

ОБЪЕДИНЕННЫЙ
ИНСТИТУТ
ЯДЕРНЫХ
ИССЛЕДОВАНИЙ
ДУБНА

E4-85-953

V.S.Melezhik, J.Wozniak

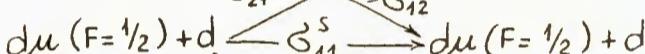
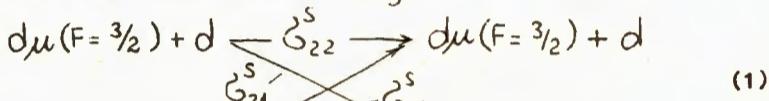
CROSS SECTIONS FOR SCATTERING
OF DEUTERIUM MESIC ATOMS
ON DEUTERIUM NUCLEI

Submitted to "Physics Letters A"

1985

1. To describe kinetics of muon catalyzed fusion which has attracted much interest recently one must know cross sections and rates of various processes involving mesic atoms (cf., e.g., ref. 1).

In this paper cross sections σ_{ij}^S for reaction



are calculated. Here σ_{11}^S and σ_{22}^S are elastic-scattering cross sections of $d\mu$ -atoms in para and ortho states ($F=1/2$ and $F=3/2$, respectively) on deuterium nuclei; σ_{12}^S and σ_{21}^S are cross sections for transitions between those states; S is the total spin of the $d\mu + d$ system; F is the spin of the $d\mu$ -atom. Considered energy region is $\varepsilon \leq 50$ eV*. The ground states of $d\mu$ -atoms are only considered, i.e., two lower reaction channels can be open: channel 1 with the $d\mu$ -atom spin $F=1/2$ and channel 2 with $F=3/2$, the gap between their thresholds being equal to the hyperfine splitting energy of the $d\mu$ -atom $\Delta \varepsilon^{hfs} = 0.049$ eV.

When describing reaction (1) we took into account the exchange interaction where the total spin $S = F + S_d$ was conserved. The spin may take values $S=1/2, 3/2$ for the transitions $F=3/2 \leftrightarrow F=1/2$. Spin-spin and spin-orbital interactions which change S and whose values are several orders of magnitude lower than that of the exchange interaction /2/ are not considered here. Therefore the elastic scattering ($F=3/2 \rightarrow F=3/2$) is only possible for the $S=5/2$ state.

Earlier /2-4/ reaction (1) was studied within the limits of low collision energy $\varepsilon \lesssim \Delta \varepsilon^{hfs}$.

*) The initial energy of deuterium mesic atoms produced in the ground state does not exceed ~ 2 eV for D_2 and D_2+T_2 mixture and ~ 50 eV for D_2+H_2 .

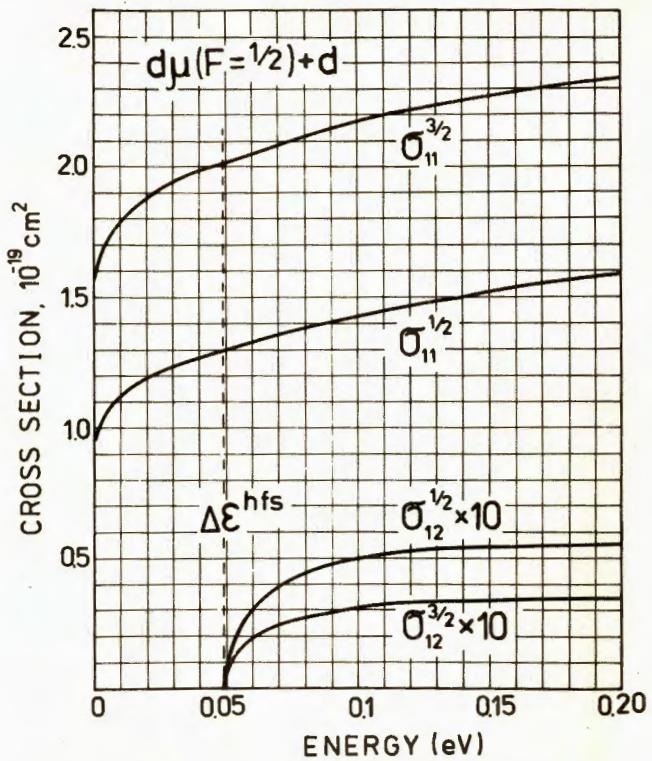


Fig. 1. Cross sections σ_{11}^S for elastic scattering of $d\mu$ -atoms in the para state ($F=1/2$) on deuterium nuclei and for para-ortho transition σ_{12}^S . Results for low collision energies (S-wave only) are shown.

2. The cross sections have been calculated in the adiabatic representation /5/ where the problem of slow collisions in a three-body system is reduced to the multichannel scattering problem /6,7/ with a great number of closed channels, $N \sim 100$:

$$\frac{d^2}{dR^2} y_i^{(r)}(R) + \left(k_i^2 - \frac{J(J+1)}{R^2}\right) y_i^{(r)}(R) - \sum_{j \neq i}^N U_{ij}^S(R) y_j^{(r)}(R) = 0 \quad (2)$$

$i = 1, \dots, N$; $r = 1, \dots, \gamma$; $\gamma = 1, 2$ is the number of open channels. Here k_i is the momentum in the channel ($k_i^2 \geq 0$ at $i \leq \gamma$ and $k_i^2 < 0$ at $i > \gamma$); J is the angular momentum of the system; the wave

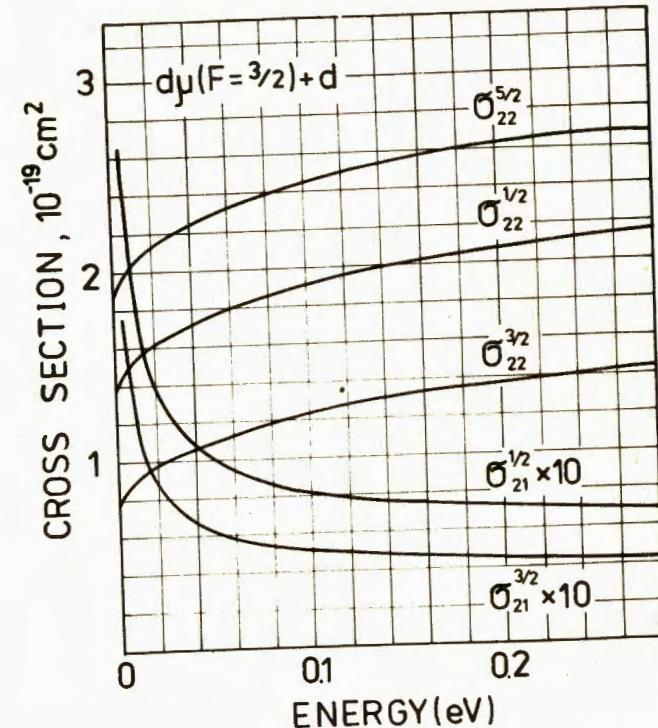


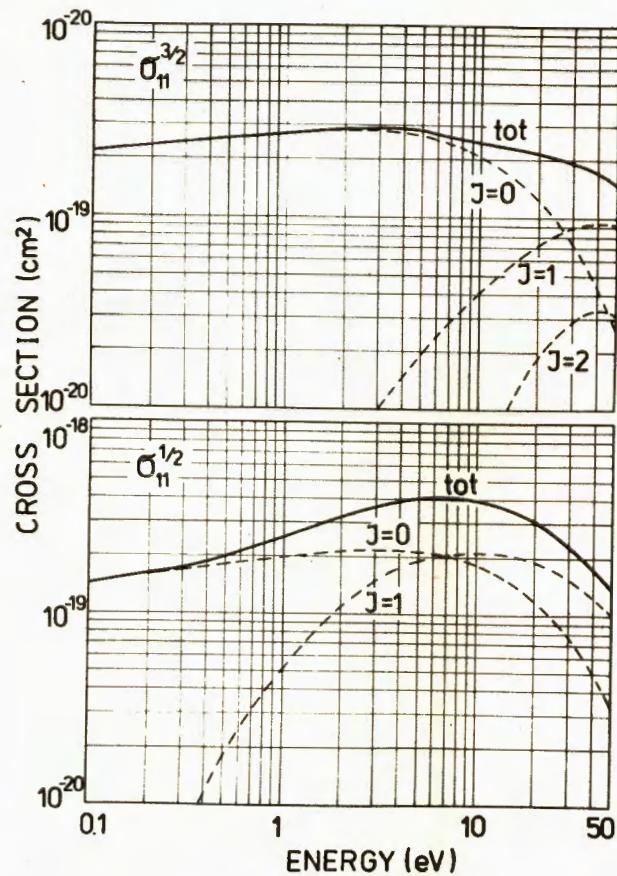
Fig. 2. Elastic scattering s-wave cross sections σ_{22}^S of ortho- $d\mu$ -atoms ($F=3/2$) and for ortho-para transitions σ_{21}^S in the low energy region.

functions $y_i^{(r)}(R) \equiv y_{1S6,F=1/2}^{(r)}(R), y_2^{(r)}(R) \equiv y_{1S6,3/2}^{(r)}(R), \dots$ of the relative nuclear motion are regular at $R=0$ and correspond to subsystems $d\mu_{1S}$ ($F=1/2$)+d and $d\mu_{1S}$ ($F=3/2$)+d in the asymptotic region $R \rightarrow \infty$ where they have the form:

$$\left\{ y_i^{(r)}(R) \right\} = \begin{cases} y_1^{(1)}, \dots, y_\gamma^{(\gamma)} \\ \vdots \\ y_N^{(1)}, \dots, y_N^{(\gamma)} \end{cases} = \begin{cases} \hat{j}_j(R) - \hat{n}_j(R) \hat{T} & \text{for } i \leq \gamma \\ \exp(-ik_i|R) & \text{for } i > \gamma \end{cases} \quad (2a)$$

Here $\hat{j}_j(R) = j_j(k_i R) \delta_{ij}$, $\hat{n}_j(R) = n_j(k_i R) \delta_{ij}$ and $j_j(k_i R), n_j(k_i R)$ are Bessel's spherical functions; \hat{T} is the reaction matrix to be found. Potentials $U_{ij}^S(R)$ are formed of the effective potentials of

Fig. 3. Total and partial (s, p, d -waves) elastic scattering cross sections σ_{ij}^S of para- $d\mu$ -atoms for $0.1 \leq E \leq 50$ eV.



the three-body problem $U_{ij}(R)$ [5]. The spin is introduced according to the rules [2] by means of the transformations $A(S, J)$:

$$U^S = A(S, J) U A^{-1}(S, J), \text{ where } U = \begin{pmatrix} U_{igjj}, U_{igju} \\ U_{iujj}, U_{iuju} \end{pmatrix}$$

$$A(\frac{1}{2}; J=0, 2) = \frac{1}{\sqrt{3}} \begin{pmatrix} -1 & \sqrt{2} \\ \sqrt{2} & 1 \end{pmatrix}; \quad A(\frac{3}{2}; J=0, 2) = \frac{1}{\sqrt{16}} \begin{pmatrix} -\sqrt{5} & 1 \\ 1 & \sqrt{5} \end{pmatrix}$$

$$A(\frac{1}{2}; J=1) = \frac{1}{\sqrt{3}} \begin{pmatrix} \sqrt{2} & 1 \\ -1 & \sqrt{2} \end{pmatrix}; \quad A(\frac{3}{2}; J=1) = \frac{1}{\sqrt{16}} \begin{pmatrix} 1 & \sqrt{5} \\ \sqrt{5} & 1 \end{pmatrix}$$
(3)

$$U^{\frac{5}{2}} = \begin{cases} U_{igjj} & \text{for } J=0, 2 \\ U_{iuju} & \text{for } J=1. \end{cases}$$

Here we take into account the wave-functions symmetry of three spin particles $d\mu + d$ and their spin-spin contact interaction [27].

The method from ref. 7 is used to solve the problem (2)-(3). The matrix cross section

$$\sigma_{ij}^S = \frac{\pi}{k_i^2} \left| \delta_{ij} - \left\{ (1+i\hat{T})(1-i\hat{T})^{-1} \right\}_{ij} \right|^2 \quad (4)$$

is found by the calculated reaction matrix T . Preliminary results of calculations were given in ref. 8.

3. In Figs 1-5 one can see the results of calculations of cross sections σ_{ij}^S versus collision energy $E = k^2/2M$ (M is the reduced mass of $d\mu + d$ system). Cross sections for elastic scattering of $d\mu$ -atoms in para states σ_{11}^S and for transitions to ortho states σ_{12}^S are shown in Fig. 1 for $E \leq 0.2$ eV. Analogous quantities $\sigma_{22}^S, \sigma_{21}^S$ for ortho state are presented in Fig. 2 for $E \leq 0.28$ eV. The S -wave cross sections are only shown for possible values of total spin S , as the contribution of higher partial waves can be neglected for the considered energy region (except for σ_{21}^S case shown more completely in Fig. 5 for $E \geq 0.1$ eV). Noteworthy is the threshold behaviour at $E = \Delta E^{hfs} = 0.049$ eV (Fig. 1).

In the limit $E \rightarrow 0$ the T -matrix for $J=0$ can be presented in the form $T_{ij} = -a_{ij} \cdot \sqrt{k_i k_j}$ via introducing the scattering lengths a_{ij} . In the Table we present our best values of scattering lengths obtained at $N=92$ and the results at $N=2$ for comparison with the results of previous calculations.

Table. Scattering lengths for $d\mu + d \rightarrow d\mu + d$ processes
(in units $a_\mu = 2.56 \cdot 10^{-11}$ cm)

| S | N | a_{11} | $a_{12} = a_{21}$ | a_{22} |
|---------------|----|--|--|--|
| | | this ref. 3 work a) ref. 4 b) | this ref. 3 work a) ref. 4 b) | this ref. 3 work a) ref. 4 b) |
| $\frac{1}{2}$ | 92 | 3.40 | -0.95 | 4.07 |
| | 2 | 3.76 3.5 3.61 | -1.09 -1.1 -0.92 | 4.53 4.7 4.26 |
| $\frac{3}{2}$ | 92 | 4.40 | -0.75 | 3.06 |
| | 2 | 4.92 5.1 4.58 | -0.86 -0.86 -0.73 | 3.38 3.5 3.27 |
| $\frac{5}{2}$ | 92 | | | 4.74 |

a) two-level approximation of adiabatic approach by the method of phase functions

b) improved two-level approximation.

The known formula $k_i/T_{ii} = -1/a_{ii} + 3\pi M/(2a_{ii}^2) \cdot k_i^2 + 3M/a_{ii} \ln(9Mk_i^2/32)$

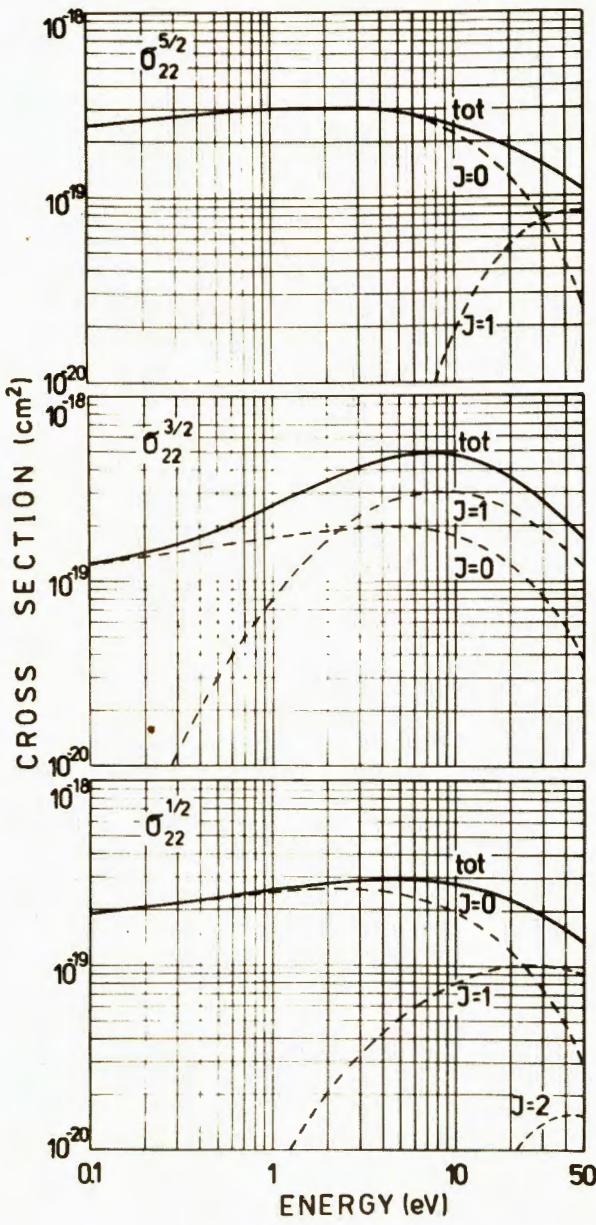


Fig. 4. Total and partial (s, p, d-waves) elastic scattering cross sections $\sigma_{22}^s(\varepsilon)$ of ortho- $d\mu$ -atoms for $0.1 \leq \varepsilon \leq 50$ eV.

can be used to obtain, with an accuracy $\leq 2\%$, the diagonal elements of T-matrix at $\varepsilon \lesssim 0.04$ eV.

Figs 3-5 show the total as well as partial (s, p and d-wave) cross sections for collision energies $0.1 \leq \varepsilon \leq 50$ eV. For the energies 25 eV the cross sections reach a unitary limit for $J=0$, since phase-shifts $\delta_{11} = \arctan T_{11}$ jump as much as π there. A similar behaviour of δ_{ij}^s ($J=1$) is observed at ~ 35 eV.

Now one can use the results obtained for δ_{ij}^s to determine total cross sections of $d\mu$ -atoms for any known population of hf-states.

Accuracy of the calculations of δ_{ij}^s is $\leq 5\%$ for the energy region $\varepsilon \lesssim 1$ eV and is a little worse ($\sim 10\%$) for $\varepsilon \gtrsim 2$ eV. It is determined by the errors of numerical integration of the system of equations (2) and by convergence of solution in the number N of functions used in expanding the wave function of the three-body system in the basis of solutions of a two-centre problem [7]. The numerical study has shown that the prescribed accuracy is achieved at $N=92$ ($J=0$) or $N=136$ ($J=1, 2$) and for the set of nodes over R : $0.1(0.1)20(1)R_{\max}$ used for numerical integration. The quadrature formulae of $O(h^2)$ order have been used, and R_{\max} has been chosen to achieve the assumed accuracy.

4. Using the cross sections δ_{12}^s and δ_{21}^s we have calculated the temperature dependence of the rates of ortho-para transitions of $d\mu$ -atoms:

$$\langle \lambda_{21} \rangle = \frac{1}{6} \int_0^\infty (\lambda_{21}^{1/2} + 2\lambda_{21}^{3/2}) \varphi(\varepsilon, T) d\varepsilon$$

$$\langle \lambda_{12} \rangle = \frac{1}{3} \int_0^\infty (\lambda_{12}^{1/2} + 2\lambda_{12}^{3/2}) \varphi(\varepsilon, T) d\varepsilon$$

$\lambda_{hf} = \langle \lambda_{21} \rangle - \frac{1}{2} \langle \lambda_{12} \rangle$,
where $\lambda_{ij}^s = N_0 U G_{ij}^s$, $N_0 = 4.25 \cdot 10^{22} \text{ cm}^{-3}$ is the liquid hydrogen density, $\varphi(\varepsilon, T)$ is the Maxwellian distribution of $d\mu$ -atoms at temperature T . The calculations have been made under the assumption of the statistical distribution of F-states ($n_{F=3/2} : n_{F=1/2} = 2$). The results are shown in Fig. 6 together with the experimental value at $T = 34$ K [9]. An interesting behaviour of ortho-to-para transition in a deuterium target is the decrease of effective rate $\lambda_{hf}(T)$ with growing temperature. This is due to the fact that $\Delta \varepsilon_{hfs}$ is of the same order of magnitude as the Maxwellian distribution width.

5. The results obtained can be employed for the calculation of kinetics of μ -atomic and μ -molecular processes in hydrogen isotope mixture.

However, when comparing the calculated elastic cross sections of mesic atoms with the experimental ones [10], one should take into

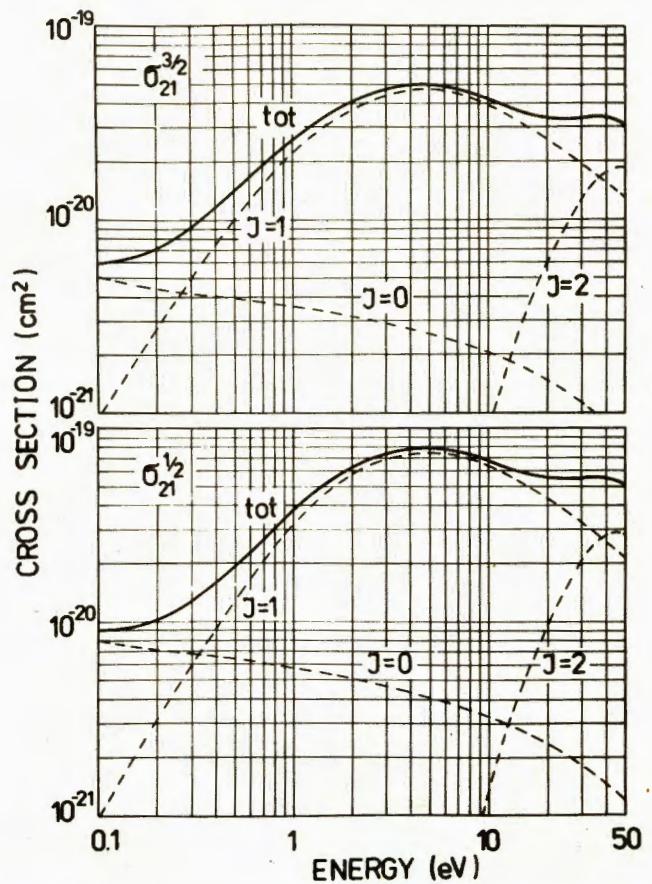


Fig. 5. Total and partial (*s*, *p*, *d*-waves) cross sections $G_{21}^s(\varepsilon)$ for ortho-para transitions. Cross sections $G_{12}(\varepsilon)$ for para-ortho transition can be obtained by using relations: $G_{12}(\varepsilon) = (1 - \Delta\varepsilon^{hf}/\varepsilon) G_{21}^s(\varepsilon)$; $\varepsilon = \varepsilon' + \Delta\varepsilon^{hf}$.

account the electron screening /11/ and the molecular structure effects. They can significantly affect the cross sections at low collision energies $\varepsilon \lesssim 0.1$ eV. On the other hand, their influence on the ortho-para transitions is insignificant, as such transition processes occur in the region of mesic-atom interactions with the nuclei, i.e. at distances $R \sim 2-3 a_0$, where the electron screening effects are weak. The obtained temperature dependence of the spin-

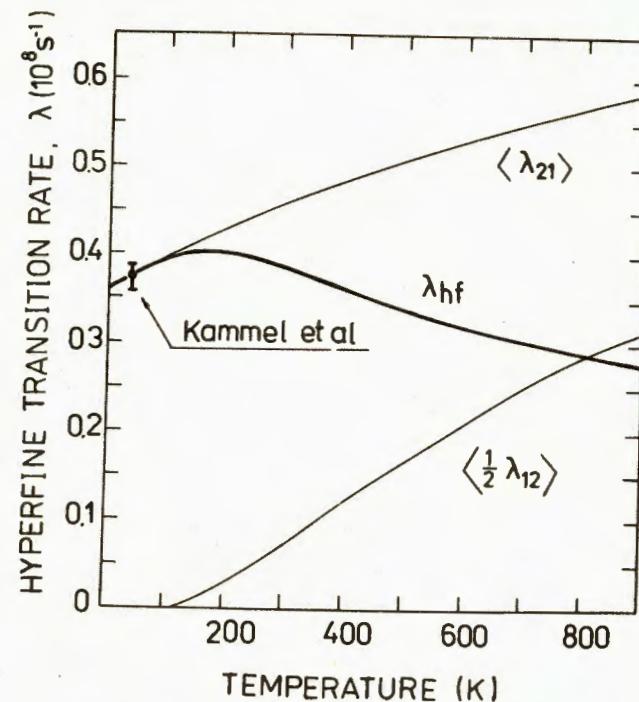


Fig. 6. Temperature dependence of the Maxwell averaged ortho-para and para-ortho transition rates (λ_{21} and λ_{12} , respectively) in a deuterium target. Effective spin flip-rate $\lambda_{hf}(T)$ (assuming the population ratio of hf-states $N_{F=3/2} : N_{F=1/2} = 2$) as well as the experimental value at $T=34$ K /10/ are also shown. Results are normalized to the liquid hydrogen density.

flip rate for $D\mu$ -atoms can be measured since it appears in the experimentally accessible region.

The authors express their gratitude to Professor L.I.Ponomarev for his interest in the work and useful discussions; to Drs. M.P. Faifman, S.I.Vinitsky and A.Adamczak for help and discussions; to Professor V.P.Dzhelepov and Professor V.G.Zinov for support.

References

1. Proc.of the III Int.Conf.on Emerging Nuclear Energy Systems, Helsinki, 7-9 June 1983, Atomkernenergie-Kerntechnik, 43 (1983);

- Contrib. Muon-Catalyzed Fusion Workshop, Jackson, 7-8 June, 1984
2. S.S.Gerstein, Zh.Eksp.Teor.Fiz. 40 (1961) 698 (Sov.Phys.JETP 13 (1961) 488)
 3. A.V.Matveenko and L.I.Ponomarev. Zh.Eksp.Teor.Fiz. 59 (1970) 1593 (Sov.Phys. JETP 32 (1971) 871);
A.V.Matveenko, L.I.Ponomarev and M.P.Faifman, Zh.Eksp.Teor.Fiz. 68 (1975) 437 (Sov.Phys.JETP 41 (1975) 212)
 4. L.I.Ponomarev, L.N.Somov and M.P.Faifman. Yad.Fiz. 29 (1979) 133 (Sov.J.Nucl.Phys. 29 (1980) 67)
 5. S.I.Vinitsky and L.I.Ponomarev. Fiz.Elem.Chastits A.Yadra 13 (1982) 1336 (Sov.J.Part.Nucl. 13 (1982) 557)
 6. V.S.Meleshik, L.I.Ponomarev and M.P.Faifman. Zh.Eksp.Teor.Fiz. 85 (1983) 434 (Sov.Phys.JETP 58 (1983) 254)
 7. V.S.Meleshik. Preprint JINR P4-84-643 (Dubna, 1984); to be published in J.Comp.Phys.
 8. V.S.Meleshik, L.I.Ponomarev and J.Wozniak, Contribution to the X-th Int.Conf.on Part.and Nucl.Book of Abstracts, Heidelberg, p. 140 (1984)
 9. P.Kammel et al. Phys.Rev. A28 (1983) 2611
 10. V.P.Dzhelepov et al. Zh.Eksp.Teor.Fiz. 49 (1965) 393 (Sov.Phys. JETP 22 (1966) 275); A.Alberigi et al. Nuovo Cim. 47B (1967) 72; A.Bertin et al. Phys.Lett. 85B (1979) 458
 11. A.B.Kravtsov, A.I.Mikhailov and N.P.Popov. Preprint LINP 1037 (Leningrad, 1985).

WILL YOU FILL BLANK SPACES IN YOUR LIBRARY?

You can receive by post the books listed below. Prices - in US \$,
including the packing and registered postage

| | | |
|---------------|--|-------|
| D1,2-82-27 | Proceedings of the International Symposium on Polarization Phenomena in High Energy Physics. Dubna, 1981. | 9.00 |
| D2-82-568 | Proceedings of the Meeting on Investigations in the Field of Relativistic Nuclear Physics. Dubna, 1982 | 7.50 |
| D3,4-82-704 | Proceedings of the IV International School on Neutron Physics. Dubna, 1982 | 12.00 |
| D11-83-511 | Proceedings of the Conference on Systems and Techniques of Analytical Computing and Their Applications in Theoretical Physics. Dubna, 1982. | 9.50 |
| D7-83-644 | Proceedings of the International School-Seminar on Heavy Ion Physics. Alushta, 1983. | 11.30 |
| D2,13-83-689 | Proceedings of the Workshop on Radiation Problems and Gravitational Wave Detection. Dubna, 1983. | 6.00 |
| D13-84-63 | Proceedings of the XI International Symposium on Nuclear Electronics. Bratislava, Czechoslovakia, 1983. | 12.00 |
| E1,2-84-160 | Proceedings of the 1983 JINR-CERN School of Physics. Tabor, Czechoslovakia, 1983. | 6.50 |
| D2-84-366 | Proceedings of the VII International Conference on the Problems of Quantum Field Theory. Alushta, 1984. | 11.00 |
| D1,2-84-599 | Proceedings of the VII International Seminar on High Energy Physics Problems. Dubna, 1984. | 12.00 |
| D17-84-850 | Proceedings of the III International Symposium on Selected Topics in Statistical Mechanics. Dubna, 1984. /2 volumes/. | 22.50 |
| D10,11-84-818 | Proceedings of the V International Meeting on Problems of Mathematical Simulation, Programming and Mathematical Methods for Solving the Physical Problems, Dubna, 1983 | 7.50 |
| | Proceedings of the IX All-Union Conference on Charged Particle Accelerators. Dubna, 1984. 2 volumes. | 25.00 |
| D4-85-851 | Proceedings on the International School on Nuclear Structure. Alushta, 1985. | 11.00 |

Orders for the above-mentioned books can be sent at the address:
Publishing Department, JINR
Head Post Office, P.O.Box 79 101000 Moscow, USSR

Received by Publishing Department
on December 27, 1985.

**Вниманию организаций и лиц, заинтересованных в получении
публикаций Объединенного института ядерных исследований**

Принимается подписка на препринты и сообщения Объединенного института ядерных исследований.

Установлена следующая стоимость подписки на 12 месяцев на издания ОИЯИ, включая пересылку, по отдельным тематическим категориям:

| ИНДЕКС | ТЕМАТИКА | Цена подписки на год |
|--------|--|-------------------------|
| 1. | Экспериментальная физика высоких энергий | 10 р. 80 коп. |
| 2. | Теоретическая физика высоких энергий | 17 р. 80 коп. |
| 3. | Экспериментальная нейтронная физика | 4 р. 80 коп. |
| 4. | Теоретическая физика низких энергий | 8 р. 80 коп. |
| 5. | Математика | 4 р. 80 коп. |
| 6. | Ядерная спектроскопия и радиохимия | 4 р. 80 коп. |
| 7. | Физика тяжелых ионов | 2 р. 85 коп. |
| 8. | Криогенника | 2 р. 85 коп. |
| 9. | Ускорители | 7 р. 80 коп. |
| 10. | Автоматизация обработки экспериментальных данных | 7 р. 80 коп. |
| 11. | Вычислительная математика и техника | 6 р. 80 коп. |
| 12. | Химия | 1 р. 70 коп. |
| 13. | Техника физического эксперимента | 8 р. 80 коп. |
| 14. | Исследования твердых тел и жидкостей ядерными методами | 1 р. 70 коп. |
| 15. | Экспериментальная физика ядерных реакций при низких энергиях | 1 р. 50 коп. |
| 16. | Дозиметрия и физика защиты | 1 р. 90 коп. |
| 17. | Теория конденсированного состояния | 6 р. 80 коп. |
| 18. | Использование результатов и методов фундаментальных физических исследований в смежных областях науки и техники | 2 р. 35 коп. |
| 19. | Биофизика | 1 р. 20 коп. |

Подписка может быть оформлена с любого месяца текущего года.

По всем вопросам оформления подписки следует обращаться в издательский отдел ОИЯИ по адресу: 101000 Москва, Главпочтамт, п/я 79.

Возняк Я., Мележик В.С.

Сечения рассеяния мезоатомов дейтерия на ядрах дейтерия

E4-85-953

Вычислены сечения упругого рассеяния мезоатомов дейтерия в орто- и параконстанциях на ядрах дейтерия в интервале энергий 0-50 эВ, а также сечения переворота спина мезоатома в таких столкновениях. Расчеты выполнены в адиабатическом представлении, в котором исходная задача о медленных столкновениях в системе трех частиц сведена к многоканальной задаче рассеяния. Получена температурная зависимость скорости орто-пара переходов $\lambda_{\text{hf}}(T)$ мезоатомов в дейтериевой мишени. Вычисленная скорость $37,2 \cdot 10^6 \text{ s}^{-1}$ при 34 К хорошо согласуется с экспериментальным значением $\lambda_{\text{exp}}(34 \text{ K}) = (37.0 \pm 0.74) \cdot 10^6 \text{ s}^{-1}$ (группа Вена-SIN, 1983г.).

Работа выполнена в Лаборатории ядерных проблем ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна 1985

Melezhik V.S., Wozniak J.

Cross Sections for Scattering of Deuterium Mesic Atoms
on Deuterium Nuclei

E4-85-953

Cross sections have been calculated for elastic scattering of deuterium mesic atoms in ortho and para states on deuterium nuclei in the 0-50 eV energy range; cross sections for spin-flip have also been calculated. The calculations have been performed in the adiabatic representation where the initial problem of slow collisions in a three-body system is reduced to the multi-channel scattering problem. The temperature dependence of the ortho-para transition rate $\lambda_{\text{hf}}(T)$ for $d\mu$ -atoms is obtained for a deuterium target. The calculated rate $\lambda_{\text{hf}}(34 \text{ K}) = 37.2 \cdot 10^6 \text{ s}^{-1}$ is in a good agreement with the recent experimental value, $\lambda_{\text{exp}}(34 \text{ K}) = (37.0 \pm 0.74) \cdot 10^6 \text{ s}^{-1}$ (Vienna-SIN group, 1983).

The investigation has been performed at the Laboratory of Nuclear Problems, JINR.

Preprint of the Joint Institute for Nuclear Research. Dubna 1985