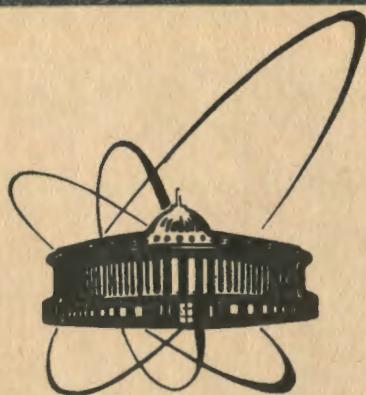


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ИНСТИТУТА  
ЯДЕРНЫХ  
ИССЛЕДОВАНИЙ  
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AN UNAMBIGUOUS WAY OF AN ESTIMATION  
OF VECTOR-MESON-NUCLEON TENSOR  
COUPLING CONSTANTS

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By treating the vector-meson resonances as point-like stable particles in some approximation, which are described by a field operator  $V_\mu(x)$ , the vector-meson-nucleon coupling of the form

$$L_{VNN} = ig_{VNN} \bar{\psi}_N \gamma_\mu \frac{1}{2} \vec{\tau} \psi_N \vec{V}_\mu + \frac{t_{VNN}}{2m_N} \bar{\psi}_N \sigma_{\mu\nu} \frac{1}{2} \vec{\tau} \psi_N (\partial_\mu \vec{V}_\nu - \partial_\nu \vec{V}_\mu) \quad (1)$$

with the so-called vector  $g_{VNN}$  and tensor  $t_{VNN}$  coupling constants is usually constructed from symmetry principles in meson field theory [1-3]. The coupling constants  $g_{VNN}$  and  $t_{VNN}$  have been evaluated numerically (for a review see Ref.[4]) in an analysis of  $\pi N$  and  $NN$  scattering and in an optimal description of the nucleon electromagnetic (e.m.) form factor (ff) data. Nevertheless, the agreement between the results of different methods (see e.g. Table 9.2 in Ref.[4]) was poor and, unexpectedly, it did not improve as the input data became more and more accurate. Therefore in this paper we propose a new unambiguous way of estimates at least of  $t_{VNN}$  which methodically differs from all previous ones and finally it can prefer some of the values presented in Table 9.2 of Ref.[4].

The method consists in an employment of a vector-meson-dominance (VMD) parametrization of isoscalar and isovector Pauli nucleon e.m. ff's [5]

$$F_2^{(s)}(t) = \sum_{s=\omega, \phi, \phi'} \frac{m_s^2 (t_{SNN}/f_s)}{m_s^2 - t} \quad (2a)$$

$$F_2^{(v)}(t) = \sum_{v=\rho, \rho', \rho''} \frac{m_v^2 (t_{VNN}/f_v)}{m_v^2 - t} \quad (2b)$$

where  $t = -Q^2$  is the four-momentum transfer squared, " $m$ " is

the mass of a vector meson under consideration and  $f_V$  is the universal vector-meson coupling constant determining the lepton width of a vector meson as follows

$$\Gamma(V \rightarrow e^+e^-) = (\alpha^2 m_V / 3) (f_V^2 / 4\pi)^{-1}. \quad (3)$$

The sums in (2a) and (2b) are saturated with only well established isoscalar and isovector mesons with quantum numbers of the photon, respectively.

The resonances  $\rho(770)$ ,  $\omega(783)$  and  $\phi(1020)$ , given by the Review of Particle Properties [6], are taken into account without any discussion. The  $\phi'(1680)$  meson, though in [6] it is declared as a not well established resonance, however, just after a new experimental study [7] of the reaction  $e^+e^- \rightarrow K^+K^-$  in the energy range  $1350 \text{ MeV} < \sqrt{s} < 2400 \text{ MeV}$  there is no doubt [8] in its existence now. Finally,  $\rho'(1450)$  and  $\rho''(1700)$  have been quite well confirmed [9-13] as the first and second radial excitations of the  $\rho(770)$  meson in the recent analysis of various processes.

Then by transforming (2a) and (2b) to a common denominator separately and forcing the resultant expressions to have the QCD (up to log-corrections) asymptotic behaviour [14,15]  $F_2^{(s)}(t) \equiv F_2^{(v)}(t) \rightarrow t^{-3} |t|^{-\omega}$  and also normalizations (see e.g. Ref.[15])  $F_2^{(s)}(0) = 1/2[\mu_P + \mu_N]$ ,  $F_2^{(v)}(0) = 1/2[\mu_P - \mu_N]$ , the following two systems of closed algebraic equations

$$\begin{aligned} (t_{\omega NN}/f_\omega) m_\omega^2 (m_\phi^2 + m_\psi^2) + (t_{\phi NN}/f_\phi) m_\phi^2 (m_\omega^2 + m_\psi^2) + (t_{\psi' NN}/f_{\psi'}) m_{\psi'}^2 (m_\omega^2 + m_\phi^2) &= 0 \\ (t_{\omega NN}/f_\omega) m_\omega^2 + (t_{\phi NN}/f_\phi) m_\phi^2 + (t_{\psi' NN}/f_{\psi'}) m_{\psi'}^2 &= 0 \quad (4) \\ (t_{\omega NN}/f_\omega) + (t_{\phi NN}/f_\phi) + (t_{\psi' NN}/f_{\psi'}) &= 1/2[\mu_P + \mu_N] \end{aligned}$$

$$\begin{aligned} (t_{\rho NN}/f_\rho) m_\rho^2 (m_{\rho'}^2 + m_{\rho''}^2) + (t_{\rho' NN}/f_{\rho'}) m_{\rho'}^2 (m_\rho^2 + m_{\rho''}^2) + (t_{\rho'' NN}/f_{\rho''}) m_{\rho''}^2 (m_\rho^2 + m_{\rho'}^2) &= 0 \\ (t_{\rho NN}/f_\rho) m_\rho^2 + (t_{\rho' NN}/f_{\rho'}) m_{\rho'}^2 + (t_{\rho'' NN}/f_{\rho''}) m_{\rho''}^2 &= 0 \quad (5) \\ (t_{\rho NN}/f_\rho) + (t_{\rho' NN}/f_{\rho'}) + (t_{\rho'' NN}/f_{\rho''}) &= 1/2[\mu_P - \mu_N] \end{aligned}$$

are obtained, from which the coupling constants  $t_{VNN}$  can be found. These are deduced by the solutions of (4) and (5), using masses of the  $\rho(770)$ ,  $\omega(783)$ ,  $\phi(1020)$ ,  $\phi'(1680)$  and anomalous magnetic moments of the proton  $\mu_p$  and the neutron  $\mu_n$  from the Review of Particle Properties [6], and masses of the  $\rho'(1450)$  and  $\rho''(1700)$  mesons from the paper by Donnachie and Mirsaie [9]. Numerical values of the universal  $\rho$ ,  $\omega$ ,  $\phi$  -meson coupling constants

$$f_\rho = 4.9905 \pm 0.1089; \quad f_\omega = 17.0550 \pm 0.2990; \quad f_\phi = 12.8866 \pm 0.2368; \quad (6)$$

are obtained from the corresponding lepton widths [6].

The lepton width of  $\phi'(1680)$  is, however, not established (see Ref.[6]) experimentally. Therefore the value

$$f_{\phi'} = 27.9030 \pm 1.5978 \quad (7)$$

has been estimated from measured combinations of several partial widths with the partial width of decay into  $e^+e^-$  and the total width (see Ref.[6]), by neglecting  $K_s^0 K\pi$  and  $3\pi$  channels at the same time. The universal  $\rho'$ ,  $\rho''$  -meson coupling constants

$$f_{\rho'} = 13.6491 \pm 0.9521; \quad f_{\rho''} = 22.4020 \pm 2.2728 \quad (8)$$

have been found from lepton widths estimated by Donnachie and O'Leeg [16].

As a result, the following numerical values of the vector-meson-nucleon tensor coupling constants are obtained

$$\begin{aligned} t_{\omega NN} &= -3.1739 \pm 0.0756; & t_{\rho NN} &= 16.0744 \pm 1.7595; \\ t_{\phi NN} &= 1.7461 \pm 0.0522; & t_{\rho' NN} &= 37.5100 \pm 9.0783; \\ t_{\phi' NN} &= 0.2651 \pm 0.0188; & t_{\rho'' NN} &= 30.9170 \pm 7.1628; \end{aligned} \quad (9)$$

where the errors have been evaluated from errors of all other quantities appearing in eqs.(4) and (5). For a comparison of our results with values estimated by other authors see Table .

One finds there an unexpected agreement between (9) and phenomenological estimates by Höhler et al [17,18], besides the value of  $t_{\rho' NN}^2/4\pi$ . However, the latter corresponds to an unexisting resonance  $\rho'$  (1250) used in Höhler's nucleon e.m. ff analysis, now.

In conclusion we would like to emphasise that the unambiguous confirmation of Höhler's results on  $t_{\rho NN}$  declared in Table particularly supports the conclusion of a recent paper by Weber [19] that relativistic effects in the constituent quark model of hadrons are large enough.

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Table

V	$t_{VNN}^2/4\pi$	Ref.
$\omega(783)$	$0.8016 \pm 0.0382$	this work
	1	[18]
	1.46	[20]
	$0.16 \pm 0.45$	[21]
$\phi(1020)$	$0.2426 \pm 0.0145$	this work
	0.2	[18]
$\phi'(1680)$	$0.0056 \pm 0.0008$	this work
$\rho(770)$	$20.5617 \pm 4.5014$	this work
	$20.5 \pm 2.1$	[17]
	10.92	[22]
	15.00	[20]
$\rho'(1450)$	$111.9655 \pm 4.1966$	this work
$\rho''(1700)$	$76.0650 \pm 35.7363$	this work

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