_{1,2}^0}^2 \\
 &\times \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2)}{\partial q^2} - \frac{1}{4} \frac{1}{R^2 r_{1,2}} \left[ 2 \sin 2\alpha \left\{ \frac{\cos(\beta+\alpha)}{\sin(\beta+\alpha)} \right\} \right. \\
 &\left. \mp \cos 2\alpha \left\{ \frac{\sin(\beta+\alpha)}{\cos(\beta+\alpha)} \right\} \right]^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2, M_{H_{1,2}^0}^2)}{\partial q^2} - \frac{1}{8} \frac{\cos^2 2\beta}{R^2 r_{1,2}} \\
 &\times \left\{ \frac{\cos^2(\beta+\alpha)}{\sin^2(\beta+\alpha)} \right\} M_{H_{1,2}^0}^2 \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, M_{H_3^0}^2, M_{H_3^0}^2)}{\partial q^2} - \frac{r_{1,2}}{8 \sin^2 \beta} \left\{ \frac{\sin^2\alpha}{\cos^2\alpha} \right\} \\
 &\times \text{Tr} \left( m_u^2 - \frac{2m_u^4}{M_{H_{1,2}^0}^2} \right) \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, m_u^2, m_u^2)}{\partial q^2} - \frac{r_{1,2}}{8 \cos^2 \beta} \left\{ \frac{\cos^2\alpha}{\sin^2\alpha} \right\}
 \end{aligned}$$

$$\text{Tr} \left( m_d^2 - \frac{2m_d^2}{M_{H_{1,2}^0}^2} \right) \frac{\partial \bar{I}_0(q^2 = -M_{H_{1,2}^0}^2, m_d^2, m_d^2)}{\partial q^2} \quad (12)$$

and for  $H_1^0 H_2^0$ -mixing the appropriate counter term is

$$\begin{aligned} \sqrt{Z_{M_1, M_2}} &= \frac{ig^2}{16\pi^2} \left\{ \left[ -\frac{11}{4} \sin 2(\beta-\alpha) + \frac{3}{8} r_{1,2} \sin 2(\beta-\alpha) \right] \right. \\ &+ \frac{1}{r_{1,2}} \left( 4 \sin 2(\beta-\alpha) + \frac{r_+}{4} (\sin 2\beta \cos 2\alpha - 3 \sin 2\alpha \cos 2\beta) \right. \\ &+ \left. \frac{r_3}{4} \sin 2(\beta-\alpha) \right) + \frac{1}{r_{1,2} R} \left( \frac{3}{2} r_+ \sin 2\alpha \cos 2\beta - \frac{3}{8} (r_1 - r_2) \right. \\ &\times \sin 4\alpha \left. \right) + \frac{1}{r_{1,2} R^2} \left( \frac{3}{2} \sin 2(\beta-\alpha) - \frac{1}{4} \sin 4\alpha \cos 2(\beta+\alpha) \right. \\ &\left. - \sin 2(\beta+\alpha) \left( \sin^2 2\alpha + \frac{\cos^2 2\alpha}{2} + \frac{3}{8} \cos^2 2\beta \right) \right) + \frac{\sin 2\alpha}{2M_W^2 \sin^2 \beta} \\ &\times \text{Tr} m_u^2 \left( 1 - \frac{6m_u^2}{M_{H_{1,2}^0}^2} \right) - \frac{\sin 2\alpha}{2M_W^2 \cos^2 \beta} \text{Tr} m_d^2 \left( 1 - \frac{6m_d^2}{M_{H_{1,2}^0}^2} \right) \Big] P_+ \\ &+ \frac{3}{8} \sin 2(\beta-\alpha) + \frac{1}{r_{1,2}} \left( \frac{9}{8} \sin 2(\beta-\alpha) + \frac{r_+}{4} (2 \sin 2\alpha \cos 2\beta \right. \\ &\left. - \sin 2\beta \cos 2\alpha) - \frac{r_+ \ln r_+}{8} (3 \sin 2\alpha \cos 2\beta - \sin 2\beta \cos 2\alpha) \right) \\ &+ \frac{1}{r_{1,2} R} \left( -\frac{r_+}{4} \cos 2\beta \sin 2\alpha + \frac{r_3}{8} (2 \sin 2\beta \cos 2\alpha - 3 \sin 2\alpha \right. \\ &\times \cos 2\beta) + \frac{r_+ \ln r_+}{4} \cos 2\beta \sin 2\alpha + \frac{r_3 \ln r_3}{8} \\ &\times (2 \sin 2\alpha \cos 2\beta - \sin 2\beta \cos 2\alpha) + \end{aligned}$$

$$\begin{aligned} &+ \frac{3}{16} \sin 4\alpha (r_1 - r_2) - \frac{3}{16} r_1 \ln r_1 \sin 4\alpha + \frac{3}{16} r_2 \ln r_2 \sin 4\alpha \Big) + \\ &+ \frac{1}{r_{1,2} R} \ln R \left( \frac{r_3}{8} \sin 2(\beta-\alpha) - \frac{1}{8R} \sin 2(\beta-\alpha) \right) + \frac{1}{r_{1,2} R^2} \frac{3}{8} \\ &\times \sin 2(\beta-\alpha) + \frac{\sin 2\alpha}{8M_W^2 \sin^2 \beta} \text{Tr} m_u^2 \left( 1 - \frac{2m_u^2}{M_{H_{1,2}^0}^2} \right) - \frac{\sin 2\alpha}{8M_W^2 \cos^2 \beta} \\ &\times \text{Tr} m_d^2 \left( 1 - \frac{2m_d^2}{M_{H_{1,2}^0}^2} \right) - \frac{\sin 2\alpha}{2M_W^2 \sin^2 \beta} \text{Tr} \frac{m_u^4}{M_{H_{1,2}^0}^2} \ln \frac{m_u^2}{M_W^2} + \\ &+ \frac{\sin 2\alpha}{2M_W^2 \cos^2 \beta} \text{Tr} \frac{m_d^4}{M_{H_{1,2}^0}^2} \ln \frac{m_d^2}{M_W^2} + \frac{1}{8} \left( \frac{(r_+ - r_{1,2} - 1)^2}{r_{1,2}} - 4 \right) \sin 2(\beta-\alpha) \\ &\times \bar{I}_0(-M_{H_{1,2}^0}^2, M_W^2, M_{H_1^0}^2) - \frac{1}{8} \left( \frac{(r_3 - r_{1,2} - 1/R)^2}{r_{1,2}} - \frac{4}{R} \right) \sin 2(\beta-\alpha) \\ &\times \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_2^0}^2, M_{H_3^0}^2) - \left( \frac{1}{2} - \frac{r_{1,2}}{8} - \frac{3}{2r_{1,2}} \right) \sin 2(\beta-\alpha) \bar{I}_0(-M_{H_{1,2}^0}^2, \\ &M_W^2, M_W^2) - \left( \frac{1}{4R} - \frac{r_{1,2}}{16} - \frac{3}{4r_{1,2} R^2} \right) \sin 2(\beta-\alpha) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_2^0}^2, \\ &M_{H_3^0}^2) + \frac{1}{r_{1,2}} \left( \cos(\beta-\alpha) \sin(\beta-\alpha) + \frac{1}{2R} \sin 2\alpha \cos 2\beta - \frac{1}{4R^2} \cos^2 2\beta \right. \\ &\times \cos(\beta+\alpha) \sin(\beta+\alpha) \left. \right) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_1^0}^2, M_{H_2^0}^2) - \frac{1}{8R^2 r_{1,2}} \cos^2 2\beta \\ &\times \cos(\beta+\alpha) \sin(\beta+\alpha) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_3^0}^2, M_{H_3^0}^2) - \frac{3 \cos 2\alpha}{8R^2 r_{1,2}} (2 \sin 2\alpha \\ &\times \cos^2(\beta+\alpha) + \frac{1}{2} \cos 2\alpha \sin^2(\beta+\alpha)) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_1^0}^2, M_{H_1^0}^2) + \frac{3 \cos 2\alpha}{8R^2 r_{1,2}} \times \end{aligned}$$



$$\begin{aligned}
& \left( 2 \sin 2\alpha \sin^2(\beta+\alpha) - \frac{1}{2} \cos 2\alpha \sin 2(\beta+\alpha) \right) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_1^0}^2, M_{H_2^0}^2) + \\
& - \frac{1}{4R^2 r_{1,2}^2} \left( 2 \sin^2 2\alpha \sin 2(\beta+\alpha) - \sin 4\alpha \cos 2(\beta+\alpha) - \frac{1}{2} \cos^2 2\alpha \cdot \right. \\
& \left. \times \sin 2(\beta+\alpha) \right) \bar{I}_0(-M_{H_{1,2}^0}^2, M_{H_1^0}^2, M_{H_2^0}^2) + \frac{\sin 2\alpha}{4M_W^2 \sin^2 \beta} \text{Tr } m_u^2 \left( 1 - \right. \\
& \left. - \frac{4m_u^2}{M_{H_{1,2}^0}^2} \right) \bar{I}_0(-M_{H_{1,2}^0}^2, m_u^2, m_u^2) - \frac{\sin 2\alpha}{4M_W^2 \cos^2 \beta} \text{Tr } m_d^2 \left( 1 - \frac{4m_d^2}{M_{H_{1,2}^0}^2} \right) \cdot \\
& \left. \times \bar{I}_0(-M_{H_{1,2}^0}^2, m_d^2, m_d^2) \right\} \quad (13)
\end{aligned}$$

The above expressions contain next abbreviations:  $r_+ = M_{H_t}^2 / M_W^2$ ,  $r_{H_i^0} = M_{H_i^0}^2 / M_W^2$ ,  $R = M_W^2 / M_Z^2$ ,  $\text{tg } \beta = v_2 / v_1$ .  $M_i$  are masses of particles in the SUSY model;  $v_1$  and  $v_2$  are expectation values;  $\alpha$  is a mixing angle in the neutral scalar sector.

So we calculated the counter terms for a renormalization in the unitary gauge. The analogous problem was solved in [2], but we didn't use any approximations as contrasted to the author of the ref. paper.

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$H^+H^-$ -взаимодействие с учетом высших порядков теории возмущений в модели с двумя хиггсовскими дублетами /контрчлены/

Вычислены контрчлены, необходимые для перенормировки однопетлевых диаграмм, которые используются для построения амплитуды  $H^+H^-$  взаимодействия в рамках модели с минимальным суперсимметричным расширением стандартной модели.

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$H^+H^-$  Interaction up to Higher Orders of Perturbation Theory in the Model with Two Higgs Doublets /Counter Terms/

The counter terms for the one-loop renormalization of the  $H^+H^-$  amplitude in the framework of the minimal SUSY extension of the standard model are calculated.

The investigation has been performed at the Laboratory of Theoretical Physics, JINR.

Communication of the Joint Institute for Nuclear Research, Dubna 1991