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Nguyen Van Hieu

A LOWER BOUND
OF THE ELASTIC SCATTERING AMPLITUDE
AT FIXED MOMENTUM TRANSFER

ЛАБОРАТОРИЯ ТЕОРЕТИЧЕСКОЙ ФИЗИКИ

1968

Нгуен Ван Хьеу

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Оценка снизу для амплитуды упругого рассеяния при
фиксированной передаче импульса

На основе аналитических свойств амплитуды рассеяния, получена
нижняя граница убывания сечений в физической области $t < 0$.

Препринт Объединенного института ядерных исследований.
Дубна, 1968.

Nguyen Van Hieu

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A Lower Bound of the Elastic Scattering Amplitude
at Fixed Momentum Transfer

A lower bound of the cross section decrease is obtained in
the physical region $t < 0$ on the basis of the analytic properties of
the scattering amplitude.

Preprint. Joint Institute for Nuclear Research.
Dubna, 1968

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**A LOWER BOUND
OF THE ELASTIC SCATTERING AMPLITUDE
AT FIXED MOMENTUM TRANSFER**

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In this work we establish some restriction on the decrease of the elastic scattering amplitude $F(s, t)$ at $s \rightarrow \infty$ for the fixed $t < 0$, where s is the squared total energy in c.m.s., $t = -2k^2(1 - \cos \theta)$, k is the value of the three-dimensional momentum and θ is the scattering angle in this system of reference. For a series of processes the amplitude

$F(s, t)$ is an analytic function of t in the ellipse E_t (called the Martin ellipse) with the foci at $t = 0$ and $t = -4k^2$ and major semiaxis $a = 2k^2 + \gamma$, $\gamma > 0$, as this was shown by Martin^{/1/} and Sommer^{/2/}. We denote the imaginary part of $F(s, t)$ by $A(s, t)$. From the results of Jin and Martin^{/3/} it follows that

$$\max_{t \in E_t} |A(s, t)| \leq \text{const } s^{1+\varepsilon}, \quad s \rightarrow \infty \quad (1)$$

for some positive $\varepsilon < 1$. Instead of $A(s, t)$ it is convenient to consider the function

$$f(s, t) = \frac{A(s, t)}{A(s, 0)}, \quad (2)$$

which has the same analytic properties in t as $A(s, t)$ has. We suppose that the total cross section has the behavior $\text{const. } s^\beta$ at $s \rightarrow \infty$. Then

$$A(s, 0) \sim \text{const. } s^{1+\beta}, \quad s \rightarrow \infty,$$

and we have

$$\max_{t \in E_t} |f(s, t)| \leq \text{const } s^{\varepsilon-\beta}, \quad s \rightarrow \infty. \quad (3)$$

By the change of the variables $w = t + 2k^2$ we transform E_t into the ellipse E_w with the foci at $w = \pm c$, $c > 2k^2$, and with the same major semiaxis. The minor semiaxis is $b = \sqrt{a^2 - c^2}$. We introduce now an arbitrary positive number $\alpha < c$ and consider the ellipse E'_w with the foci at $w = \pm c'$, $c' = c - \alpha$, and with the minor semiaxis b' . This ellipse, of course, is contained inside E_w . The major semiaxis a' of E'_w is determined by equation

$$a'^2 = a^2 + c'^2 - c^2.$$

We shall choose c' (i.e. α) in such a manner that the points $w = \pm c$ (i.e. $t = 0$ and $t = -4k^2$) are contained inside E'_w . Then the physical region is contained completely inside E'_w . This condition is fulfilled if $a' > c'$, i.e. if the following inequality holds

$$(2k^2 + \gamma)^2 + (2k^2 - \alpha) - (2k^2)^2 > (2k^2)^2.$$

For large λ from this inequality we get

$$\alpha < \gamma. \quad (4)$$

By means of the conformal mapping

$$\xi = \frac{w + \sqrt{w^2 - c'^2}}{c'}$$

we transform the interval $[-c', c']$ in the w plane into the unite circumference in the ξ plane, following Cerulus and Martin^{4/}. In this conformal mapping the ellipse E'_w is transformed into the ring with the internal radius 1 and the external radius R ,

$$R = \frac{a' + \sqrt{a'^2 - c'^2}}{c'}. \quad (5)$$

The point $w = c$ (i.e. $t = 0$) is transformed to the point $\xi = \tau$,

$$\tau = \frac{c + \sqrt{c^2 - c'^2}}{c'}. \quad (6)$$

Denote by m the maximum of $|f(s, t)|$ in the interval $-c' \leq w \leq c'$ i.e. in the interval $-4k^2 + \alpha \leq t \leq -\alpha$, and by M that of $|f(s, t)|$ on the boundary of E'_t . According to the Hadamard's three circle theorem (see ref.^{5/}, theorem 5.32) we get

$$\ln |f(s, 0)| \leq \left(1 - \frac{\ln \tau}{\ln R}\right) \ln m + \frac{\ln \tau}{\ln R} \ln M.$$

Putting into (5) and (6) the values of a' , c , c' we get at the limit $\lambda \rightarrow \infty$

$$\frac{\ln \tau}{\ln R} \approx \sqrt{\alpha/\gamma}.$$

On the other hand $f(s, 0) = 1$. Therefore we have

$$m \geq \left(\frac{1}{M}\right)^{\frac{\sqrt{\alpha/\gamma}}{1 - \sqrt{\alpha/\gamma}}}.$$

For πN -scattering we have $\gamma = 4m_\pi^2$. Using the condition (3) we get now

$$\max_{-4k^2 + \alpha \leq t \leq -\alpha} |f(s, t)| \geq \text{const } s^{-(\epsilon - \rho)} \phi(\alpha), \quad (7)$$

where

$$\phi(\alpha) = \frac{\sqrt{\alpha/4m_\pi^2}}{1 - \sqrt{\alpha/4m_\pi^2}}. \quad (8)$$

If we assume that $f(s, t)$ is analytic and uniformly polynomially bounded in the whole cut plane (the Mandelstam representation) then we have

$$\max_{-4k^2 + \alpha \leq t \leq -\alpha} |f(s, t)| \geq \text{const } s^{-n} \psi(\alpha), \quad (9)$$

where

$$\psi(\alpha) = \frac{\left[1 - \left(1 + \frac{\alpha}{4m_\pi^2}\right)^{-1/2}\right]^{1/2}}{1 - \left[1 - \left(1 + \frac{\alpha}{4m_\pi^2}\right)^{-1/2}\right]^{1/2}}, \quad (10)$$

and n is such a constant that

$$|f(s, t)| \leq \text{const } s^n, \quad s \rightarrow \infty$$

for any t .

If the amplitude has the Regge behavior then the inequalities (7), (9) are the lower bounds of the corresponding Regge trajectories.

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