

ОБЪЕДИНЕННЫЙ  
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ДУБНА

A 29

E13-87-309

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**COMBINED SCINTILLATION  
AND TRACK TECHNIQUE  
TO SEARCH FOR  $\beta\beta$ -DECAY**

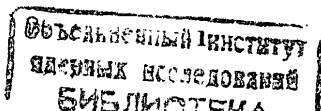
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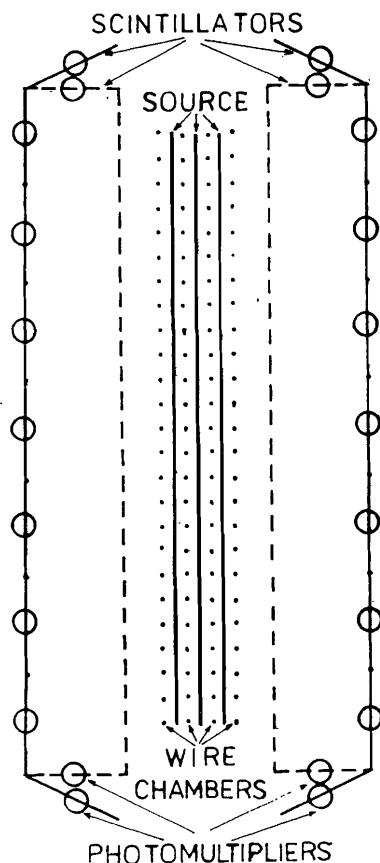
**1987**

Double beta decay is one of the most powerful ways to test conservation of the lepton number. The highest half-life ( $T_{1/2}$ ) limits for the neutrinoless ( $0\nu$ )  $\beta\beta$ -decay were obtained using <sup>11/</sup> germanium solid state detectors. <sup>76</sup>Ge-isotope can give the  $0\nu \beta\beta$ -decay with the total energy of two electrons  $E_0 = 2$  MeV. But this technique does not exclude other methods of investigation with different sources of  $\beta\beta$ -decay. For instance a probability of the  $0\nu \beta\beta$ -decay for <sup>150</sup>Nd is much higher (by two orders <sup>12/</sup>) than for <sup>76</sup>Ge. Besides it is not necessary to use a high resolution detector to search for  $\beta\beta$ -decay with  $2 \nu$  or Majoron ( $0\nu \chi^0$ ), because  $E \neq \text{const}$  for these modes. New limits for the  $0\nu \chi^0$ -mode will have some important consequences in astrophysics <sup>13/</sup>. According to ref. <sup>14/</sup>, a probability of the  $0\nu \chi^0$ -mode is even higher than that of the  $0\nu$ -mode, if Majoron exists. A scintillation counter with a  $\sim 20\%$  energy resolution is an adequate detector for a search for the  $0\nu$  - or  $2\nu$ -modes.

A scintillation spectrometer with a <sup>150</sup>Nd-source was used in the underground experiment <sup>15/</sup>. The Nd<sub>2</sub>O<sub>3</sub>-sample (<sup>150</sup>Nd - 92.5%) was placed between scintillators and a two-dimensional spectrum was measured. The main background in the experiment is due to electrons produced in the scintillators by gamma-rays from natural radioactivity and n-capture.

The time-of-flight and track technique can be used to exclude such a background. This approach is used in the version of the apparatus shown in the Figure 1.





Double beta decay apparatus.

There are three planes of the source ( $\sim 3 \text{ M}^2$ ) in the center of the set-up. Four wire chambers measure particle tracks in the source region. An accuracy  $\Delta \theta = \pm 5^\circ$  for measurement of the angle  $\theta$  between two tracks is quite sufficient. The thickness of each plane is about  $6 \text{ mg/cm}^2$ . It can be  $\text{Nd}_2\text{O}_3$  precipitated on  $6\mu\text{K}$ -mylar (aluminized).  $\text{Nd}$  thickness is  $\sim 4 \text{ mg/cm}^2$  in this case. To simplify preparation of the source precipitation can be done on smaller surfaces ( $\sim 0.1 \text{ M}^2$ ) in succession. Two planes are for  $^{150}\text{Nd}$  ( $M = 75 \text{ g}$ ) and the third plane is a

dummy source. Each chamber measures two coordinates. The average amount of substance ( $\text{Cu}$ ) in each chamber is  $\sim 1 \text{ mg/cm}^2$ .

Scintillation counters are made of long plates. Typical thickness is  $1 \text{ cm}$ , width is  $4 \times 5 \text{ cm}$ . Scintillations in each plate are detected by a small photomultiplier, taking off  $\sim 10\%$  of collected light. Four such plates form a counter with two fast photomultipliers on both edges. A time resolution  $2.35 \sigma \leq 1 \text{ ns}$  is expected for the energy of electrons  $T = 0.5 \times 3 \text{ MeV}$ . The time of flight of electrons between the scintillators is  $2 \times 5 \text{ ns}$  depending on the angle of emission. The energy of two electrons  $E_0 = 3,4 \text{ MeV}$  for the  $0\nu \beta\beta$ -decay of  $^{150}\text{Nd}$ . The most probable energy of one electron is  $0.5 E_0$ . Timing allows determining a position of the scintillation in the counter.

Background electrons with the energy about  $E_0$  produced in the scintillators will have the r.m.s. scattering angle  $\sim 10^\circ$  after traversing the source. A major part of these electrons can be rejected if opening angles  $\theta > 155^\circ$  are excluded. Only 10% of the true events are found at  $\theta > 155^\circ$  in the distribution with the angular correlation  $(1 - \cos \theta)^{1/6}$ . Timing can make the electron background negligible for  $\theta > 155^\circ$ . The main background will be caused by pair production in the source. But there are 90% of pairs and only 20% of the true events for  $\theta < 80^\circ$ . Multiple scattering will not change the angle distribution very much because of small thickness of the source planes.

The apparatus efficiency for  $80^\circ < \theta < 155^\circ$  is about  $1/3$ . If  $T_{1/2} = 10^{22} \text{ y.}$  and  $M = 75 \text{ g}$  one will record  $N = 7$  events of the  $0\nu \beta\beta$ -decay for the measuring time  $t = 1 \text{ y.}$

Let's take data of  $^{151}$  to estimate the background. Events occurred at a rate of  $1.5/\text{h}$  for  $E = 3 \times 3.5 \text{ MeV}$  with the scintillator of mass  $M_s = 6.2 \text{ kg}$ . Taking into account the ratio  $M_{\text{Nd}}/M_s$ , the pair production cross section and efficiency for  $80^\circ < \theta < 155^\circ$

(~ 6%) one can expect ~ 0.5 N of  $e^+e^-$ -pairs for  $t = 1$  y. But the real background can be higher. Active scintillation shielding ( $20 \div 30$  g/cm<sup>2</sup>) can be made for suppressing the background. There will be a high probability of detecting one of the two annihilation gamma-quanta (0.5 MeV).

In the case of the  $0\nu\beta\beta$ -decay two electrons have a total energy  $E = (\sim 0.5-1)E_0$ . Gamma-quanta must have energies  $E_\gamma \geq 1.5 E_0$  to produce pairs with  $E \geq 0.5 E_0$ . It is higher than the natural radioactivity boundary and the background is much lower in this region. The number of single background electrons will be larger than for the  $0\nu$  -mode. But this kind of background can be effectively excluded by the time-of-flight and track technique.

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Received by Publishing Department  
on May 5, 1987.

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E13-87-309

Комбинированная сцинтилляционная и трековая техника для поиска  $\beta\beta$ -распада

Рассматривается вариант сцинтилляционного спектрометра с проволочными камерами для поиска  $\beta\beta$ -распада. Предложена времяпролетная техника для подавления фона. Источник ( $\sim 3$  М<sup>2</sup>), расположенный в центре установки, состоит из трех плоскостей. Четыре проволочные камеры измеряют треки частиц в центре источника. Весь объем заполнен гелием. Сцинтилляционные счетчики изготовлены из длинных пластин. Четыре таких пластины образуют счетчик с двумя быстрыми фотомножителями на обоих торцах. Время пролета электронов между сцинтилляторами составляет 2-5 нс. Если масса источника (<sup>150</sup>Nd) равна 75 г, то может быть измерен период полураспада  $T_{1/2} = 10^{22}$  лет.

Работа выполнена в Лаборатории ядерных проблем ОИЯИ. Препринт Объединенного института ядерных исследований. Дубна 1987

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E13-87-309

Combined Scintillation and Track Technique to Search for  $\beta\beta$ -Decay

A version of the scintillation spectrometer with wire chambers for the search for  $\beta\beta$ -decay is considered. The time-of-flight technique for suppression of a background is proposed. There are three planes of the source ( $\sim 3$  m<sup>2</sup>) in the center of the set-up. Four wire chambers measure particle tracks in the source region. The whole volume is filled with helium. Scintillation counters are made of long plates. Four such plates form a counter with two fast photomultipliers on both edges. The time of flight of electrons between the scintillators is  $2 \div 5$  ns. If a source mass (<sup>150</sup>Nd) is 75 g, a half-life  $T_{1/2} = 10^{22}$  y. can be measured.

The investigation has been performed at the Laboratory of Nuclear Problems, JINR.

Preprint of the Joint Institute for Nuclear Research. Dubna 1987