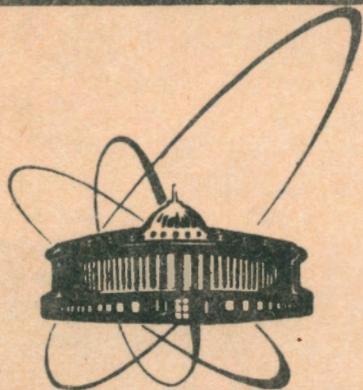


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СООБЩЕНИЯ
ОБЪЕДИНЕННОГО
ИНСТИТУТА
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ИССЛЕДОВАНИЙ
ДУБНА

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NEUTRON PRODUCTION IN LEAD TARGETS
BY HIGH-ENERGY LIGHT-MASS
HEAVY IONS

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INTRODUCTION

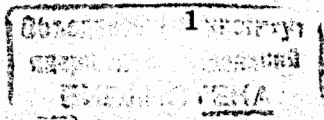
A problem of forming medium- and high-energy nuclear data file is an important task at present. A successful development of a lot of applications in science and practice depends on this problem. With reference of neutron data they have the most cost. But nowadays there is very poor experimental nuclear data set at the energy region above 100 MeV/nucleon. This prevents to produce the evaluated experimental data library. But recently the attempt to widen the ENDF/B-VI library to 1 GeV was undertaken by Pearlstein^{11/}. However, a comparison of these results with theoretical predictions showed the unaccountable discrepancy between the experiment and the theory^{2-4/}. So an accumulation of experimental information is the main task at present.

The systematic investigations of neutron production in thin and thick lead targets was performed by groups from KfK^{5-7/} and LANL^{8-10/} in the energy region of incident protons 100-1100 MeV. The proton spectra were also measured by KfK group. The beam of protons with energy above 1 GeV was used in the measurements of cross-sections of high-energy neutron and proton production by ITEP group^{11,12/}. But there are no any systematic investigations of neutron production in lead targets by relativistic heavy ions besides of the results of Cecil et al.^{13/} obtained with 640 and 710 MeV alpha-particles.

During last years we measured the neutron and charged particle energy distributions for thin and thick lead targets bombarded by 2.0 GeV and 2.55 GeV protons and d, α , ^6Li and ^{12}C ions with the energy of 1 GeV per nucleon. The main our purpose was to obtain the neutron data for the accelerator-based neutron spallation sources design. This research was supported by μCF program.

EXPERIMENTAL METHOD

The measurements were performed by the time-of-flight technique on our hadron spectrometer based on the JINR synchrophasotron. The schematic drawing of the target-detectors arrangement at the JINR synchrophasotron is given in fig.1. The block diagram of the electronics used in the TOF-spectrometer is given in fig.2. A neutron emission was studied at three different angles (at present after the reconstruction the spectrometer has four identical arms) using three types of detectors with stylben crystals $\phi 40 \times 10 \text{ mm}^3$ and $\phi 50 \times 50 \text{ mm}^3$ and plastic scintillators $\phi 120 \times 200 \text{ mm}^3$. This detector set overlapped very wide energy region from 0.2 MeV to about the beam energy containing more than 95% of all emitted neutrons. Using one of them we could study all stages of neutron emission — cascade,



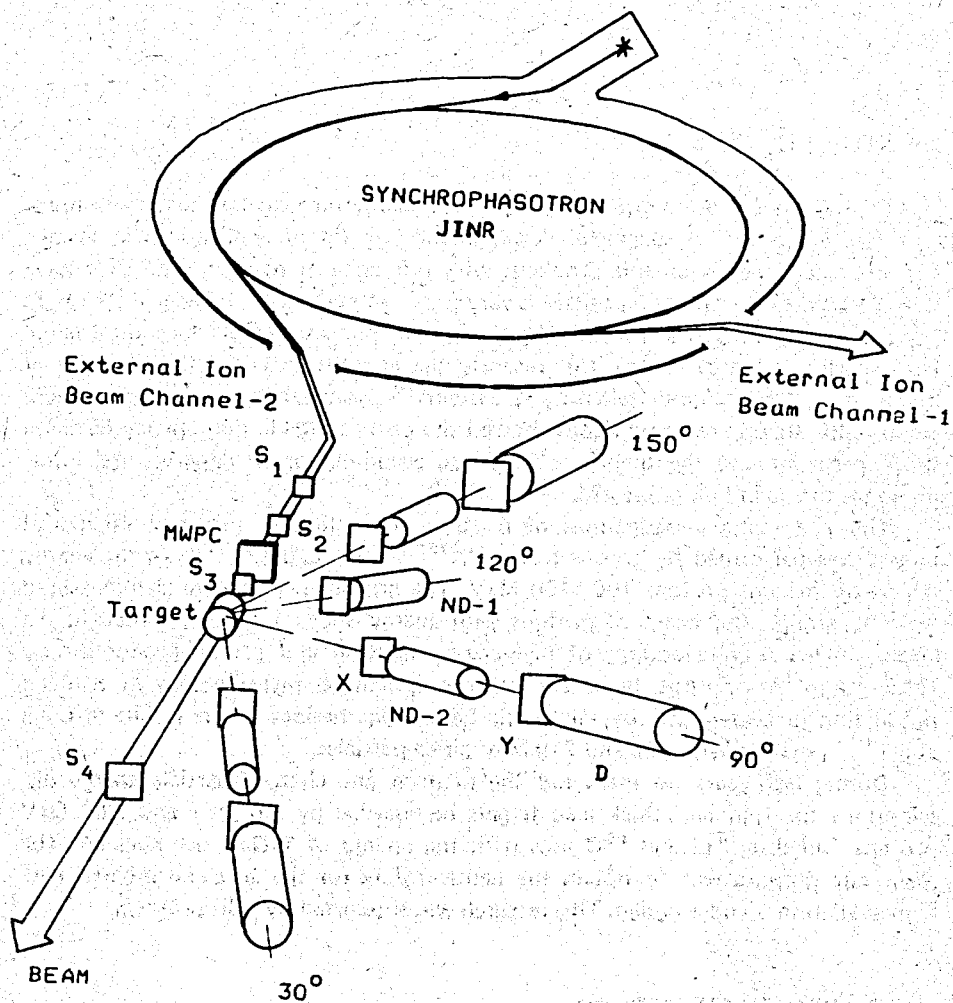


Fig. 1. Schematic drawing of the target-detectors arrangement at the JINR synchrotron: S1-S4 - beam counters and X, Y - veto-counters with thin plastic scintillators, ND-1 - detector with stylben $\phi 40 \times 10 \text{ mm}^3$, ND-2 - detector with stylben $\phi 50 \times 50 \text{ mm}^3$, D - detector with plastic scintillator $\phi 120 \times 200 \text{ mm}^3$.

pre-equilibrium and equilibrium. Thin 3 or 5 mm plastic scintillators as veto-counters for charged particles were placed in front of all detectors. For stylben detectors n- γ pulse-shape discrimination was employed. A separation of neutron- and gamma-events was performed in a stage of off-line data analysis. As example, two

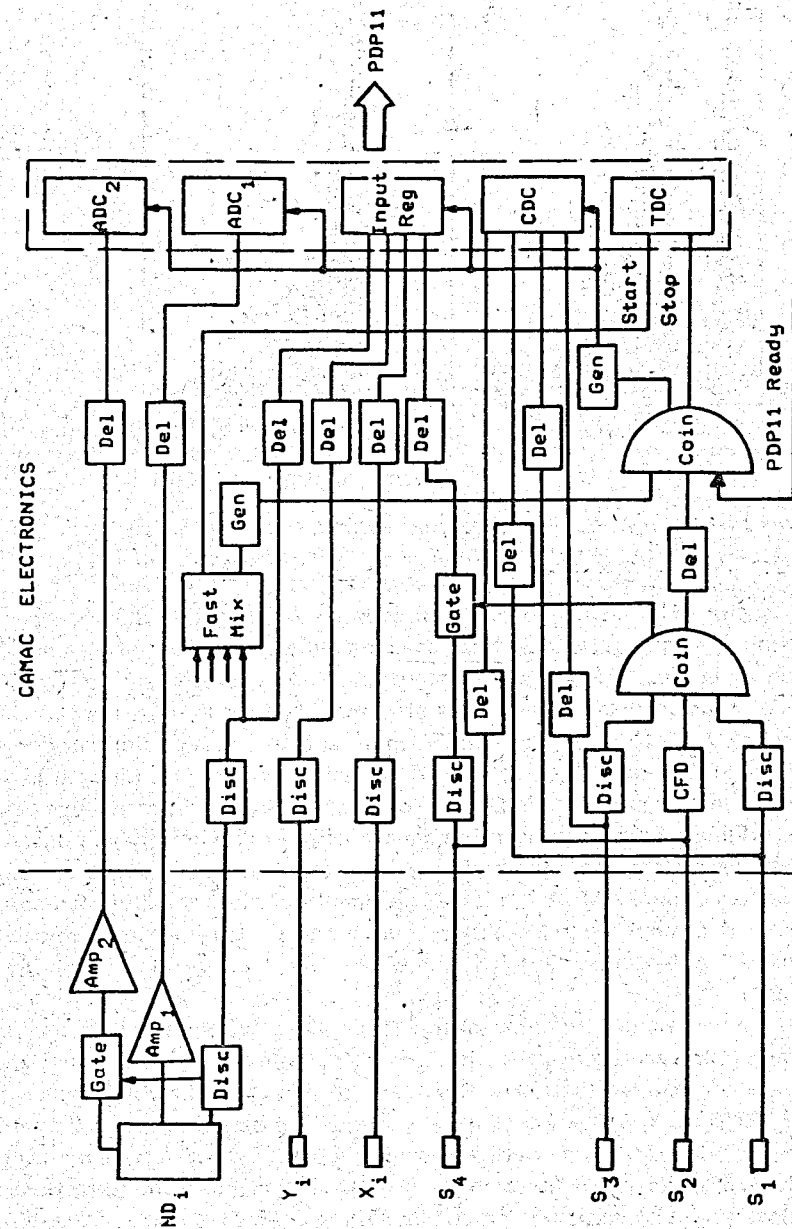


Fig. 2. Block diagram of electronics used in the time-of-flight spectrometer. The abbreviations are as follows: Disc - discriminator, Amp - linear amplifier, Del - delay, CFD - constant fraction discriminator, Coin - coincidence, Gen - generator, ADC - analog-to-digital converter, Input Reg - input register, CDC - charge-to-digital converter, TDC - time-to-digital converter.

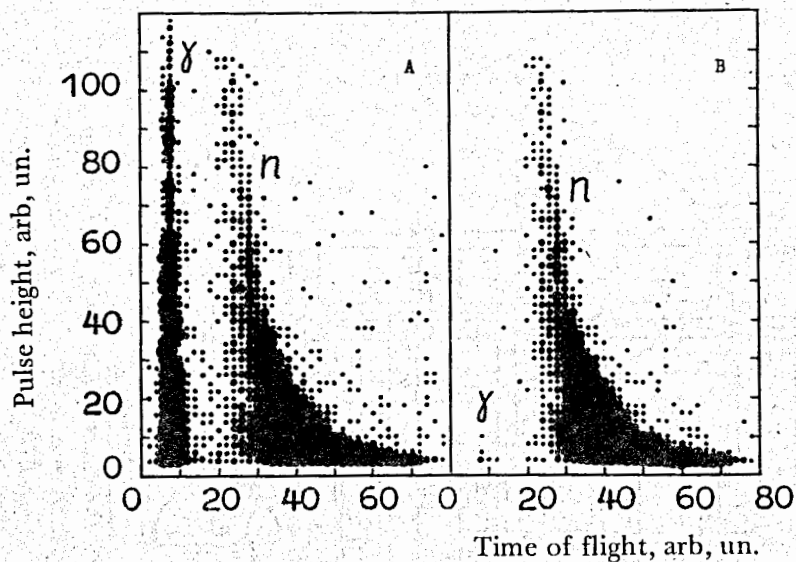


Fig.3. Two time-of-flight spectra measured with the fast ionization chamber with a layer of ^{252}Cf : (A) – without and (B) – with n - γ discrimination.

time-of-flight spectra (without and with n - γ discrimination) measured using the fast ionization chamber with a layer of ^{252}Cf as the source of neutrons and gamma-quanta are shown in fig. 3. Many efforts were undertaken to accurately determine the efficiency of neutron detectors¹⁴. The efficiency data for the detectors based on stylben crystals $\phi 40 \times 10 \text{ mm}^3$ and $\phi 50 \times 50 \text{ mm}^3$ and the plastic scintillator for values of threshold 0.1, 1 and 30 MeV, respectively, are given in fig. 4. Charged particle spectra were measured with X-Y-D detector telescopes. The different types of particles were off-line separated using t-E (time-of-flight – pulse height) plot analysis. In fig. 5 the example of t-E plot is shown for charged particles emitted at 90° from thick lead target by $1 \text{ GeV} \cdot \text{A}$ ^6Li ions. The time resolution was about 0.5 and 1 ns/m (FWHM of prompt gamma-peak) for plastic and stylben detectors, respectively. The typical flight paths were 0.7, 1.3 and 2.0 m for the detectors ND-1, ND-2 and D respectively.

The ion beam had a pulse structure that consisted of macropulse that were about 350 ms long and following through 9 s with about 10^5 ion/macropulse. Within each macropulse were micropulses that were about 20 ns wide and following one after another through 300 ns. The position and the profile of the beam in front of the target was monitored by means of multiwire proportional chamber (MWPC). Beam size was typically 20 mm (FWHM) both in the vertical and horizontal directions.

Data accumulation was accomplished employing the on-line computer. The experimental method and the setup were described in detail in ref.¹⁵. The list of measurements performed for three last years is given in the table.

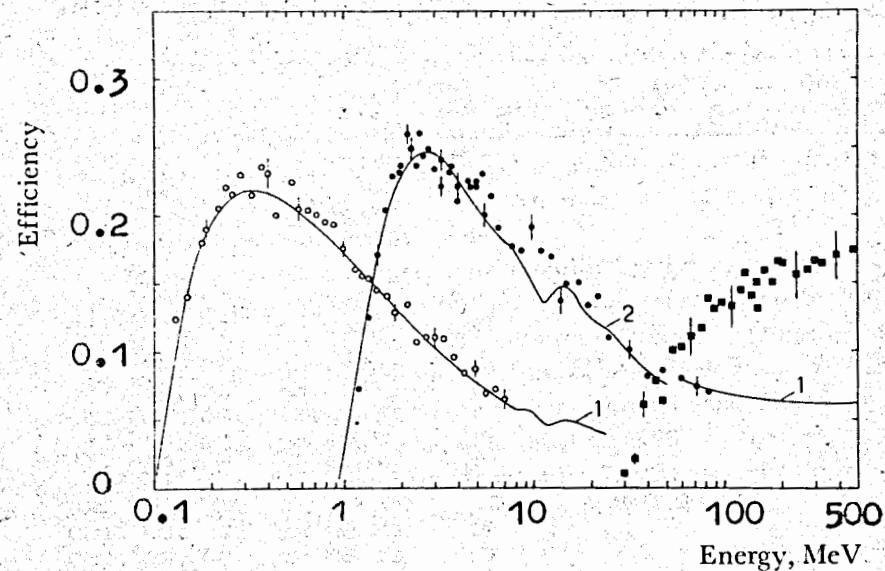


Fig.4. Efficiency for three types of neutron detectors with values of thresholds 0.1 MeV (stylben $\phi 40 \times 10 \text{ mm}^3$), 1 MeV (stylben $\phi 50 \times 50 \text{ mm}^3$) and 30 MeV (plastic scintillator $\phi 120 \times 200 \text{ mm}^3$). The points – experimental results, the curves – results of calculations.

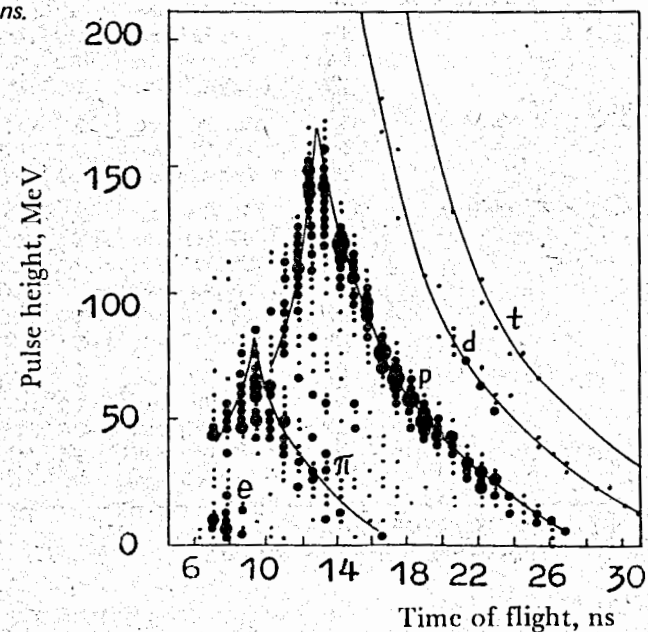


Fig.5. t-E plot for charged particles measured at 90° on $1 \text{ GeV} \cdot \text{A}$ ^6Li ion beam with lead target $8 \times 8 \times 8 \text{ cm}^3$.

Table. Our last time-of-flight measurements (beams, targets, angles)

Number of experiment	Beam ion	Beam energy	Target nucleus*	Target size**	Angle
1	${}^6\text{Li}$	1 GeV A	C	23.3 mm	90°
			C	100 mm	90°
			Pb	8x8x8 cm	90°
			Pb	$\phi 20 \times 20$ cm	90°
2	p	2.55 GeV	Al	$\phi 20 \times 20$ cm	90°
			Pb	8x8x8 cm	90°
			Pb	$\phi 20 \times 20$ cm	$30^\circ, 90^\circ, 150^\circ$
3	p	2.00 GeV	Be	8x8x4 cm	$30^\circ, 90^\circ$
			Al	17.32 mm	$30^\circ, 90^\circ, 150^\circ$
			Cu	11.43 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	6.93 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	8x8x8 cm	$30^\circ, 90^\circ, 150^\circ$
4	${}^{12}\text{C}$	2 GeV A	C	30.60 mm	$30^\circ, 90^\circ, 150^\circ$
			Al	23.09 mm	$30^\circ, 90^\circ, 150^\circ$
			Cu	11.55 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	13.86 mm	$30^\circ, 90^\circ, 150^\circ$
5	d	1 GeV A	Pb	6.93 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	8x8x8 cm	$30^\circ, 90^\circ, 150^\circ$
			Pb	$\phi 20 \times 20$ cm	$30^\circ, 90^\circ, 150^\circ$
6	${}^4\text{He}$	1 GeV A	C	31.18 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	6.93 mm	$30^\circ, 90^\circ, 150^\circ$
			Pb	8x8x8 cm	$30^\circ, 90^\circ, 150^\circ$
			Pb	$\phi 20 \times 20$ cm	$30^\circ, 90^\circ, 150^\circ$
7	${}^{12}\text{C}$	1 GeV A	Cu	15.60 mm	$30^\circ, 60^\circ, 90^\circ$
			Pb	12.00 mm	$30^\circ, 60^\circ, 90^\circ$
			Pb	8x8x8 cm	$30^\circ, 60^\circ, 90^\circ$
7	${}^{12}\text{C}$	2 GeV A	C	24.25 mm	$30^\circ, 60^\circ, 90^\circ$
			Al	17.21 mm	$30^\circ, 60^\circ, 90^\circ$
			Cu	15.60 mm	$30^\circ, 60^\circ, 90^\circ$
			Cd	12.00 mm	$30^\circ, 60^\circ, 90^\circ$
			Pb	12.00 mm	$30^\circ, 60^\circ, 90^\circ$
			Pb	8x8x8 mm	$30^\circ, 60^\circ, 90^\circ$
8	p	2.00 GeV	Al	17.21 mm	$30^\circ, 60^\circ, 90^\circ, 120^\circ$
			Cu	15.60 mm	$30^\circ, 60^\circ, 90^\circ, 120^\circ$
			Cd	12.00 mm	$30^\circ, 60^\circ, 90^\circ, 120^\circ$
			Pb	12.00 mm	$30^\circ, 60^\circ, 90^\circ, 120^\circ$
			Pb	8x8x8 cm	$30^\circ, 60^\circ, 90^\circ, 120^\circ$

* - natural mixture of isotopes

** - effective thickness along beam axis

RESULTS AND DISCUSSION

A transformation of hadron spectra with the variation of angle, target size and atomic mass of incident ion was studied in our experimental data analysis. For the proton beam a calculation of neutron distributions was also carried out by the code SITHA^{16/}. This computer code simulated a secondary hadrons transport using EVAP and D2N2 approximations^{17/} of characteristics of emitted particles in hadron-nucleus inelastic interactions.

The high-energy tail of neutron spectrum and spectra of hydrogen isotopes nuclei at 30° measured on 1 GeV·A ${}^{12}\text{C}$ ion beam with lead target $8 \times 8 \times 12 \text{ cm}^3$ are given in fig. 6. The yield of high-energy neutrons exceeds the proton yield more than by an order and the neutron spectrum overlaps very wide energy region (more than the beam ion energy per nucleon). A comparison of experimental results for high-energy neutron, proton and pion emission at angle 90° from thick lead targets bombarded by 2.55 GeV protons are shown in fig. 7. The yield of charged hadrons falls with an increase of target dimensions but the neutron yield rises.

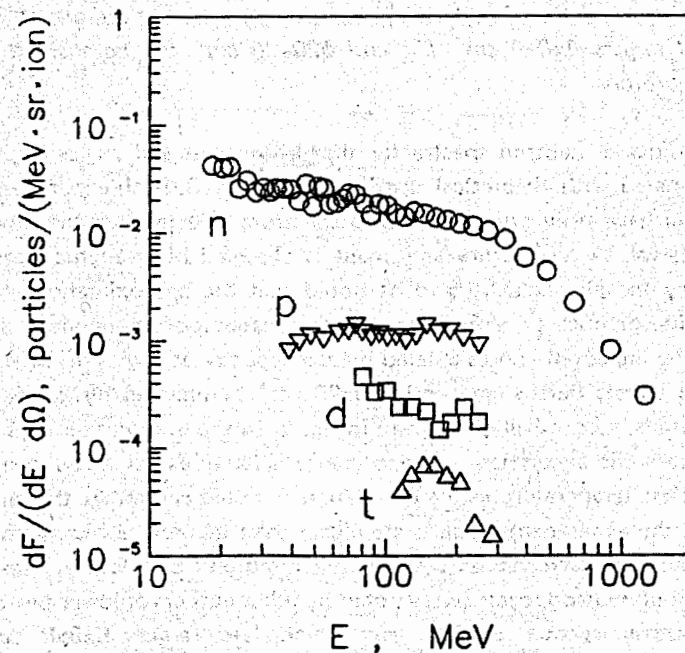


Fig. 6. Energy spectra of n, p, d and t at 30° for lead target $8 \times 8 \times 12 \text{ cm}^3$ bombarded by 1 GeV/N ${}^{12}\text{C}$ ions.

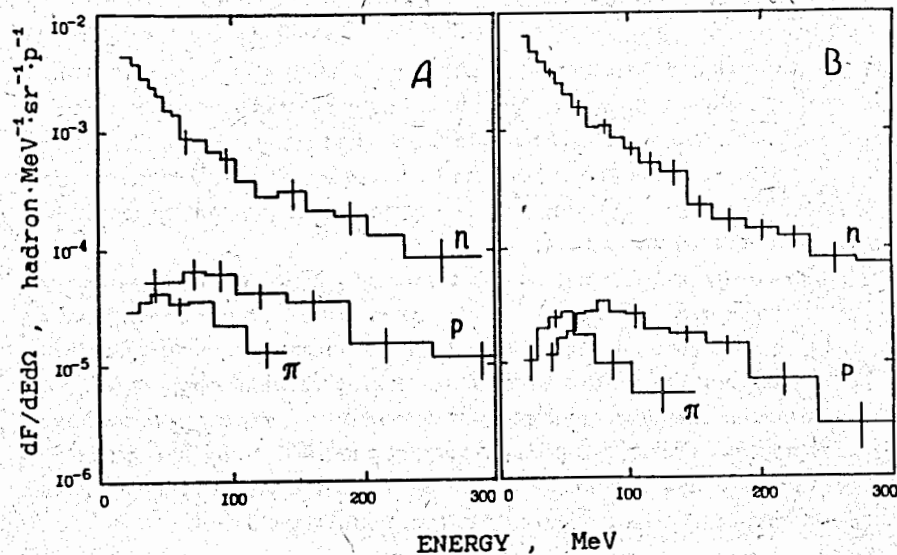


Fig. 7. Energy spectra of high-energy neutrons, protons and pions at 90° from lead targets $8 \times 8 \times 8 \text{ cm}^3$ (A) and $\phi 20 \times 20 \text{ cm}^3$ (B) bombarded by 2.55 GeV protons.

The experimental neutron spectra for thin lead target and angles 30° , 90° , and 150° are compared with theoretical distributions in fig. 8. Both experimental and theoretical data have been normalized on a spectrum integral over the experimental energy interval. A satisfactory agreement is observed between the experiment and the theory for all angles. It's to be noted that the approximations¹⁷ give more high value of nuclear temperature for the evaporation component than the experiment. The measured and calculated neutron spectra at angles of 30° and 150° for thick lead targets $8 \times 8 \times 8 \text{ cm}^3$ and $\phi 20 \times 20 \text{ cm}^3$ bombarded by 2.0 GeV and 2.55 GeV protons, respectively, are given in fig. 9. In the evaporation region a discrepancy between the experiment and the theory is likely explained by using a high value of a nuclear temperature in the calculation as noted earlier. At the same time for all targets the experimental points are higher than theoretical histograms above 20 MeV. Analogous discrepancies between measured and calculated values of the high energy tail of neutron spectra were noted by other authors at lower proton energies. The neutron spectra at 90° from thick lead targets $8 \times 8 \times 8 \text{ cm}^3$ and $\phi 20 \times 20 \text{ cm}^3$ bombarded by different ions are shown in fig. 10. All spectra have a similar shape. For incident protons the results of calculation are shown too. This experimental results are preliminary and at present the data analysis is performed.

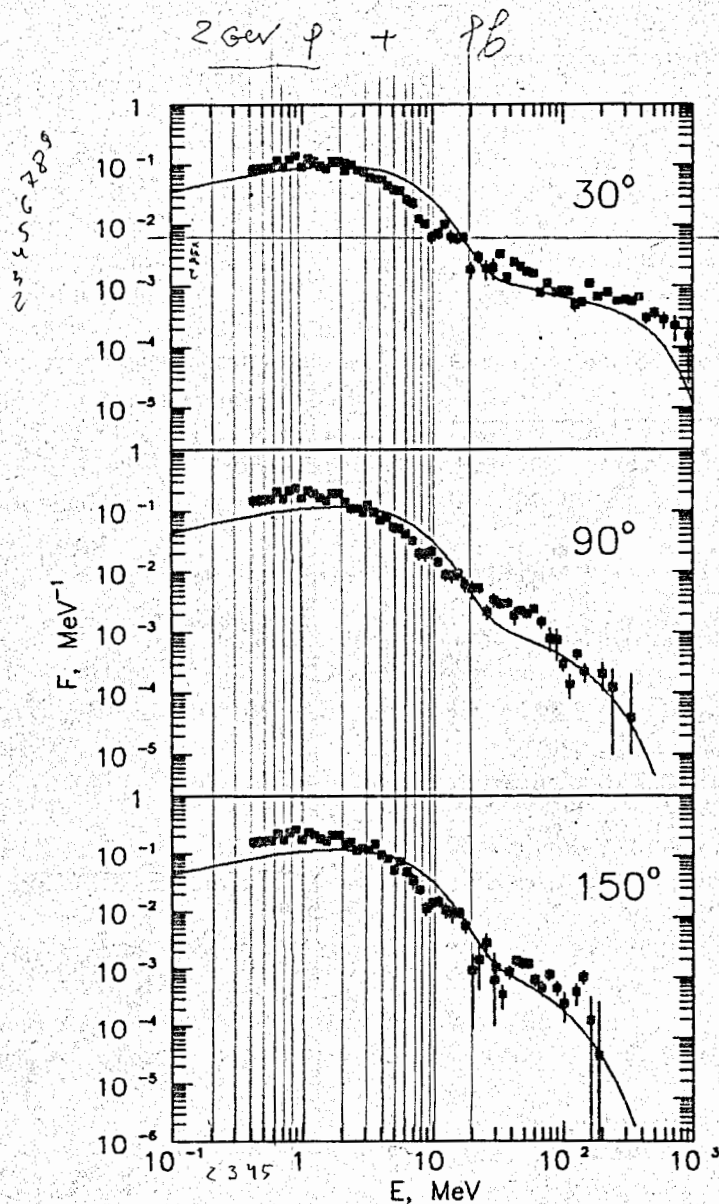


Fig. 8. Normalized neutron energy spectra for thin lead target bombarded by 2 GeV protons: points — the experiment, histograms — the calculation.

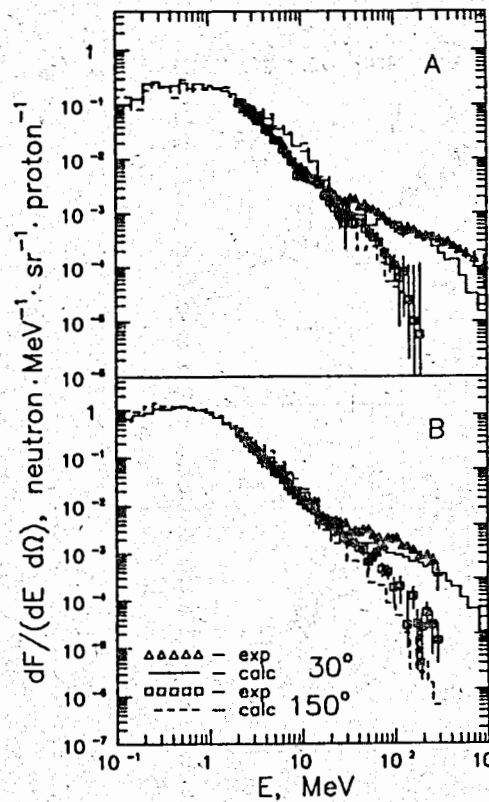


Fig. 9. Neutron spectra at 30° and 150° for lead targets 8x8x8 cm³ (A) and φ20x20 cm³ (B) bombarded by 2 GeV and 2.55 GeV protons respectively: points — the experiment, histograms — the calculation.

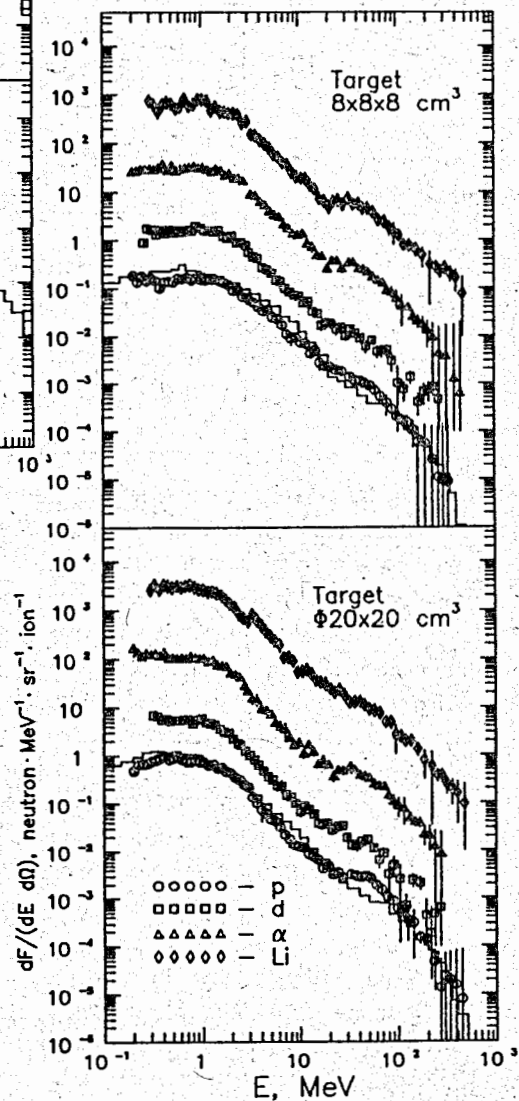


Fig. 10. Neutron spectra at 90° for lead targets 8x8x8 cm³ and φ20x20 cm³ bombarded by protons and 1 GeV-A d, α, and ⁶Li ions: points — the experiment, the histogram — the calculation for a proton beam. The values for d, α and ⁶Li ions have been increased by factors of 10, 100 and 1000, respectively.

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