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ЛАБОРАТОРИЯ ТЕОРЕТИЧЕСКОЙ ФИЗИКИ

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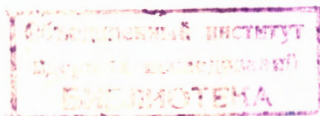
THE ORBITAL FRACTIONAL PARENTAGE
COEFFICIENTS FOR TWO AND THREE
NUCLEONS IN THE 2 SLD SHELL IN THE
SU(3) CLASSIFICATION SCHEME

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It is well-known^{1/}, that it is possible to introduce a basis in the space of orbital wave functions for k nucleons in the 2 sld shell of the harmonic oscillator, so that under transformations belonging to $U(6) \supset SU(3) \supset SO(3)$ the wave functions of this basis

$$\phi_{1\dots k} ([f] \beta (\lambda) \alpha LM(r)) \quad (1)$$

transform according to the irreducible representations^{x)} $[f], (\lambda)$ and L of these groups. The transformation properties of these functions with respect to nucleon permutations are determined by the Yamanouchi symbol r . Here α and β are the so-called additional quantum numbers.

The wave functions (1) may be expressed in terms of those for k-1 nucleons by means of the orbital fractional parentage coefficients

$$\begin{aligned} & \phi_{1\dots k} ([f] \beta (\lambda) \alpha LM | r) = \\ & = \sum_{\substack{(\bar{f} \bar{\beta} (\bar{\lambda}) \bar{\alpha} \bar{L}; \bar{\ell}) \\ (\bar{\beta} \bar{\lambda}) \bar{\alpha} \bar{\ell}}} [f] \beta (\lambda) \alpha L [\phi_{1\dots k-1} ([\bar{f}] \bar{\beta} (\bar{\lambda}) \bar{\alpha} \bar{L} (r) \phi_k (\bar{\ell})]_M^L \end{aligned} \quad (2)$$

These coefficients are the products of two factors, each of which depends only on a smaller number of variables and has a simple group-theoretical meaning :

$$\begin{aligned} & ([\bar{f}] \bar{\beta} (\bar{\lambda}) \bar{\alpha} \bar{L}; \bar{\ell}) [f] \beta (\lambda) \alpha L = \\ & = ([f] \beta (\lambda); [1] (20) [[\bar{f}] \bar{\beta} (\bar{\lambda})] (\bar{\alpha} \bar{L}; (20) \bar{\ell} | (\lambda) \alpha L). \end{aligned} \quad (3)$$

^{x)} We denote an irreducible representation of $SU(3)$ $(\lambda) = (\lambda_1 \lambda_2)$ instead of Elliott's notation $(\lambda \mu)$. As far as $SO(3)$ is concerned, we follow Fano and Racah^{2/}, except of some unimportant changes in notation.

Table 1

k=3 orbital fractional parentage coefficients

		(21)	(11)				(21)	(11)									
		[F]	[2]				[41]	[2]									
		(λ)	(21)			(40)	(02)	(21)	(40)			(02)					
		L	F	D	P	G	D	S	F	D	P	G	D	S	D	S	
(321) [111]	(30)	F	$2/\sqrt{3.5}$	$2/3$	$\sqrt{4/3.5}$				$-1/3\sqrt{2}$								
	(03)	P	$-2/\sqrt{3.5}$	$-1/3\sqrt{2}$	$2/\sqrt{3.5}$				$1/3$								
(211) (121) [21]	(41)	H	1			1											
		G	$\sqrt{4/2.3}$	$\sqrt{5/2.3}$		$-1/3\sqrt{11/2.7}$	$-\sqrt{5/2.7}$		$\sqrt{1/3}$			$1/3\sqrt{5}$					
		F	$-\sqrt{4/2.5}$	$\sqrt{4/2.3}$	$\sqrt{5/3}$	$-\sqrt{5/2.7}$	$3\sqrt{4/2.7}$								$\sqrt{4/2.3}$		
		D	$-2/3\sqrt{5}$	$-1/3\sqrt{2}$	$\sqrt{7/3.5}$	$2/\sqrt{2.7}$	$\sqrt{4/2.7}$	$-\sqrt{7/3.5}$		$1/3\sqrt{2}$							
	(22)	G	$\sqrt{5/2.3}$	$-\sqrt{2/2.3}$		$-\sqrt{11/2.7}$	$\sqrt{1/2.5.7}$		$\sqrt{1/3}$				$\sqrt{5/4}$				
		F	$\sqrt{4/2.5}$	$\sqrt{4/2.3}$	$-\sqrt{2/3}$	$-3\sqrt{4/2.7}$	$-\sqrt{5/2.7}$										
		D'	$-\sqrt{7/2.3.5}$	$1/2\sqrt{7/3}$	$\sqrt{4/2.3}$	$-2/3\sqrt{1/2}$	$2/3\sqrt{1/5}$	$1/5\sqrt{1/2}$	$-2/3$	$-1/2\sqrt{1/3}$			$\sqrt{5/4}$	$1/2\sqrt{7/3}$		$-2/3\sqrt{1/2}$	$\sqrt{7/2.3.5}$
	(11)	D''	$-2/3\sqrt{1/2}$	$1/2\sqrt{1/5}$	$-1/3\sqrt{3/2}$	$3\sqrt{3/2.5.7}$	$5/2\sqrt{1/3.7}$	$\sqrt{7/2.3.5}$		$1/2\sqrt{7/3}$				$-1/2\sqrt{1/3}$			
		S		-1		$-\sqrt{7/3.5}$								$2/\sqrt{3.5}$			$-\sqrt{5/2.3}$
		D	$2/\sqrt{3.5}$	$1/3\sqrt{1/3}$	$1/3\sqrt{1/3}$	$\sqrt{7/3.5}$											
(00)	P	$-2/3\sqrt{3/5}$	$-1/3\sqrt{2}$	$-1/3\sqrt{3/3}$												$-2/\sqrt{3.5}$	
(411) [3]	(60)	I				1											
		G				$-2/3\sqrt{1/2}$	$\sqrt{2.11/2.7}$										
		D				$2/5\sqrt{1/7}$	$-2\sqrt{2/2.7}$	$1/3\sqrt{7}$									
	(22)	S				$2/\sqrt{2.3.5}$											
		G				$\sqrt{11/2.7}$	$\sqrt{1/2.5.7}$										
		F				$-3\sqrt{4/2.7}$	$-\sqrt{5/2.7}$										
		D'				$-2/3\sqrt{1/2}$	$2/2\sqrt{1/5}$	$1/3\sqrt{1/2}$	$\sqrt{5/3}$	$-1/2$							
	(00)	D''				$3\sqrt{3/2.5.7}$	$5/2\sqrt{1/3.7}$	$\sqrt{7/2.3.5}$									
		S				$-\sqrt{7/3.5}$											
		S				$\sqrt{5/2.3}$											$\sqrt{1/2.3}$

I = 2

I = 0

The orbital coefficients of fractional parentage for $k=3$ and $k=2$ were calculated by the usual methods (see, e.g. ^{3/}, and the references quoted there) and are presented in tables 1 and 2. The arrangement of our tables is similar to that of Jahn ^{4/}. The coefficients with $\bar{l}=2$ are written on the left, the coefficients with $\bar{l}=0$ on the right-hand side.

Besides the three components of angular momentum, we use five operators $Q(\mu) = \sum_{\mu=1}^5 Q_{\mu}(\mu)$ where

$$Q_{\mu}(\mu) = -[a_{\mu}^{\dagger}(1)a_{\mu}(1)]_{\mu}^2 \quad (4)$$

as the infinitesimal operators of the algebra of $SU(3)$. The used standard form of the irreducible representations of $SU(3)$ is specified completely by the double bar matrix $(\alpha L \| Q \| \alpha' L')$ in table 3.

Further, the matrix elements of the Wigner central force, between our two-particle states, expressed in terms of the Talmi integrals are presented (table 4).

References

1. J.P.Elliott, Proc. Roy. Soc. A 245, 128 (1958).
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Table 2

$k=2$ orbital fractional parentage coefficients

	\bar{F} [\bar{F}] ($\bar{\lambda}$)	\bar{L}	(1) [1] (20)	(1) [1] (20)
			D S	D S
r [f] (λ)		L		
(21) [11] (21)		F	$\begin{bmatrix} 1 & \\ & -\sqrt{\frac{1}{2}} \end{bmatrix}$	$\begin{bmatrix} \sqrt{\frac{1}{2}} \\ \\ \end{bmatrix}$
		D		
		P	$\begin{bmatrix} -1 & \\ & \end{bmatrix}$	
		G	$\begin{bmatrix} 1 & \\ -\frac{1}{3}\sqrt{2} & \frac{1}{3}\sqrt{\frac{7}{2}} \\ \frac{2}{3} & \end{bmatrix}$	$\begin{bmatrix} \frac{1}{3}\sqrt{\frac{7}{2}} \\ \frac{1}{3}\sqrt{5} \\ \end{bmatrix}$
(11) [2]		D		
		S	$\begin{bmatrix} \frac{1}{3}\sqrt{7} & \frac{1}{3} \\ \frac{1}{3}\sqrt{5} & \end{bmatrix}$	$\begin{bmatrix} \frac{1}{3} \\ -\frac{2}{3} \\ \end{bmatrix}$
			$\bar{l}=2$	$\bar{l}=0$

Table 3

Standard form of the double bar matrix ($aL||Q||a'L$)

(02) $\begin{matrix} D & S \\ D & \begin{matrix} \sqrt{\frac{5}{3}} & -2\sqrt{\frac{5}{3}} \\ -2\sqrt{\frac{5}{3}} \end{matrix} \\ S & \end{matrix}$

(03) $\begin{matrix} F & P \\ F & \begin{matrix} 3\sqrt{\frac{2}{5}} & -2\sqrt{\frac{3}{5}} \\ -2\sqrt{\frac{3}{5}} & 3^2\sqrt{\frac{1}{5}} \end{matrix} \\ P & \end{matrix}$

(11) $\begin{matrix} D & P \\ D & \begin{matrix} & \sqrt{3 \cdot 5} \\ -\sqrt{3 \cdot 5} \end{matrix} \\ P & \end{matrix}$

(20) $\begin{matrix} D & S \\ D & \begin{matrix} -\sqrt{\frac{5}{3}} & 2\sqrt{\frac{5}{3}} \\ 2\sqrt{\frac{5}{3}} \end{matrix} \\ S & \end{matrix}$

(30) $\begin{matrix} F & P \\ F & \begin{matrix} -3\sqrt{\frac{2}{5}} & 2\sqrt{\frac{3}{5}} \\ 2\sqrt{\frac{3}{5}} & -3^2\sqrt{\frac{1}{5}} \end{matrix} \\ P & \end{matrix}$

(40) $\begin{matrix} G & D & S \\ G & \begin{matrix} -\sqrt{\frac{2 \cdot 3 \cdot 5 \cdot 11}{7}} & 2 \cdot 3 \sqrt{\frac{2 \cdot 3}{7}} \\ 2 \cdot 3 \sqrt{\frac{2 \cdot 3}{7}} & -11 \sqrt{\frac{5}{3 \cdot 7}} & 2\sqrt{\frac{2 \cdot 7}{3}} \end{matrix} \\ D & & \\ S & & \end{matrix}$

(22) $\begin{matrix} G & F & D' & D'' & S \\ G & \begin{matrix} \sqrt{2 \cdot 3 \cdot 5} & 2\sqrt{2 \cdot 3} \\ -\sqrt{2 \cdot 3 \cdot 5} & & 2\sqrt{2 \cdot 5} \end{matrix} \\ F & & & & \\ D' & 2\sqrt{2 \cdot 3} & & & \\ D'' & & -2\sqrt{2 \cdot 5} & 5 & 2^2 \\ S & & & 2^2 & \end{matrix}$

(41) $\begin{matrix} H & G & F & D & P \\ H & \begin{matrix} -\sqrt{\frac{11 \cdot 13}{5}} & 2\sqrt{\frac{3 \cdot 11}{5}} & 2\sqrt{11} \\ -2\sqrt{\frac{3 \cdot 11}{5}} & -2 \cdot 3 \sqrt{\frac{2 \cdot 3 \cdot 11}{5 \cdot 7}} & \sqrt{2 \cdot 3} & 2 \cdot 3 \sqrt{\frac{5}{7}} \end{matrix} \\ G & & & & \\ F & 2\sqrt{11} & -\sqrt{2 \cdot 3} & -2\sqrt{\frac{2 \cdot 7}{3}} & 2 \cdot 3 & 2^2 \cdot 3 \sqrt{\frac{1}{5}} \\ D & & 2 \cdot 3 \sqrt{\frac{5}{7}} & -2 \cdot 3 & -3\sqrt{\frac{3 \cdot 5}{7}} & 2\sqrt{2 \cdot 3} \\ P & & & 2^2 \cdot 3 \sqrt{\frac{1}{5}} & -2\sqrt{2 \cdot 3} & 3^2 \sqrt{\frac{1}{5}} \end{matrix}$

(60) $\begin{matrix} I & G & D & S \\ I & \begin{matrix} -\sqrt{\frac{3 \cdot 5 \cdot 7 \cdot 13}{11}} & 2\sqrt{\frac{3 \cdot 5 \cdot 13}{11}} \\ 2\sqrt{\frac{3 \cdot 5 \cdot 13}{11}} & -3 \cdot 5 \sqrt{\frac{2 \cdot 3 \cdot 5}{7 \cdot 11}} & 2^2 \sqrt{\frac{3 \cdot 11}{7}} \end{matrix} \\ G & & & \\ D & & 2^2 \sqrt{\frac{3 \cdot 11}{7}} & -5\sqrt{\frac{3 \cdot 5}{7}} & 2 \cdot 3 \\ S & & & 2 \cdot 3 & \end{matrix}$

(21) $\begin{matrix} F & D & P \\ F & \begin{matrix} -\sqrt{\frac{2 \cdot 7}{3}} & 2\sqrt{\frac{2 \cdot 7}{3}} & 2\sqrt{\frac{2 \cdot 7}{5}} \\ -2\sqrt{\frac{2 \cdot 7}{3}} & -\sqrt{\frac{5 \cdot 7}{3}} & 2\sqrt{2} \\ 2\sqrt{\frac{2 \cdot 7}{5}} & -2\sqrt{2} & 7\sqrt{\frac{1}{5}} \end{matrix} \\ D & & \\ P & & \end{matrix}$

Table 4

Wigner central force matrix elements between two-nucleon states

$$\frac{3}{23} I_4 + \frac{1}{2^2} I_2 + \frac{3}{23} I_0$$

$$\frac{1}{2} I_3 + \frac{1}{2} I_1$$

$$\frac{3^3}{2^4} I_4 - \frac{3 \cdot 7}{2^3} I_3 + \frac{41}{2^3 \cdot 3} I_2 - \frac{7}{2^3 \cdot 3} I_1 + \frac{5^2}{2^4 \cdot 3} I_0$$

$$0$$

$$-\frac{1}{2^2 \cdot 3} \sqrt{\frac{7}{2}} I_2 + \frac{1}{2 \cdot 3} \sqrt{\frac{7}{2}} I_1 - \frac{1}{2^2 \cdot 3} \sqrt{\frac{7}{2}} I_0$$

$$\frac{7}{2^3} I_3 - \frac{3}{2^2} I_2 + \frac{7}{2^3} I_1$$

$$0$$

$$\frac{7}{2 \cdot 3} I_2 - \frac{1}{3} I_1 + \frac{1}{2 \cdot 3} I_0$$

$$\frac{7}{2^2} I_3 - \frac{5}{2} I_2 + \frac{7}{2^2} I_1$$

$$\frac{3^3 \cdot 7}{2^6} I_4 - \frac{3 \cdot 5 \cdot 7}{2^4} I_3 + \frac{17 \cdot 3^7}{2^5 \cdot 3} I_2 - \frac{5 \cdot 31}{2^4 \cdot 3} I_1 + \frac{13 \cdot 19}{2^6 \cdot 3} I_0$$

$$\frac{\sqrt{5}}{2 \cdot 3} I_2 - \frac{\sqrt{5}}{3} I_1 + \frac{\sqrt{5}}{2 \cdot 3} I_0$$

$$\frac{5}{3} I_2 - \frac{2^2}{3} I_1 + \frac{2}{3} I_0$$