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ОБЪЕДИНЕННЫЙ ИНСТИТУТ ЯДЕРНЫХ ИССЛЕДОВАНИЙ

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DETERMINATION OF PARAMETERS OF THE ROTATIONAL DYNAMICS OF GROUPS NH IN NH CIO AND H O IN H OCIO, BY MEANS OF INELASTIC SCATTERING OF NEUTRONS

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DETERMINATION OF PARAMETERS OF THE ROTATIONAL DYNAMICS OF GROUPS NH IN NHCIQAND H30 IN H30CI0, BY MEANS OF INELASTIC SCATTERING OF NEUTRONS

Объеднасизнай инстрало AGENIAN ECCOCHESS. **БИБЛИОТЕН/**-

## 1. Introduction

The existence of an  $H_30$  group in perchloric acid monohydrate may be considered as an experimentally well-established fact  $^{1-4/}$ . It has also been proved that  $H_30Cl0_4$  and  $NH_4Cl0_4$ , which belong to the orthorombic crystallographic system, possess practically identical lattices of  $Cl0_4 - ions^{/1,3/}$ . This means that the  $H_30$  and  $NH_4$  groups in the two substances are situated in almost the same surrounding conditions.

It is interesting to compare the dynamic properties of the NH<sub>4</sub> and H<sub>3</sub>0 groups in these crystals, paying special attention to the rotational possibilities of the groups. The information which has been obtained in connection with this problem comes from three types of investigation i.e. those made by applying the Raman technique<sup>/5/</sup>, the NMR technique <sup>/6-9/</sup>, and the neutron scattering technique<sup>/10-13/</sup>.

The following results have been obtained for NII, Clo, : a relatively narrow proton resonance line leads to low values of the so-called second moment  $(\Delta H)^2$ . These values are: 2,42  $Oe^2$  for the temperature 20<sup>o</sup>K (Richards and Schaefer  $\frac{7}{7}$  ), and 1,27  $\text{Oe}^2$  and 1,18  $\text{Oe}^2$  for temperatures 70°K and  $298^{\circ}$ K (Ibers  $\binom{6}{6}$ ). The barrier to rotation evaluated for NH<sub>4</sub>ClO<sub>4</sub> from these results is, according to Ibers, 2 kcal/mol; according to Richards and Schaefer on the other hand it is less than 1 kcal/mol. Ibers' evaluation leads to the conclusion that the rotation of the NH, group is hindered, Richards and Schaefer's value suggests freedom of rotation. Neutron investigations give futher arguments strongly supporting a low value of the barrier to rotation; these arguments are: a high value (13 b/ $\mathcal{R}$ ) of the slope of linear dependence between the total neutron scattering cross-section and the neutron wavelength in the subthermal neutron energy region (Rush et al.  $^{/11/}$ ) leads to the conclusion that the rotation of NH, groups in NH, ClO, is free, and to an evaluation of the barrier value for 0,1 - 0,2 kcal/ mol, Also the neutron spectrum obtained by the scattering of cold neutrons from a crystalline NH<sub>4</sub>ClO<sub>4</sub> (Janik et al.  $^{/12/}$ ) shows no sharp peaks from which the existence of discrete torsional oscillation levels could be deduced. Thus, if the high barrier value suggested by Ibers is not taken into account, all othere experiments lead to the conclusion that the rotation of NH, groups in NH, Clo, is free.

Much less clear is the situation in H, OCIO, : although from experiments performed by applying the Raman technique  $\frac{2}{3}$ , 14-16, 5 many frequencies may be identified as those caused by the  $H_{00}^{+1}$  ion, there is a lack of data in the low frequency region (several hundred  $cm^{-1}$ ), in which the H<sub>0</sub> hindered rotatation frequency (i.e. its torsional vibration frequency) could appear, Taylor and Vidale  $\frac{5}{100}$  only suggest that the frequency 995 cm<sup>-1</sup> may have the torsional vibration as its origin; however at the same time they write that this frequency is much too high for this type of vibration. On the other hand, a relatively narrow NMR line in the temperature region from 150°K to room temperature, with a broadening at lower temperatures, led Kakiuchi et al. 9/ to the suggestion that in perchloric acid monohydrate, starting from 150°K, the H.0<sup>+</sup> rotation is relatively free. Neutron experiments, however, do not support this suggestion; the slope of the linear dependence between the total neutron scattering crosssection and the neutron wavelenght is rather low (8 b/A), which might be considered as an argument that the barrier to rotation is c 1,8 kcal/mol (J.M.Janik $^{/13/}$ ). Also the neutron spectrum obtained by the scattering of cold neutrons from a crystalline H.OCIO, (Janik et al.  $^{12/}$ ) shows a sharp peak. which could be interpreted as being caused by a torsional vibration of the group with a corresponding energy  $\hbar\omega = 0.056$  eV.

On the basis of the situation described above we decided to perform for  $NH_4ClO_4$  and  $H_3OClO_4$  additional neutron measurements, this time by applying a so-called "inverted geometry". The reason for this these experiments may be explained as follows. First - in the "inverted geometry" technique neutron energy losses (caused by the molecular degrees of freedom of the sample) are measured, whereas in the "normal geometry" neutron energy gains caused by excited energy levels of the sample are obtained. From this point of view, previous results  $\frac{12}{12}$  may be compared to anti-Stokes lines in the Raman technique. It is interesting to support these results by obtaining the "Stokes-lines" characteristic for the "inverted geometry" technique. Second - the "inverted geometry" technique works much better than the "normal geometry" when applied to samples in lower temperatures. Therefore it was then possible to investigated the dynamics of  $NH_4$  group in  $NH_4ClO_4$  in a wide temperature interval: from liquid nitrogen to room temperature.

#### 2. Experimental technique

The inverted berylium filter arrangement constructed at the pulsed reactor IBR of the Joint Institute of Nuclear Research at Dubna was applied. A detailed description of the arrangement with an analysis of the resolution of the instrument has been published elsewhere  $^{/17/}$ ; at present only the main parameters of the apparatus will be given. In the pulsed reactor pulses of fast neutrons with a repetition ratio 8 pulses per second were produced. The duration of an impulse was c. 36  $\mu$  sec. Neutrons were slowed down in a moderator. This caused an increase in the pulse duration to c. 120  $\mu$  sec. Then, after passing a 20 m long flight path, the neutrons reached the sample, scattered under the 90<sup>°</sup> angle, and were transmitted through a nitrogen cooled Be-filter to a scintillation counter. In this arrangement the counter could register only neutrons whose wavelengh was above the Be-cut-off i.e. neutrons of the most probable energy 3,7 meV. From the flight-path distances: moderator-sample (L<sub>1</sub>) , sample-detector (L<sub>2</sub>), (L<sub>1</sub>>>L<sub>2</sub>) and the measured time distribution of neutrons arriving at the detector at the end of the flight-path, the incident neutron energy may be computed. In this way, for each incident neutron energy it was possible to measure the number of neutrons which in the scattering process obtained energy below the Be-cut-off.

Neutron spectra obtained in this way were however distorted in several narrow energy intervals because of the fact that, besides the main neutron pulses, some additional, much weaker neutron pulses were produced in the reactor with a 10 times greater repetition ratio. These fast neutron pulses naturally caused sharp peaks in periodic intervals of the time of flight scale i.e. in some narrow energy intervals of incident neutrons. The distorted parts of the spectrum were, of course, worthless for analysis of results and therefore in the corresponding plots experimental points were not plotted there and empty places were left.

Measurements for  $NH_4ClO_4$  were made in a thin wall aluminium sample holder, placed in a cryostat. The measurement at  $78^{\circ}K$  was made with the cryostat filled with liquid nitrogen. Then, after the evaporation of the nitrogen, two hour measurement were performed at higher temperatures. Background measurements were made with a Cd shield in the beam of incident neutrons.

Measurements for  $H_{3}$  OClO<sub>4</sub> were performed in another aluminium sample holder with thicker (1 mm) walls. This was necessary be cause of a design solution which assured the hermetic sealing of the sample holder. Background measurements were made with an identical, empty sample holder replasing the sample.

The thicknesses of the two substances were sufficiently small to justify the neglect of the multiple scattering correction.

4

#### 3. Results

Fig.1 presents results for the crystalline  $NH_4ClO_4$  in the temperature region from liquid nitrogen to room temperature. The values plotted are the differences  $I_0 - I_{back}$  between the total number of observed counts and the background (without normalization to constant incident flux). Each plot shows an elastic part in the form of a berylium cut-off and an inelastic part in the form of a broad distribution. All five plots are normalized to the same height of the inelastic part and pushed aside along the vertical axis with an arbitrary scale.

As the main result for  $NH_4ClO_4$  we consider the fact one obtaines in the inelastic part at all temperatures broad distributions without any sharp peaks. These distributions are after all identical. The elastic part decreases regularly with the temperature increase and is very weak at room temperature.

Fig. 2 presents the results obtained for the crystalline  $H_{3}OCIO_{4}$  at room temperature. In the same Fig. is shown for comparison the result for  $NH_{4}CIO_{4}$  at room temperature. This other plot is normalized to the same height of the inelastic parts and pushed aside along the vertical axis with an arbitrary scale.

As the main result for  $H_3OCIO_4$  we consider the fact that the inelastic part shows a sharp peak at incident neutron energy 0,058 eV. The elastic part for  $H_3OCIO_4$  is much higher than for  $NH_4CIO_4$  at the same temperature.

# 4. Discussion

The lack of sharp peaks in the neutron spectra obtained in whole temperature region for NH<sub>4</sub>ClO<sub>4</sub> excludes the existance of torsional vibration levels and therefore supports the hypothesis of free rotation of NH<sub>4</sub> groups. This result agrees with the interpretation of data obtained by Richards and Schaefer<sup>/7/</sup> by NMR method. It agrees also neutron investigations<sup>/10-12/</sup>. As compared with the previous neutron measurements<sup>/12/</sup> the present data are an assential extential extension because the measurements were carried out not only at room temperature but in a broad temperature region including liquid nitrogen temperature.

The result obtained for  $H_80Cl0_4$  at room temperature supports the interpretation of the previous neutron experiments  $^{/12,13/}$  which suggests hindered rotation of the  $H_80$  groups. On the basis of the experimental data of the present work an attempt has been made to present them in a form independent from the neutron scattering dynamics but dependent only on the molecular dynamics of the substance. This was done by a transformation of data from a form in which the scattered

neutron intensity is presented in dependence on the incident neutron energy to a form  $p_{u}(\epsilon)e^{-2W_{H}}$ , where p<sub>H</sub> is the frequency distribution of hydrogen atoms,  $\epsilon = \hbar \omega$  corresponds to the neutron energy loss and  $e^{2W_{\rm H}}$  is the Debye-Waller factor. In this representation is also included normalization to a constant value of incident flux. It must be pointed out that if the torsional vibration peaks are not very sharp they will not appear in this representation in the same places as in the Fig. 2 representation. This explains why the value  $\Delta E = e^{-0.054eV}$ obtained from Fig. 2, although being in agreement (within the limits of experimental errors) with  $\Delta E = \epsilon = 0,056$  eV obtained for a torsional vibration peak in H. OCIO, by Janik et al.  $\frac{12}{12}$ , cannot be directly considered as  $\hbar\omega$ , (where  $\omega_1$  is the frequency of torsional vibration). Fig. 3 presents the function  $p_{\rm H}(\mathbf{r}) e^{-2W_{\rm H}}$ for H, OCIO, calculated on the basis of the data of Fig. 2. This function has a maximum at ( = 0,062 eV which, according to our interpretation, should correspond to the energy  $1\omega_L$ ; this gives for  $\omega_L = 497$  cm<sup>-1</sup>.

It is worth while to note that the Raman frequency 995 cm<sup>-1</sup> obtained by Taylor and Vidale<sup>/5/</sup> was by them interpreted as  $\omega_L$ . If, however, it is not  $\omega_L$  but  $2\omega_L$ , one obtains for  $\omega_L = 497$  cm<sup>-1</sup> in an excellent agreement with our results.

With regards to the suggestion of free rotation of  $H_{s}0$  groups in  $H_{s}0Cl0$ , made on the basis of the NMR measurements 7,9, it should be pointed out that a narrow resonance line cannot be treated simply as a proof of free rotation but only as evidence of a reorientation of the  $H_{s}0$  group which occurs with a frequency greater than c. 50 kc/sec. This reorientation does not exclude the existence of a torsional vibration (with reorientational tunnel transitions).

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7

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8





Fig. 1. The intensity of neutrons scattered by the crystalline NII<sub>4</sub>ClO<sub>4</sub>vs. the incident neutron energy. Temperatures of the sample: 78°K, 94°K, 165°K, and 283°K. Scattering angle - 90°.



Fig. 2. The intensity of neutrons scattered by the crystalline  $NH_4ClO_4$  and the crystalline  $H_3OClO_4$  vs. the incident neutron energy. Both samples at room temperatures. Scattering angle - 90°.

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![](_page_6_Figure_3.jpeg)

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Function  $P_{H}(t)e^{2\frac{W_{H}}{2}}$  for a crystalline  $H_{3}OCIO_{4}$  at room temperature. Fig. 3.