

ОБЪЕДИНЕННЫЙ  
ИНСТИТУТ  
ЯДЕРНЫХ  
ИССЛЕДОВАНИЙ  
ДУБНА

4639/83

29/VIII-83

E1-83-417

**A STUDY OF THE  $\Lambda_c^+$  PRODUCED  
IN nC INTERACTIONS AT  $\sim 58$  GeV**

**Collaboration BIS-2: Berlin - Budapest -  
Dubna - Moscow - Prague -  
Sofia - Tbilisi**

Submitted to the International European  
Conference on High Energy Physics  
(Brighton, July 20-27 1983)

**1983**

A.N.Aleev, V.A.Arefiev, V.P.Balandin, V.K.Berdyshev, V.D.Cholakov,  
 A.S.Chvyrov, I.I.Evsikov, T.S.Grigalashvili, B.N.Gus'kov,  
 I.M.Ivanchenko, I.N.Kakurin, M.N.Kapishin, N.N.Karpenko,  
 D.A.Kirillov, I.G.Kosarev, V.R.Krastev, N.A.Kus'min, B.A.Kulakov,  
 M.F.Likhachev, A.L.Lyubimov, A.N.Maksimov, P.V.Moisenz,  
 A.N.Morozov, Nguyen Mong Zao, V.V.Pal'chik, L.V.Sil'vestrov,  
 L.A.Slepets, G.G.Sultanov, P.T.Todorov, R.K.Trayanov, N.V.Vlasov  
 Joint Institute for Nuclear Research, Dubna, USSR

K.Hiller, H.Nowak, Z.Nowak, A.V.Pose, H.-E.Ryseck  
 Institute of High Energy Physics of the GDR Academy  
 of Sciences, Berlin-Zeuthen, GDR

A.S.Belousov, A.M.Fomenko, E.I.Malynovsky, S.V.Rusakov,  
 Yu.V.Soloviev, P.N.Shareiko, L.N.Shtarkov, Ya.A.Vazdik  
 Lebedev Physical Institute of the USSR Academy of Sciences,  
 Moscow, USSR

E.A.Chudakov  
 Institute of Nuclear Physics, Moscow State University, USSR

J.Hladky, S.Nemecek, M.Novak, A.Prokes  
 Institute of Physics of the Czechoslovak Academy of Sciences,  
 Prague, Czechoslovakia

M.V.Tosheva  
 Higher Engineering and Electrotechnical Institute, Varna,  
 Bulgaria  
 V.J.Zayachky  
 Higher Chemical-Technological Institute, Sofia, Bulgaria

D.T.Burilkov, V.I.Genchev, I.M.Geshkov, P.K.Markov  
 Institute of Nuclear Research and Nuclear Energetics  
 of the Bulgarian Academy of Sciences, Sofia, Bulgaria

N.S.Amaglobeli, V.P.Dzhordzhadze, V.D.Kekelidze, N.L.Lomidze,  
 R.G.Shanidze  
 Institute of High Energy Physics, Tbilisi State University,  
 USSR

Here we report the result of a search for inclusive  $\Lambda_c^+$  pro-  
 duction by  $\sim 58$  GeV neutrons on carbon at the Serpukhov accele-  
 rator using the spectrometer BIS-2<sup>1/1</sup>. We have observed the  
 following decay modes:

$$\Lambda_c^+ \rightarrow \bar{K}^0 p \pi^+ \pi^-, \quad (1)$$

$$\Lambda_c^+ \rightarrow \Lambda^0 \pi^+ \pi^+ \pi^-. \quad (2)$$

Our result is a first statistically significant direct ob-  
 servation of charmed particle production in strong interactions  
 at these low energies. A previous stage of this investigation  
 has been reported in<sup>2,3/</sup>.

The layout of the BIS-2 spectrometer (fig.1) allows the re-  
 gistration of secondary particles mainly produced in the high  
 energy fragmentation region of the neutron beam. The momentum  
 spectrum of the neutron beam is shown in fig.2. A 7.8 and  
 6. g/cm<sup>2</sup> carbon targets (T) were used. The magnet (M) changed  
 the transverse momentum of charged particles by 0.64 GeV/c. To  
 trigger the spectrometer, four or more charged particles were  
 required to pass the whole spectrometer. The spectrometer ope-  
 rated on-line with an EC-1040 computer.

The results are based on the data obtained at three diffe-  
 rent configurations of the spectrometer. Main differences were  
 the polarity of the magnetic field in the magnet, positions of  
 the target, and some geometrical parameters of the proportio-  
 nal chambers (PC). All the changes of the experimental condi-  
 tions allowed us to estimate possible systematic errors.  
 5.3·10<sup>6</sup>, 1.7·10<sup>6</sup>, and 4.4·10<sup>6</sup> events were recorded at each  
 of the configurations, respectively. These events corresponded  
 to the integral neutron flux  $M_n \approx 6 \cdot 10^{11}$  through the target.

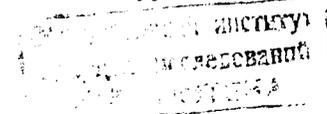
We searched for inclusive  $\Lambda_c^+$  production in the process

$$n + C \rightarrow \Lambda_c^+ + \dots \quad (3)$$

$$\downarrow$$

$$\Lambda^0 (\bar{K}^0) h^+ h^+ h^-$$

This process possesses the topology with a  $\Lambda^0 (\bar{K}^0)$  "Vee", having  
 its vertex outside the target, and with three charged particles  
 ( $h^+$ ,  $h^+$  and  $h^-$ ) emitted directly from it. The "Vee" vertex was  
 required to have the closest approach less than 1 cm and to lie



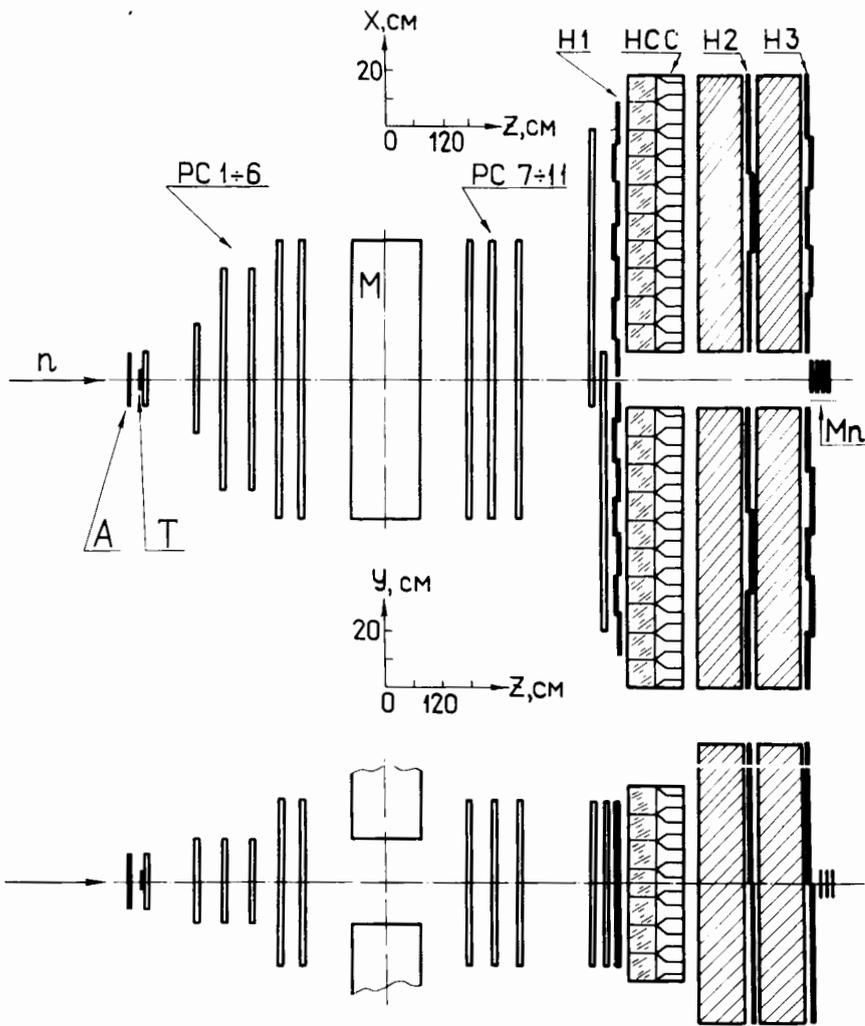


Fig.1. A scheme of the BIS-2 set up: M - magnet; PC - two-coordinate multiwire proportional chambers; A - scintillation anticoincidence counter; T - carbon target; H1 - scintillation hodoscope;  $M_n$  - neutron beam monitor; H2 and H3 -  $\mu$  detector hodoscopes; HCC - lead glass walls. No information from H2, H3 and HCC have been used in data analysis presented.

more than 15 cm downstream the centre of the target. Figure 3 illustrates the "Vee" invariant mass spectra for  $K_S^0$  and  $\Lambda^0$ . We identified the "Vee", having the invariant mass of a two-

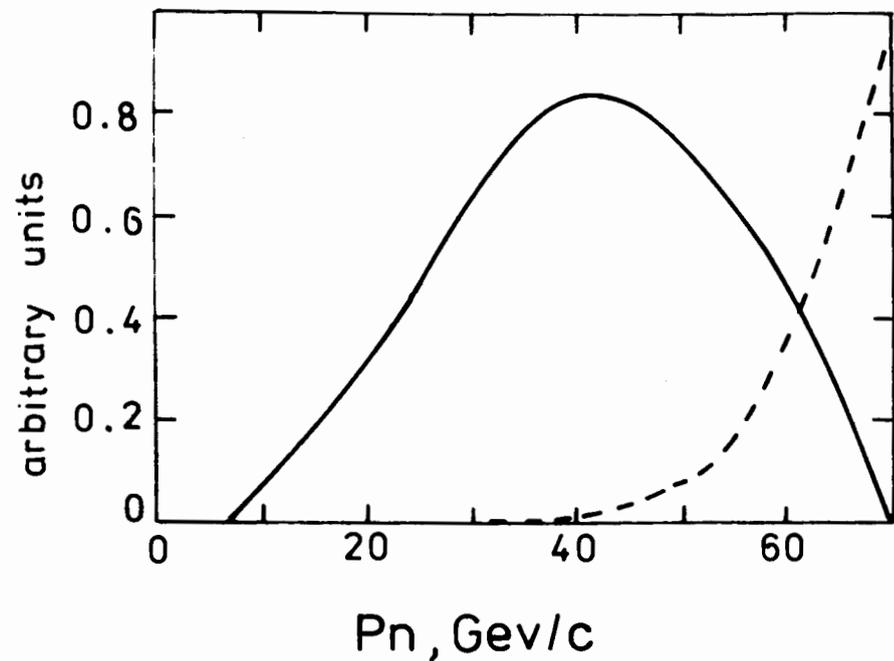


Fig.2. Neutron beam momentum spectrum (solid curve) and neutron momentum acceptance for  $\Lambda_c^+$  (dashed curve).

pion system  $M(\pi^+\pi^-)$  with  $+10$  MeV/ $c^2$  from  $M(K_S^0) = 497.67$  MeV/ $c^2$  with  $K_S^0$ , and the "Vee" having the invariant mass of a system of a proton and negative pion  $M(p\pi^-)$  with  $+7$  MeV/ $c^2$  from  $M(\Lambda^0) = 1115.6$  MeV/ $c^2$  with  $\Lambda^0$ .

We have selected events containing at least one  $K_S^0 h^+ h^-$  or  $\Lambda^0 h^+ h^-$  combination. The  $M(K_S^0 p \pi^+ \pi^-)$  invariant mass distribution for  $K_S^0 h^+ h^-$  combinations, assuming that the largest momentum of the positive particles belongs to the proton, is presented in fig.4a (solid line distribution). The width of the bins ( $= 20$  MeV/ $c^2$ ) was chosen equal to  $\sim 2$  times the experimental resolution. Figure 4b (solid line distribution) presents the invariant mass  $M(\Lambda^0 \pi^+ \pi^+ \pi^-)$  distribution for  $\Lambda^0 h^+ h^-$  combinations. The width of the bins ( $= 30$  MeV/ $c^2$ ) is  $\sim 4$  times the experimental resolution for this invariant mass. Explicit peaks are seen in both distributions in the mass region of the charmed baryon  $\Lambda_c^+$ . To estimate the number of events in the peaks a polynomial background, was fitted to these distributions (dotted distributions in fig.4).  $130 \pm 18$  and  $57 \pm 14$  events above the background correspond to  $\sim 10$  and  $\sim 5$  standard deviations. We consider these peaks as a baryon state with the width less than  $30$  MeV/ $c^2$ . If the observed baryon is a hyperon resonance,

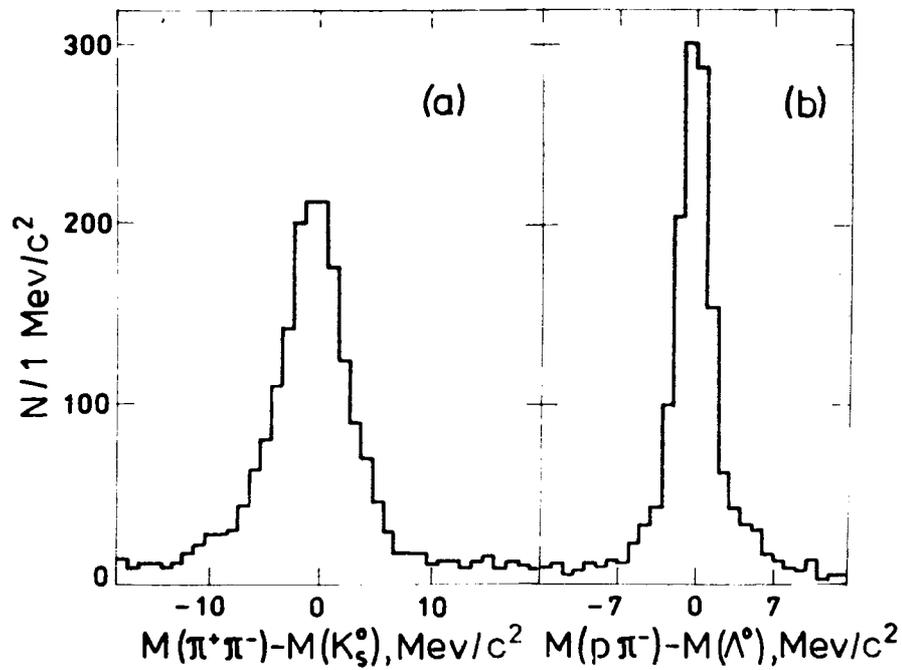


Fig.3. "Vees" invariant mass spectra illustrating  $K_S^0$  and  $\Lambda^0$  detection at the part of the statistics. (a) is the  $M(\pi^+\pi^-)$  spectrum in a mass region around  $M(K_S^0) = 497.67 \text{ MeV}/c^2$ ; (b) is the  $M(p\pi^-)$  spectrum in a mass region around  $M(\Lambda^0) = 1115.6 \text{ MeV}/c^2$ .

in the neutron beam fragmentation region one could observe, with a higher probability, the isotopically conjugated negative state (as, for example, the production of the  $\Sigma(1385)^{4/}$ ). But no significant peaks are seen in the invariant mass distributions for  $K_S^0 p \pi^+ \pi^-$  and  $\Lambda^0 \pi^+ \pi^- \pi^-$  systems selected similarly to  $K_S^0 p \pi^+ \pi^-$  and  $\Lambda^0 \pi^+ \pi^+ \pi^-$  (the crossline distributions in fig.4). So, we are led to identify the observed peaks with the Cabibbo favoured charmed baryon  $\Lambda_c^+$  decays (1) and (2). The latter decay mode has been detected elsewhere<sup>5/</sup>.

The partial cross section is given by  $\sigma B = NA / (N_A T M_n B_0 \epsilon)$ , where  $N$  is the number of  $\Lambda_c^+$  decays observed;  $A$  is the atomic number of the carbon nucleus;  $N_A$  is the Avogadro number;  $T$  is the target length;  $M_n$  is the neutron flux;  $B_0$  is the decay ratio of  $K^0(\Lambda^0)$  observed;  $B$  is the decay ratio of  $\Lambda_c^+$  observed;  $\epsilon$  is the detection efficiency of  $\Lambda_c^+$  obtained by the Monte-Carlo simulation taking into account all experimental and data processing conditions. The  $\Lambda_c^+$  decays (1) and (2) simulated according to the phase space of four particles. The dashed line

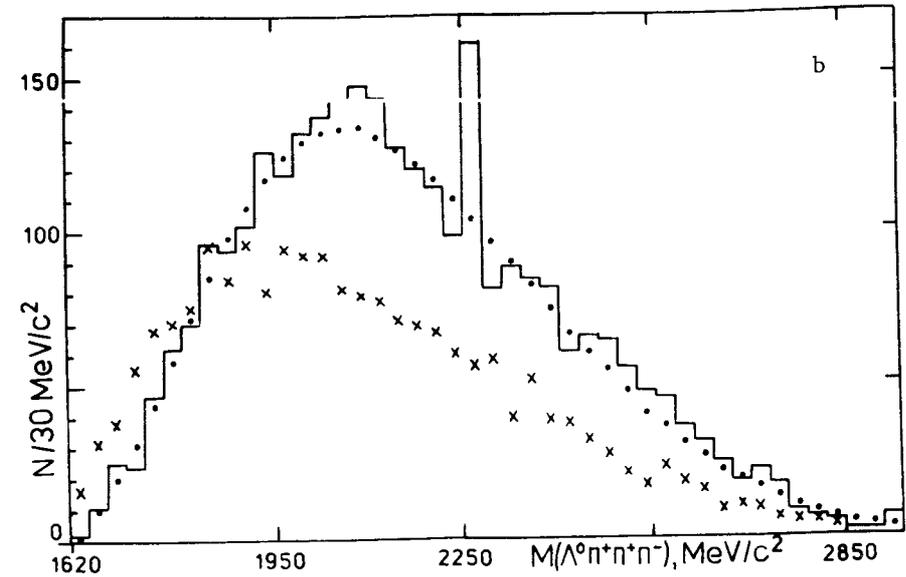
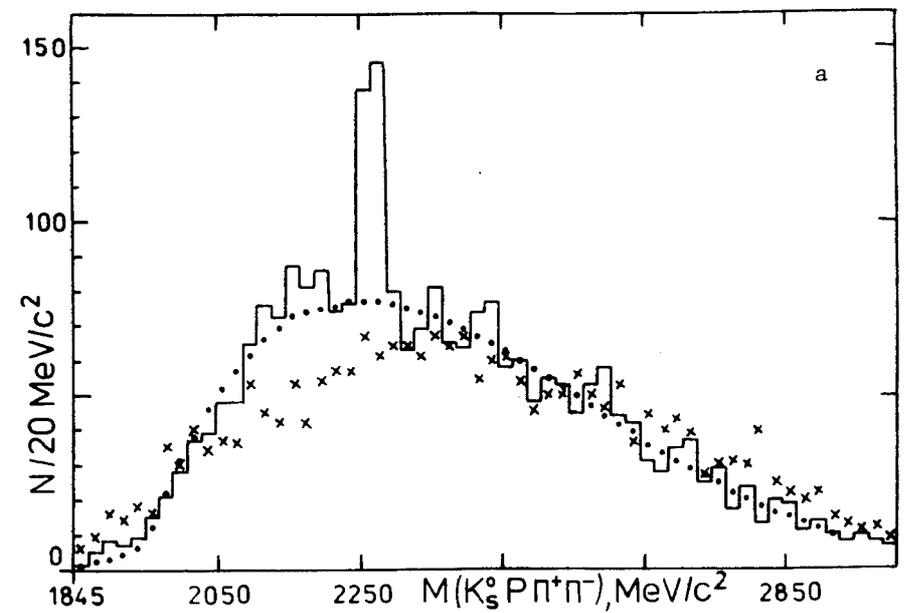


Fig.4. The invariant mass distributions: (a) for  $K_S^0 p \pi^+ \pi^-$  (solid line) and  $K_S^0 p \pi^- \pi^-$  (cross-line) systems; (b) for  $\Lambda^0 \pi^+ \pi^+ \pi^-$  (solid line) and  $\Lambda^0 \pi^+ \pi^- \pi^-$  (cross-line) systems. The dotted lines show the result of the polynomial fit to these distributions.

in fig.2 represents the beam momentum acceptance for  $\Lambda_c^+$ . The momentum of neutrons having produced  $\Lambda_c^+$  registered is greater than  $\sim 40$  GeV/c (the mean value is about 58 GeV/c). The  $\Lambda_c^+$  observed have the following transverse and longitudinal momentum ranges:  $P_{\perp} \leq 1$  GeV/c and Feynman x-variable greater than  $\sim 0.5$ .

We have obtained the partial cross section of  $\Lambda_c^+$  production for  $X_F > 0.5$  times the branching ratios of the decay (1)  $B_1$

$$\sigma(X_F > 0.5)B_1 = (6.2 \pm 1.3) \mu\text{b} \quad \text{per carbon nucleus}$$

and of the decay (2)  $B_2$

$$\sigma(X_F > 0.5)B_2 = (2.0 \pm 0.6) \mu\text{b} \quad \text{per carbon nucleus.}$$

To estimate the partial cross section in the complete region ( $X_F > 0$ ) we apply a variety of models for  $\Lambda_c^+$  production in neutron-nucleon interactions.

The A model imply the invariant inclusive cross section in the form:

$$Ed^3\sigma/d\vec{p}^3 \sim \exp(-b \cdot P_{\perp}) \cdot f(X_F), \quad (4)$$

where  $f(X_F)$  is the  $X_F$  distribution function and  $b = 2.5(\text{GeV}/c)^{-1/3}$ . The latter value is also in accordance with the experiment<sup>6,2'</sup>.  $f(X_F)$  has the same shape as for inclusive production of  $\Lambda^0$  hyperons by 69 GeV protons in pp interactions<sup>7'</sup>. This should be expected from the quark model for hadron fragmentation<sup>8'</sup> and has been also confirmed by an ISR experiment<sup>9'</sup>.

The B model: the system R is produced by diffraction dissociation of neutrons on quasifree nucleons in the carbon nucleus according to the invariant cross section  $Ed^3\sigma/d\vec{p}^3 \sim 1/M_R^2 \cdot \exp(-\beta \cdot t)$ , where  $M_R$  is the invariant mass of the R system,  $\beta = 6(\text{GeV}/c)^{-2}$ . The R system decays into  $\Lambda_c^+ D^-$  with Gottfried-Jackson angle distributions as calculated in<sup>10'</sup> for the diffraction dissociation  $p \rightarrow \Lambda^0 K^+$ .

Results close to the model B are obtained for  $\Lambda_c^+$  inclusive production using the cross section (4) with a distribution for  $f(X_F)$  calculated in<sup>11-13'</sup>. If using the B model we assume a flat distribution of the Gottfried-Jackson angles, the value of  $\sigma \cdot B$  is close to the same value in the A model.

The invariant production cross section for  $X_F > 0.5$  is described by

$$Ed^3\sigma/d\vec{p}^3 \sim (1 - X_F)^{(0.9 \pm 0.5)} \cdot \exp[-(2.5 \pm 0.6) \cdot P_{\perp}] .$$

In conclusion we summarize the results obtained. We have found  $\sim 190 \Lambda_c^+$  baryons. Their average mass is  $(2266 \pm 4) \text{ MeV}/c^2$ . Our mass value is in good agreement with the values obtained in other experiments<sup>14'</sup>.

The ratio of the branching ratios  $B_1(\bar{K}^0 p \pi^+ \pi^-)/B_2(\Lambda^0 \pi^+ \pi^+ \pi^-)$  is  $3.1 \pm 1.0$ .

The partial cross sections  $\sigma \cdot B$  for  $\Lambda_c^+$  production by  $\sim 58$  GeV neutrons on carbon for  $X_F > 0.5$  and calculated for  $X_F > 0$  within the framework of the models A and B are presented in the table. The possible systematic errors of  $\sigma \cdot B$  are mainly due to the error of the momentum spectrum measurements. We estimate that this error does not exceed 35%. We could not obtain values of  $\sigma \cdot B$  lower than the results of the B-model taking into account all considered models. Using the result of the B-model and the upper limit for the branching ratio of  $B_2(\Lambda^0 \pi^+ \pi^+ \pi^-)$  3.1%<sup>14'</sup> (90% confidence level), we estimate the inclusive production cross section  $\sigma \geq (70 \pm 20) \mu\text{b}$  per carbon nucleus for  $X_F > 0$ .

Table

Partial cross sections of  $\Lambda_c^+$  production times branching ratios  $\sigma \cdot B$  (per carbon nucleus)

Decay	$X_F > 0.5$	$\sigma \cdot B, \mu\text{b}$ per carbon nucleus	
		$X_F > 0$	
		model A	model B
$\bar{K}^0 p \pi^+ \pi^-$	$6.2 \pm 1.3$	$18.3 \pm 3.7$	$6.7 \pm 1.4$
$\Lambda^0 \pi^+ \pi^+ \pi^-$	$2.0 \pm 0.6$	$6.0 \pm 2.0$	$2.2 \pm 0.7$

The authors are grateful to A.M.Baldin, P.A.Cherenkov, N.N.Govorun, Kh.Ya.Khristov, K.Lanius, A.A.Logunov, M.G.Meshcheryakov, I.A.Savin, L.D.Soloviev, A.N.Tavkhelidze, I.S.Zlatev, I.F.Kolpakov, and E.I.Maltsev for their support of these studies, to V.Barger, S.S.Gershtein, A.B.Kaidalov, A.K.Likhoded for useful discussions and remarks, to the staff of the Serpukhov accelerator, to the staff of SSED JINR for providing the experiment during the preparation of BIS-2 and runs at the accelerator, to all specialists at different stages of its preparation and running, to E.M.Likhacheva for her constant participation in the experiment.

REFERENCES

1. Eichner G. et al. JINR, 1-80-644, Dubna, 1980; Aleev A.N. et al. JINR, 1-81-67, Dubna, 1981; Maksimov A.N. JINR, 1-81-574, Dubna, 1981; Burilkov D.T. et al. JINR, 10-80-656, Dubna, 1980; Burilkov D.T. et al. JINR, 10-81-772, Dubna, 1981.
2. Aleev A.N. et al. Yad.Fiz., 1982, 35, p.1175; Sov.J.Nucl. Phys., 1982, 35, p.687; Yad.Fiz., 1983, 37, p.1474; JINR, P1-81-693, Dubna, 1981; JINR, D1-82-895, Dubna, 1982.
3. Aleev A.N. et al. JINR, P1-82-343, E1-82-759, Dubna, 1982.
4. Aleev A.N. et al. JINR, P1-82-353, Dubna, 1982; Aleev A.N. et al. PHE 82-7, Berlin-Zeuthen, 1982.
5. Gazzoli E.G. et al. Phys.Rev.Lett., 1975, 344, p.1125; Knapp B. et al. Phys.Rev.Lett., 1976, 37, p.882; Baltay C. et al. Phys.Rev.Lett., 1979, 42, p.1721; Giboni K.L. et al. Phys.Lett., 1979, 85B, p.437; Lockman W. et al. Phys.Lett., 1979, 85B, p.443.
6. Basile M. et al. Nuovo Cim.Lett., 1981, 30, p.481.
7. Ammosov V.V. et al. Nucl.Phys., 1976, B115, p.269.
8. Anisovich V.V., Shekhter V.M. Nucl.Phys., 1973, B55, p.455.
9. Basile M. et al. Nuovo Cim.Lett., 1981, 30, p.487; Irion J. et al. Phys.Lett., 1981, 99B, p.495.
10. Bohringer T. CERN-EP 77-18, Geneva, 1977.
11. Brodsky S.J. et al. Phys.Lett., 1980, 93B, p.451; Brodsky S.J. et al. Phys.Rev., 1981, D23, p.2745; Bertsch G. et al. Phys.Rev.Lett., 1981, 47, p.297.
12. Mazzanti P., Wada S. Phys.Rev., 1982, D26, p.602.
13. Boreskov K.G., Kaidalov A.B. Preprint ITEP, 130, Moscow, 1982.
14. Roos M. et al. Phys.Lett., 1982, 111B, p.1.

Алеев А.Н. и др.

E1-83-417

Изучение  $\Lambda_c^+$ , рожденных в nC взаимодействиях при ~58 ГэВ

Наблюдается рождение очарованного бариона  $\Lambda_c^+$  в процессе  $n + C \rightarrow \Lambda_c^+ + \dots$ . Эксперимент выполнен в нейтронном пучке серпуховского ускорителя с помощью спектрометра БИС-2. В распределениях по эффективным массам систем  $K_S^0 p \pi^+ \pi^-$  и  $\Lambda^0 \pi^+ \pi^+ \pi^-$  около значения массы 2270 МэВ/с<sup>2</sup> наблюдаются узкие пики 10 и 5 стандартных отклонений от уровня фона соответственно. Найдено, что масса  $M(\Lambda_c^+) = 2266 \pm 4$  МэВ/с<sup>2</sup>. Произведения сечения рождения на вероятность распада по наблюдаемым каналам, определенные для  $X_F > 0,5$ , равны  $6,2 \pm 1,3$  и  $2,0 \pm 0,6$  мб на ядро углерода соответственно, а их отношение равно  $3,1 \pm 1,0$ . Инвариантное сечение рождения  $\Lambda_c^+$  описывается выражением:

$$E \frac{d^2\sigma}{P_{\perp} dp_{\perp} dp_{\parallel}} \sim (1 - X_F)^{(0.9 \pm 0.5)} \exp[(-2.5 \pm 0.6) P_{\perp}]$$

Работа выполнена в Лаборатории высоких энергий ОИЯИ, Preprint Объединенного института ядерных исследований, Дубна 1983

Aleev A.N. et al.

E1-83-417

A Study of the  $\Lambda_c^+$  Produced in nC Interactions at ~58 GeV

The production of the charm baryon  $\Lambda_c^+$  has been observed in the reaction  $n + C \rightarrow \Lambda_c^+ + \dots$ . The experiment has been performed in a neutron beam of the Serpukhov accelerator, using the spectrometer BIS-2. The effective mass distributions of the  $(K_S^0 p \pi^+ \pi^-)$  and  $(\Lambda^0 \pi^+ \pi^+ \pi^-)$  systems show a narrow peak of ~10 and ~5 standard deviations above the background level, respectively, at an average mass value  $M(\Lambda_c^+) = 2266 \pm 4$  MeV/c<sup>2</sup>. The cross sections times the branching ratios measured for  $X_F > 0.5$ , for each  $\Lambda_c^+$  decay channel, are equal to  $(6.2 \pm 1.3)$  and  $(2.0 \pm 0.6)$   $\mu$ b per carbon nucleus, respectively, and their ratio is  $3.1 \pm 1.0$ . The invariant production cross section for  $X_F > 0.5$  is described by:

$$E \frac{d^3\sigma}{d\vec{p}^3} \sim (1 - X_F)^{(0.9 \pm 0.5)} \cdot \exp[-(2.5 \pm 0.6) P_{\perp}]$$

The investigation has been performed at the Laboratory of High Energies, JINR.

Preprint of the Joint Institute for Nuclear Research, Dubna 1983