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ИНСТИТУТ  
ЯДЕРНЫХ  
ИССЛЕДОВАНИЙ  
ДУБНА

3555/2-80

4/8-80

E2-80-421

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ON "FAST" AND "SLOW" NEUTRINOS

Submitted to "Letters Nuovo Cimento".

1980

E2-80-421

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О "быстрых" и "медленных" нейтрино

Релятивистские нейтрино ведут себя как  $\nu_e, \nu_\mu, \nu_\tau \dots$ ; нерелятивистские - как  $\nu_1, \nu_2, \nu_3 / \nu_i$  - нейтрино с массой  $m_i$ .

Работа выполнена в Лаборатории теоретической физики ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна 1980

E2-80-421

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The possible existence of neutrino oscillations<sup>1/</sup> raises some questions. Lepton mixing implies finite neutrino masses<sup>2/</sup>. In such a case the field of neutrinos  $\nu_\ell$  ( $\ell = e, \mu, \tau \dots$ ) taking part in the weak charged current is

$$\nu_\ell = \sum_i U_{\ell i} \nu_i, \quad (1)$$

where  $\nu_i$  is the field operator of neutrino with definite mass  $m_i$ , and  $U$  is a unitary mixing matrix. For neutrinos we are working with in a Laboratory, the coherence conditions are fulfilled and the state vector of these "phenomenological" neutrinos is

$$|\nu_\ell \rangle_{\vec{p}} = \sum_i U_{\ell i}^* |\nu_i \rangle_{\vec{p}}, \quad (2)$$

where  $|\vec{p}| \gg m_i$  is the neutrino momentum. As is well-known relation (2) implies neutrino oscillations.

The phenomenon of neutrino oscillations has a well-known analog: the  $K^0 \rightarrow \bar{K}^0$  oscillations. The eigenstates of the strong interaction are  $K^0$  and  $\bar{K}^0$  just as, for the case of neutrinos, the eigenstates of the weak interaction are  $\nu_e, \nu_\mu, \nu_\tau \dots$ . The mass eigenstates are  $K_L$  and  $K_S$ , whereas in the case of neutrinos, the mass eigenstates are  $\nu_1, \nu_2, \nu_3, \dots$ . Since it is perfectly clear under what conditions one should talk either of  $K^0, \bar{K}^0$  or  $K_L, K_S$ , we will not discuss this point further. Less obvious is the question under what conditions neutrinos behave as  $\nu_1, \nu_2, \nu_3, \dots$ .

In this note we make an attempt to reply to this question, which in our opinion, has at least some instructional significance. As far as relativistic neutrinos are concerned (as already stated) the  $\nu_e, \nu_\mu, \nu_\tau, \dots$  are eigenstates of weak interaction, whereas neutrinos  $\nu_1, \nu_2, \nu_3 \dots$  have no physical meaning in the sense that they are not detectable as such. On the other hand if neutrino masses are different from zero there must exist nonrelativistic neutrinos (for example, the relic neutrino). Here the situation is inverted. In the non-relativistic case neutrinos  $\nu_e, \nu_\mu, \nu_\tau, \dots$  lose their physical

meaning, as they cannot undergo the usual weak interaction, changing them into charged leptons. They are not detectable as such. Nonrelativistic particles  $\nu_1, \nu_2, \nu_3, \dots$  with definite masses, on the contrary, acquire a definite physical meaning: they undergo the neutral current interaction without changing nature. Nonrelativistic neutrinos are just  $\nu_1, \nu_2, \nu_3, \dots$  \*.

In conclusion we stress again that if neutrinos are relativistic they will behave in general as  $\nu_e, \nu_\mu, \nu_\tau, \dots$  (in the conditions where  $\nu_1, \nu_2, \nu_3, \dots$  as such lose their physical meaning). If they are not relativistic they will behave in general as  $\nu_1, \nu_2, \nu_3, \dots$ , (whereas  $\nu_e, \nu_\mu, \nu_\tau, \dots$  as such lose their physical meaning). The function of (non-relativistic)  $\nu_1, \nu_2, \nu_3, \dots$  is gravity, the function of relativistic  $\nu_e, \nu_\mu, \nu_\tau, \dots$  is weak interaction.

We are grateful to V.N.Gribov for useful discussion.

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Received by Publishing Department  
on June 17 1980.

\* Of course, it is possible to imagine a situation where particles  $\nu_1$  (or  $\nu_2, \dots$ ) are relativistic. Such particles can be analysed in terms of  $\nu_e, \nu_\mu, \nu_\tau, \dots$  components but will not undergo oscillations.

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